

Optical Characterisation of Materials

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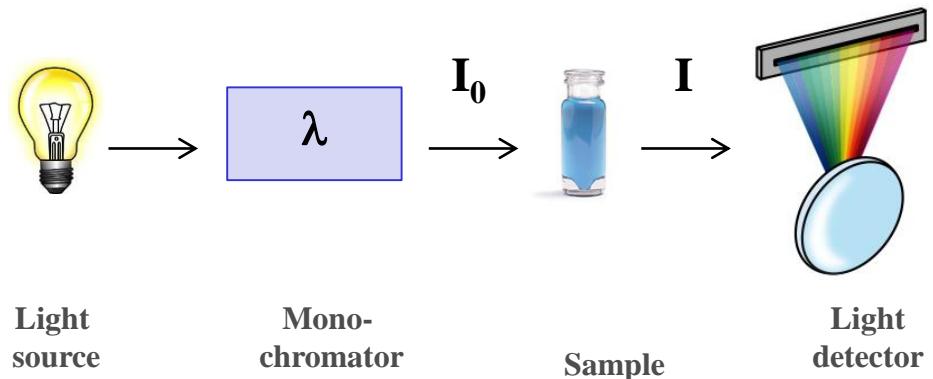
3. Experimental Setup

4. Methods of Optical Material Characterization

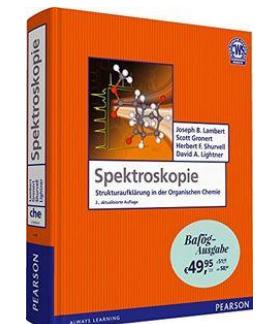
4.1 Absorption Spectroscopy

4.2 Luminescence Spectroscopy

4.3 Reflection Spectroscopy (Remission spectroscopy)



Literature

- **B. Schröder, J. Rudolph**
Physikalische Methoden in der Chemie, VCH, 1985
- **H. Naumer, W. Heller**
Untersuchungsmethoden in der Chemie, Georg Thieme Verlag, 1990
- **E.A.V. Ebsworth, D.R.H. Ranklin, S. Cradock,**
Structural Methods in Inorganic Chemistry, Blackwell Scient. Publications, 1991
- **W. Schmidt**
Optische Spektroskopie, VCH, 1993
- **M. Hesse, H. Meier, B. Zeeh**
Spektroskopische Methoden in der organischen Chemie, VCH, 2005
- **J.B. Lambert, S. Gronert, H.F. Shurvell, D.A. Lightner**
Spektroskopie - Strukturaufklärung in der organischen Chemie, Pearson, 2012

Relevance of Optical Spectroscopy

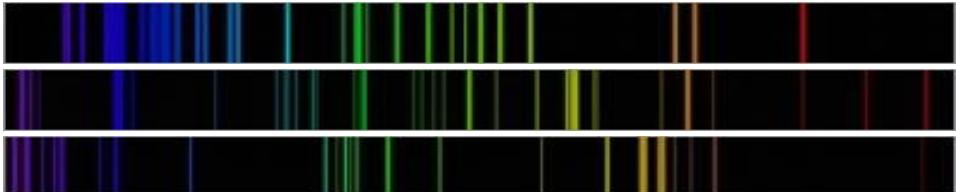
Optical Radiation: UV, visible, and IR radiation ~ 10 nm - 1000 µm (THz - PHz)

- Fingerprint of the elements



As
Pb
Hg

(Astro)physics, analytical chemistry (OES)



- Determination of elements or molecules

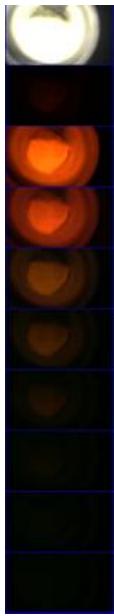
Analytical chemistry: Photometry, AAS, OES

- Characterisation of molecules

Physics, chemistry, biochemistry

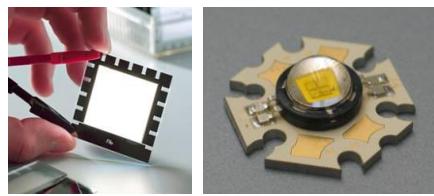
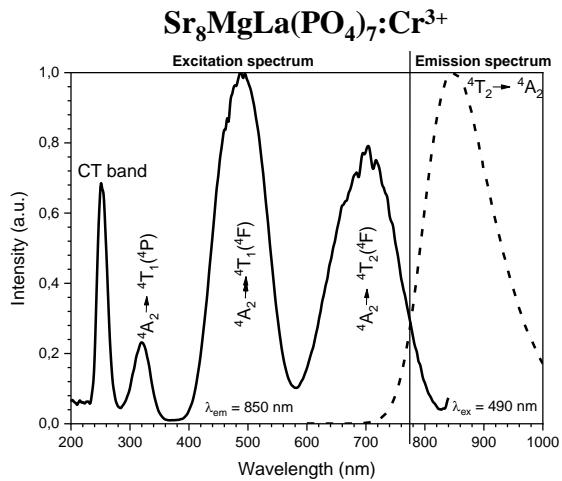
- Characterisation of materials

- Pigments
- Luminescent materials
- Scintillators



- Characterisation of devices

- Light sources
- IR and UV sources
- Displays



1. Methods of Material Characterization

1. Chemical composition

Elementary analysis	Main and secondary components
X-ray fluorescence analysis	Main and secondary components
Atomic absorption spectroscopy	Trace elements, dopants
Powder diffractometry	Phase composition
Infrared spectroscopy (IR)	Detection of functional groups
Raman spectroscopy	Detection of functional groups
Differential thermal analysis	Phase transformations, crystallization water

2. Morphological characterization (micro and nanostructure)

(Dynamic) Light scattering	Particle size distribution
Optical microscopy	Particle morphology
Scanning electron microscopy	Particle size distribution, surface structure
Transmission electron microsc.	Crystallinity, texture
BET-isotherm	Specific surface area

1. Methods of Material Characterization

3. Structural characterization

Powder diffractometry	Phase identification, crystal system, space group
X-ray structure analysis	Space group, atomic coordinates
X-ray absorption spectroscopy	Chemical environment of heavy atoms
NMR spectroscopy	Molecular structure
Infrared spectroscopy (IR)	Molecule identification Arrangement of functional groups
Raman spectroscopy	Arrangement of functional groups
Resonance-Raman spectroscopy	Arrangement of functional groups

4. Surface characterization

Electrosonic Amplitude (ESA)	Surface charge, isoelectric point
Mass spectroscopy (SIMS)	Adsorption of organic molecules at surfaces
X-ray photoelectron spectroscopy (ESCA/XPS)	Surface composition

1. Methods of Material Characterization

5. Determination of physical properties

Electronic properties

Absorption spectroscopy

EPR spectroscopy

Mössbauer spectroscopy

Ultraviolet photoelectron spectroscopy (UPS)

Electronic ground state

Oxidation states, electronic ground state

Oxidation states, coordination geometry

Band structure

Optical properties

Absorption spectroscopy

Reflectance spectroscopy

Luminescence spectroscopy

Absorption coefficient, electronic ground state

Band gap, color point

Excitation and emission spectrum

Color point

Quantum efficiency

Luminous efficacy

Thermal quenching



Magnetic properties

Susceptibility measurement

NMR spectroscopy

Dia-, para-, ferro-, ferri-, antiferromagnetism

Ground state

2. Physical Basis

2.1 Classification of Optical Spectroscopy

2.2 The Electromagnetic Spectrum

2.3 Time Scale of Physical Processes

2.4 Interaction of Light and Matter

Types of Electromagnetic Interaction

Electronic States in Atoms

Electronic States in Solids

Reflection

Absorption

Refraction

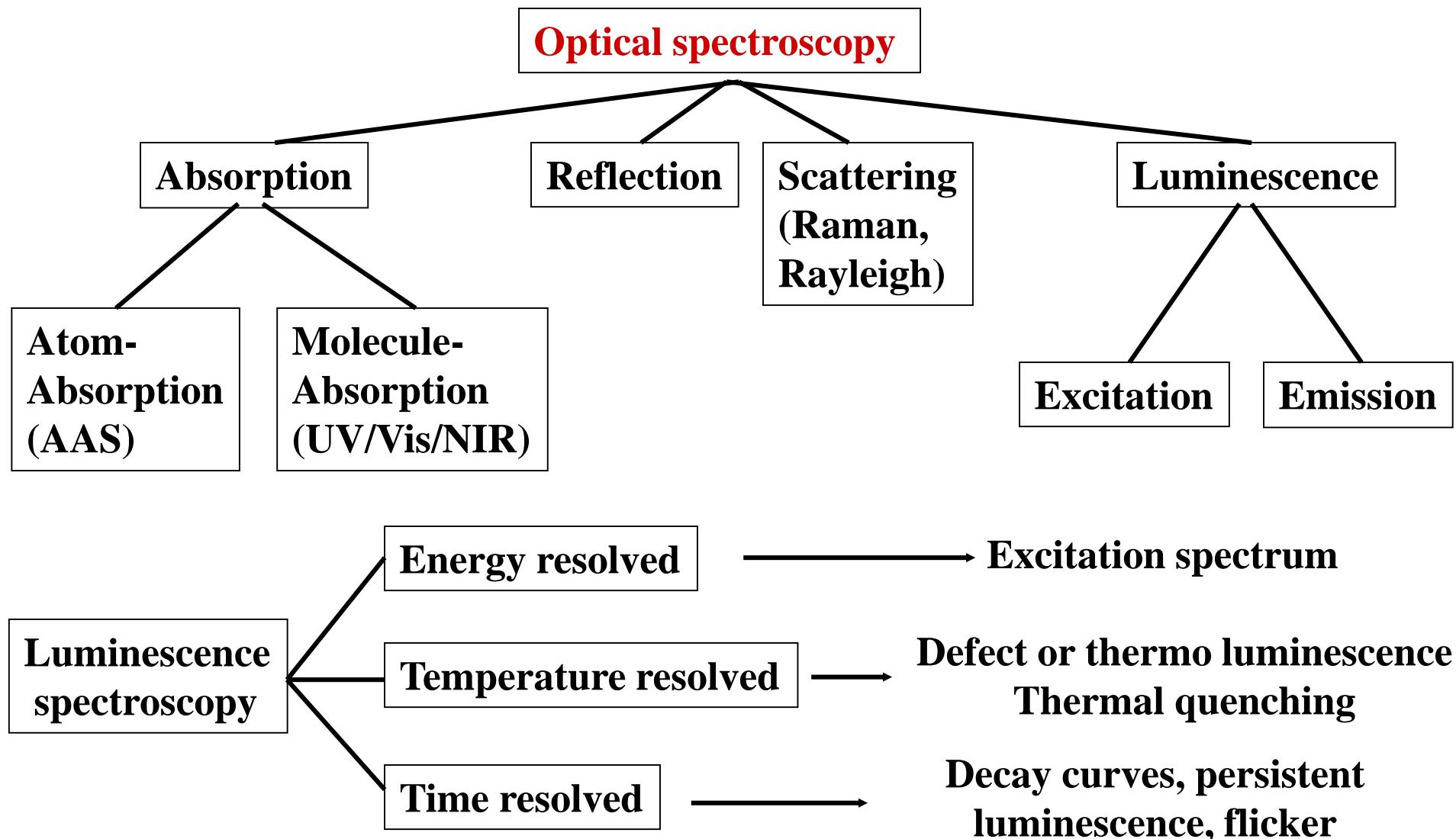
Luminescence

2.5 Radiometric Quantities

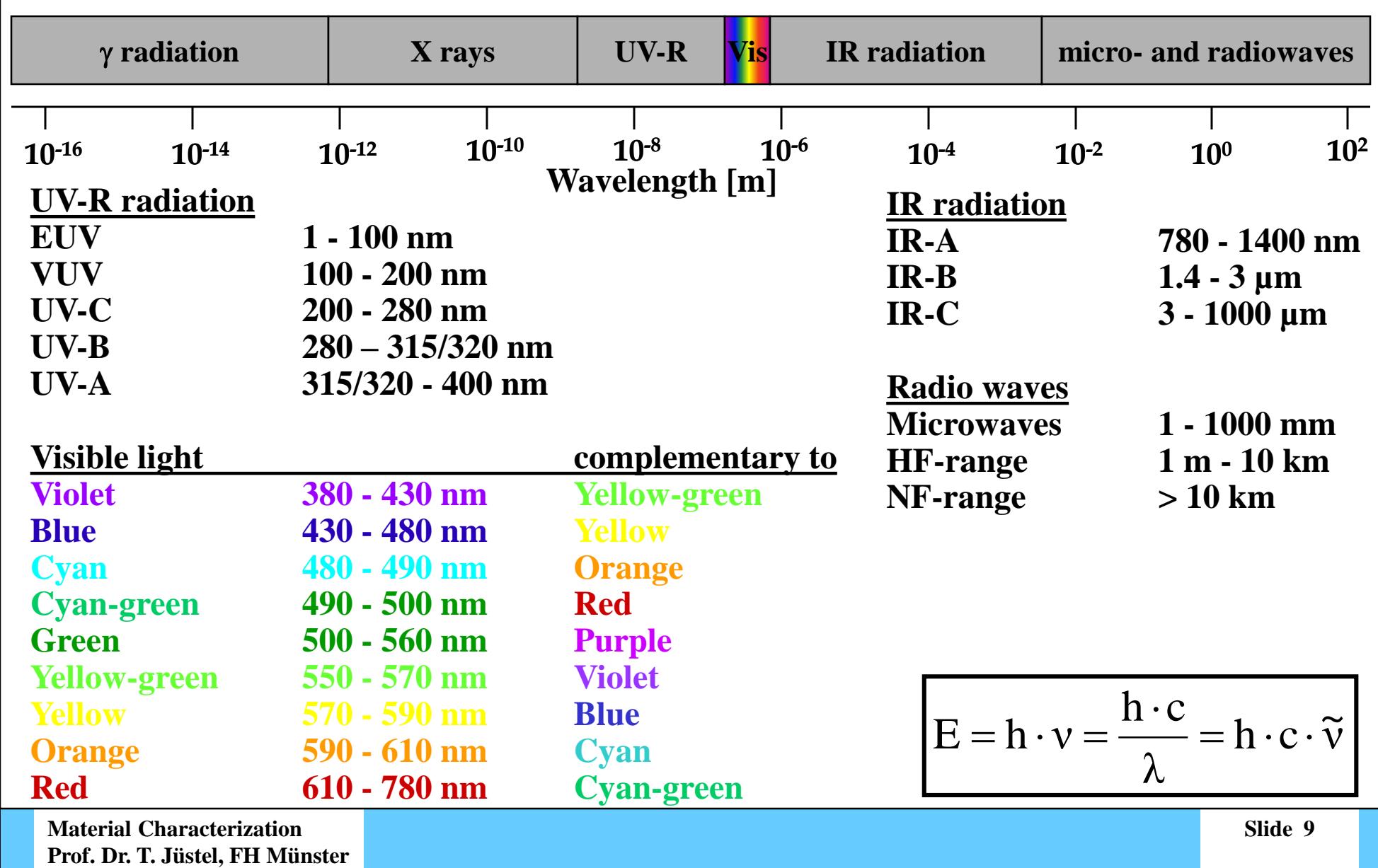
2.6 Photometric Quantities

2.7 Photosynthetically Active Radiation

2.1 Classification of Optical Spectroscopy



2.2 The Electromagnetic Spectrum



2.2 The Electromagnetic Spectrum

Energy of photons

$$E = h\nu$$

$$h = 6.626 \cdot 10^{-34} \text{ Js}$$

$$E = hc/\lambda$$

$$c = 2.9979 \cdot 10^8 \text{ m/s}$$

$$E = N_A hc/\lambda = 119226 / \lambda \text{ kJ mol}^{-1}$$

The energy is thus proportional to the reciprocal length, to so-called wavenumber $\tilde{\nu}$

$$\text{with } \tilde{\nu} = 1/\lambda$$

$$E = h \cdot c \cdot \tilde{\nu}$$

$$= 1.240 \cdot 10^{-4} \cdot \tilde{\nu} [\text{eV}]$$

$$= 1240/\lambda [\text{eV}]$$

$$\text{or } 1 \text{ eV} \sim 8065 \text{ cm}^{-1}$$

Parameters proportional to energy

1. Frequency ν [Hz]

2. Wavenumber $\tilde{\nu}$ [cm^{-1}]

Wavelength λ	Frequency ν	Wavenumber $\tilde{\nu}$	Energy E
10^{-16} m	$3 \cdot 10^{24} \text{ Hz}$	$1 \cdot 10^{14} \text{ cm}^{-1}$	12.4 GeV
10^{-14} m	$3 \cdot 10^{22} \text{ Hz}$	$1 \cdot 10^{12} \text{ cm}^{-1}$	124 MeV
$10^{-12} \text{ m (1 pm)}$	$3 \cdot 10^{20} \text{ Hz}$	$1 \cdot 10^{10} \text{ cm}^{-1}$	1.24 MeV
10^{-10} m	$3 \cdot 10^{18} \text{ Hz}$	$1 \cdot 10^8 \text{ cm}^{-1}$	12.4 keV
10^{-8} m	$3 \cdot 10^{16} \text{ Hz}$	$1 \cdot 10^6 \text{ cm}^{-1}$	124 eV
$10^{-6} \text{ m (1 } \mu\text{m)}$	300 THz	$1 \cdot 10^4 \text{ cm}^{-1}$	1.24 eV
10^{-4} m	3 THz	100 cm^{-1}	12.4 meV
10^{-2} m	30 GHz	1 cm^{-1}	0.124 meV
10^0 m (1 m)	300 MHz	0.01 cm^{-1}	1.24 μeV
10^2 m	3 MHz	$1 \cdot 10^{-4} \text{ cm}^{-1}$	12.4 neV

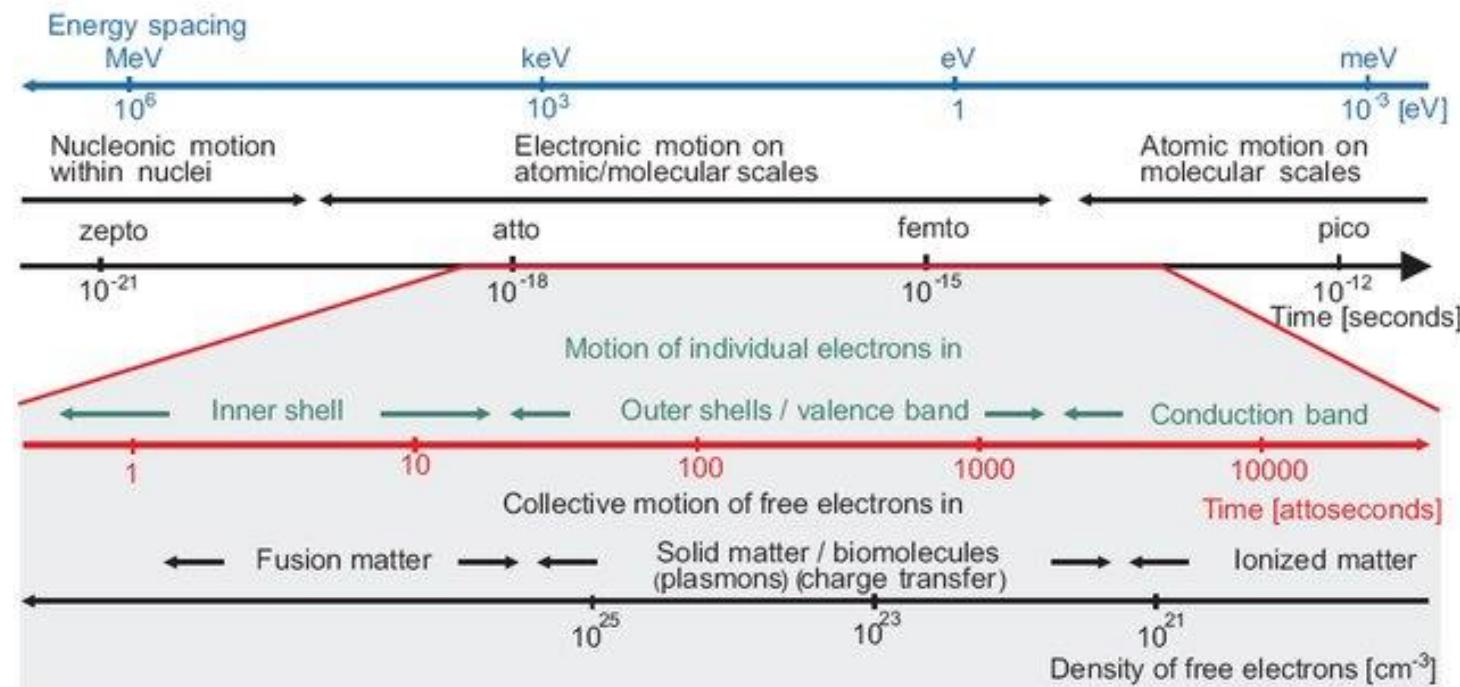
2.3 Time Scale of Physical Processes

Chemical reactions < molecular rotation < molecular vibration < electron motion

Franck-Condon Principle

Since the nuclei are much heavier than the electrons, the electronic transition takes place so rapidly that the nuclei have no time to react → Born-Oppenheimer approximation

→ Separation of electron and nucleus motion



Lit.: Ferenc Krausz and Misha Ivanov. Attosecond physics. Reviews of Modern Physics 81 (2019) 163

2.4 Interaction of Light and Matter

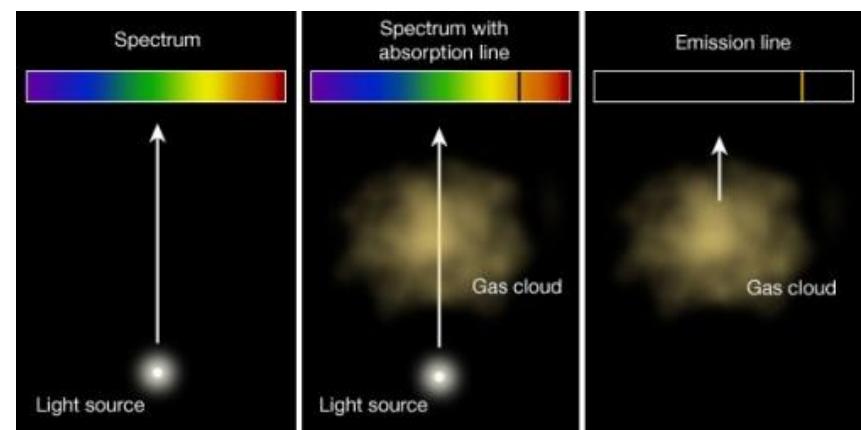
Types of electromagnetic interaction

Radiation type	Interaction with	Physical process
γ -radiation	Nuclei	Excitation of nuclear states
X-rays	Inner electrons	Excitation/diffraction
UV/Vis	Valence electrons	Excitation of electronic states
Infrared	Molecules	Vibrational excitation (phonons)
Microwaves	Molecules, electrons	Spin flip, rotational excitation
Radio waves	Nuclear spin	Spin flip

Macroscopic phenomena: absorption, refraction, emission, reflection (remission), transmission

Energy conservation: $A + R + T = 1$ or $E + R + T = 1$

Kirchhoff's law: $A(\lambda) = E(\lambda)$ (G.R. Kirchhoff 1859)



2.4 Interaction of Light and Matter

Electronic states in multi electron atoms

<u>Atom</u>	<u>Electron configuration</u>	<u>Ground state</u>	<u>First excited state</u>
Na ⁰	1s ² 2s ² 2p ⁶ 3s ¹	² S _{1/2}	² P _{3/2} , ² P _{1/2}
<i>(yellow Na-double line at 589.0 and 589.6 nm ⇒ Spectrometer resolution test)</i>			
Cr ³⁺	1s ² 2s ² 2p ⁶ 3s ² 3p ⁶ 4s ² 3d ³	⁴ F	⁴ P
Ce ³⁺	[Xe]4f ¹	² F _{5/2} , ² F _{7/2}	² D
Eu ³⁺	[Xe]4f ⁶	⁷ F ₀	⁵ D ₀
Gd ³⁺	[Xe]4f ⁷	⁸ S _{7/2}	⁶ P _J

$S = \sum s$
 $L = |\sum l|$
 $J = L+S, L+S-1, \dots |L-S|$
Russell-Saunders-Terms (energy levels)
 $2S+1L_J$

Selection rules for electric dipole transitions

1. Spin selection rule $\Delta S = 0$
2. Orbital angular momentum $\Delta l = +/-1, \Delta m_l = 0, +/-1$
3. Laporte selection rule $g \rightarrow u, u \rightarrow g, \text{ but not } g \rightarrow g, u \rightarrow u$

With Russell-Saunders (LS) coupling: $\Delta L, \Delta J = 0, +/-1, \text{ but not } J = 0 \rightarrow J = 0$

2.4 Interaction of Light and Matter

Electronic transitions in multi electron atoms

Type		ϵ [l mol ⁻¹ cm ⁻¹]	Example	Electron from	to
s-p transitions	$\Delta l = 1$	$10^3 - 10^4$	Na ⁰	3s	\rightarrow 3p
d-d transitions	$\Delta l = 0$	$< 10^2$	Cr ³⁺	3d (t_2)	\rightarrow 3d (e)
f-d transitions	$\Delta l = 1$	$10^4 - 10^5$	Ce ³⁺	4f	\rightarrow 5d
f-f transitions	$\Delta l = 0$	$< 10^2$	Eu ³⁺	4f	\rightarrow 4f
			Gd ³⁺	4f	\rightarrow 4f
			Tb ³⁺	4f	\rightarrow 4f

Forbidden due to large change in orbital angular momentum:

s-d $\Delta l = 2$

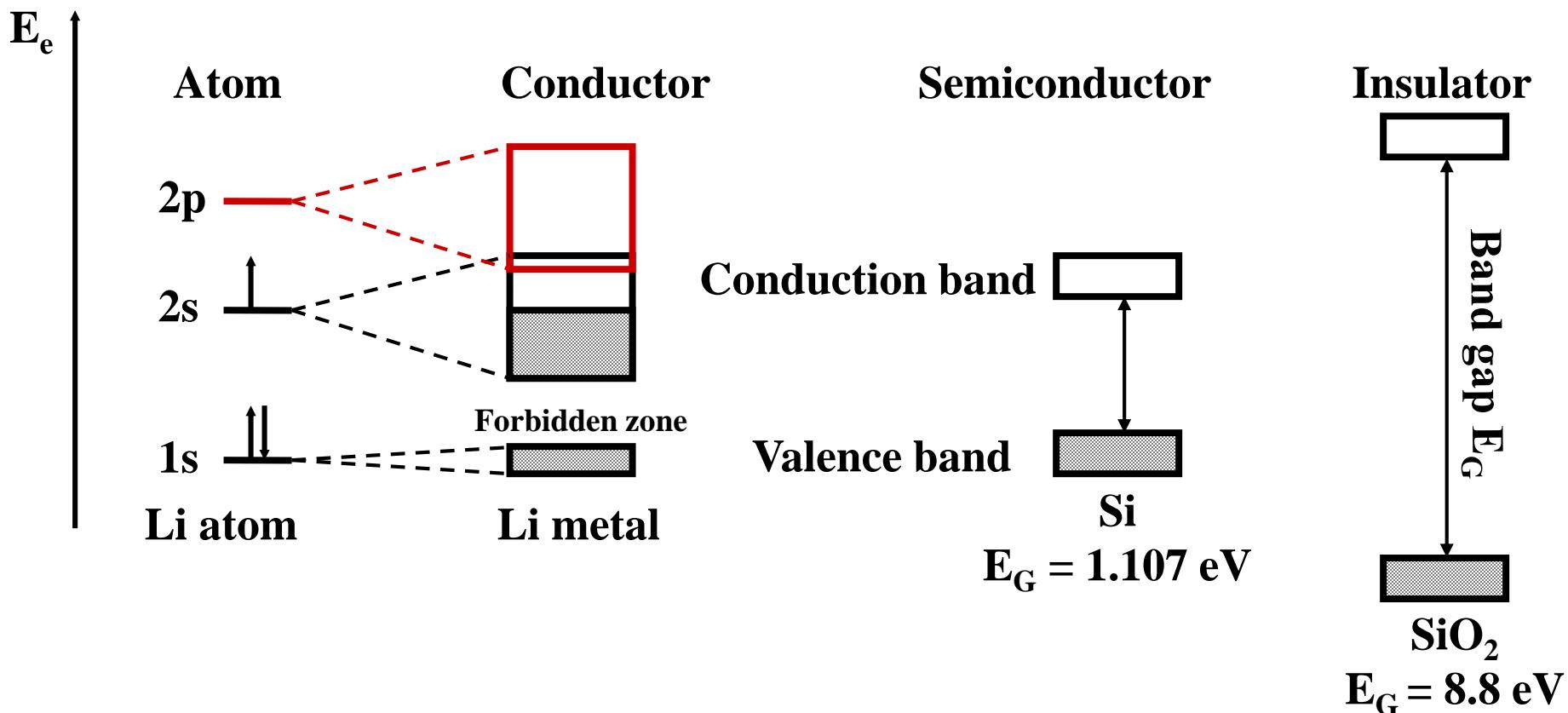
p-f $\Delta l = 2$

s-f $\Delta l = 3$ parity allowed

- these transitions are strongly forbidden!

2.4 Interaction of Light and Matter

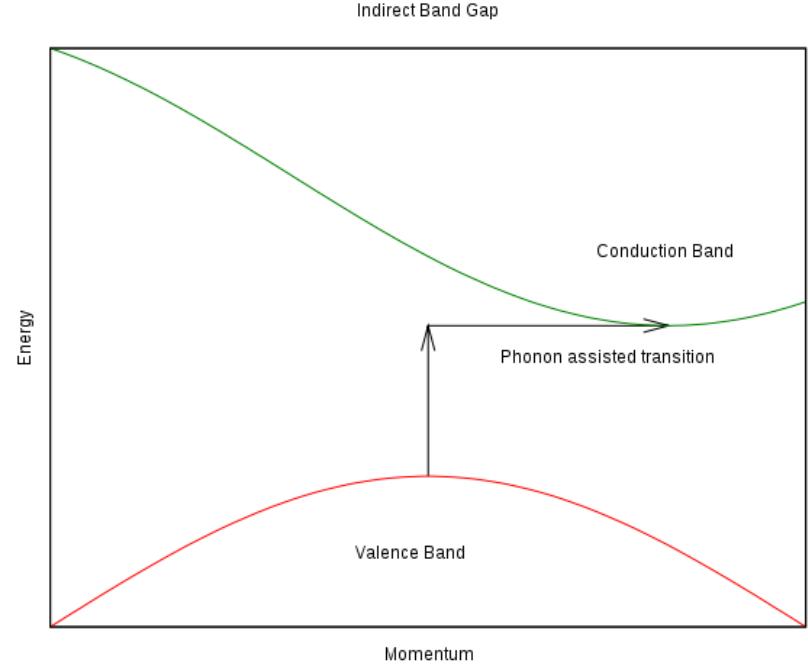
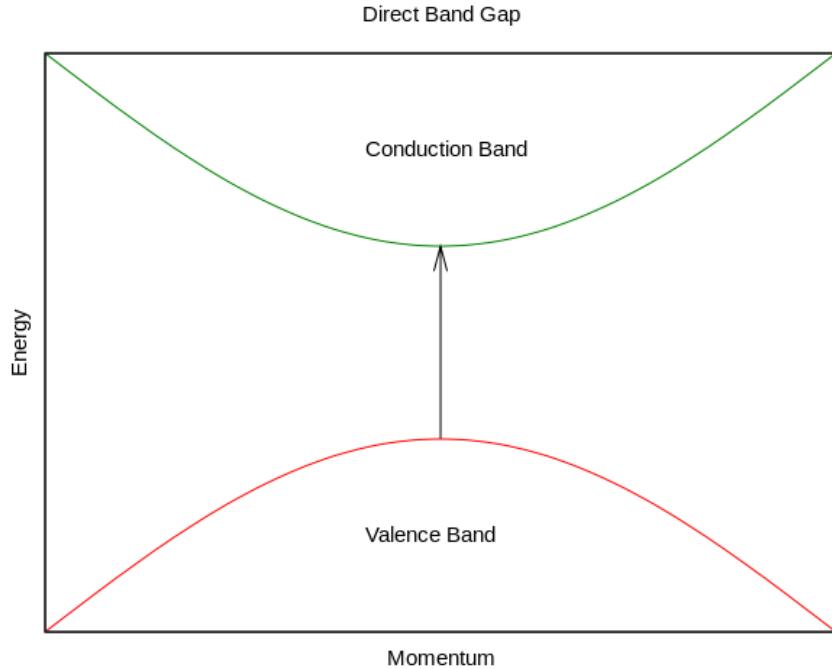
Electronic states in solids (metals, semiconductors, ceramics)



Semiconductors are materials with a small band gap (< 3.0 eV) whose conductivity are increased by structural defects, heating, or by irradiation

2.4 Interaction of Light and Matter

Electronic states in solids (metals, semiconductors, ceramics)



Direct semiconductors:

Absorption or emission of a photon

→ strong: GaAs, GaN

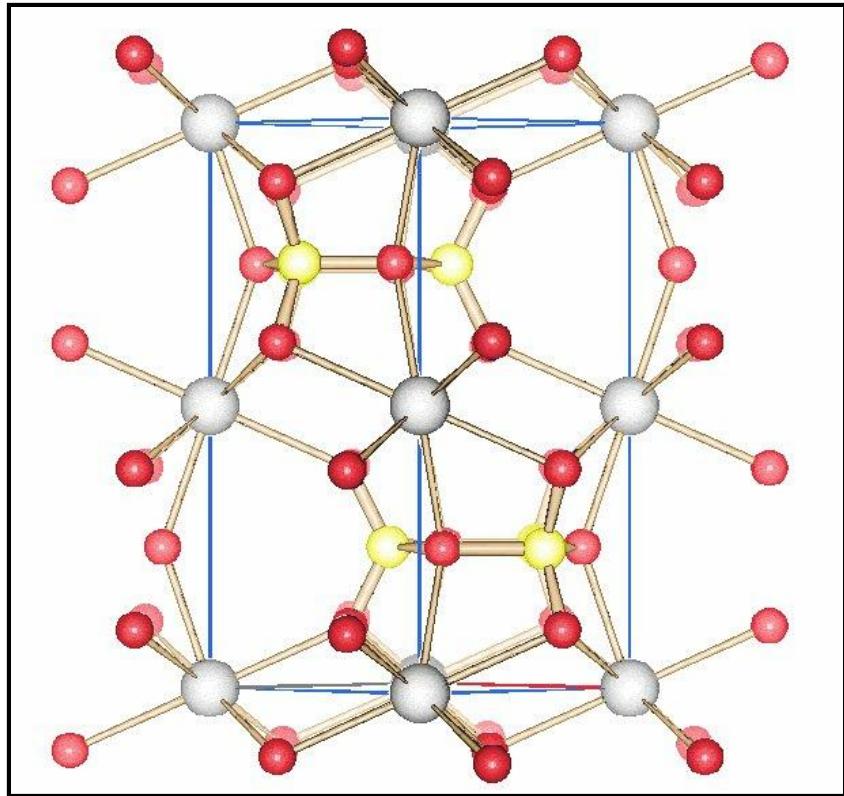
Indirect semiconductors:

Phonon assisted absorption or emission of a photon

→ weak: Si, Ge

2.4 Interaction of Light and Matter

Electronic states in solids



YBO_3 Vaterit

Band gap $E_g = 6.5 \text{ eV}$

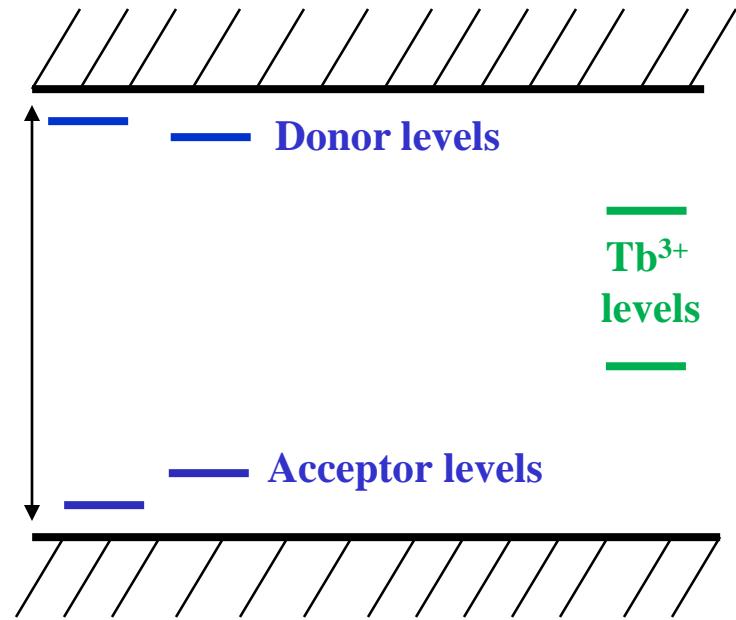
Conduction band

Band gap E_g

Valence band

Absorption by

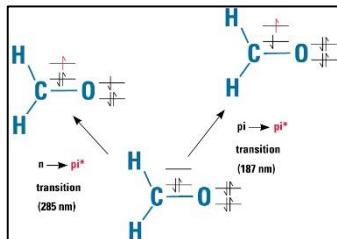
- Excitation band: valence band-conduction band
- Defects (color centers) “blue ice”
- Transitions between energy levels of dopants (Tb^{3+}) or undesired impurities



2.4 Interaction of Light and Matter

Electronic transitions in solids and molecules

Type	ϵ [$\text{lmol}^{-1}\text{cm}^{-1}$]	Example	Transition type
Band gap transitions	$10^5 - 10^6$	TiO_2	$\text{O}^{2-} \rightarrow \text{Ti}^{4+}$ (VB \rightarrow CB)
CT transitions	$10^4 - 10^6$	WO_4^{2-}	$\text{O}^{2-} \rightarrow \text{W}^{6+}$ „LMCT“
		$\text{Fe}(\text{CO})_5$	$\text{Fe}^0 \rightarrow \text{CO}$ „MLCT“
		$\text{KFe}[\text{Fe}(\text{CN})_6]$	$\text{Fe}^{2+} \rightarrow \text{Fe}^{3+}$ „MMCT“
		$[(\text{cp})_2\text{Zr}(\text{biq})]^{2+}$	$\text{cp}^- \rightarrow \text{biq}$ “LLCT”
$\sigma \rightarrow \sigma^*$	$> 10^3$	CH_4	$\text{HOMO} \rightarrow \text{LUMO}$
$\pi \rightarrow \pi^*$	$10^4 - 10^6$	Butadiene	$\text{HOMO} \rightarrow \text{LUMO}$
$n \rightarrow \pi^*$	$10^1 - 10^2$	$(\text{CH}_3)_2\text{C=O}$	$\text{HOMO} \rightarrow \text{LUMO}$



In addition

- Excitation of vibrational states, i.e. phonons in the solid state or valence and deformation vibrations in molecules
- Excitation of rotational states (only molecules)

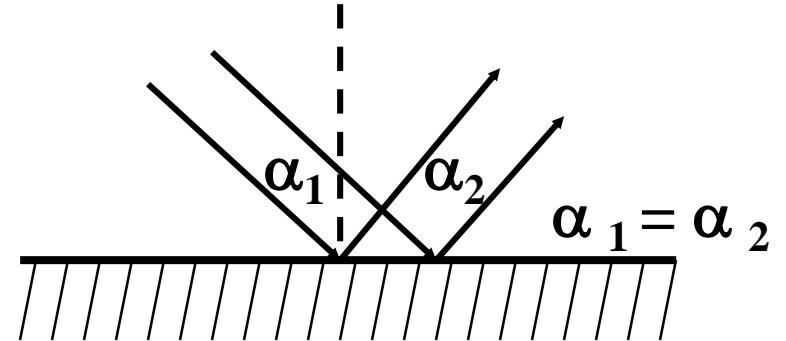
2.4 Interaction of Light and Matter

Reflection

Specular reflection (regular)

Reflection from a polished (smooth) surface
in one direction (reflecting surface)

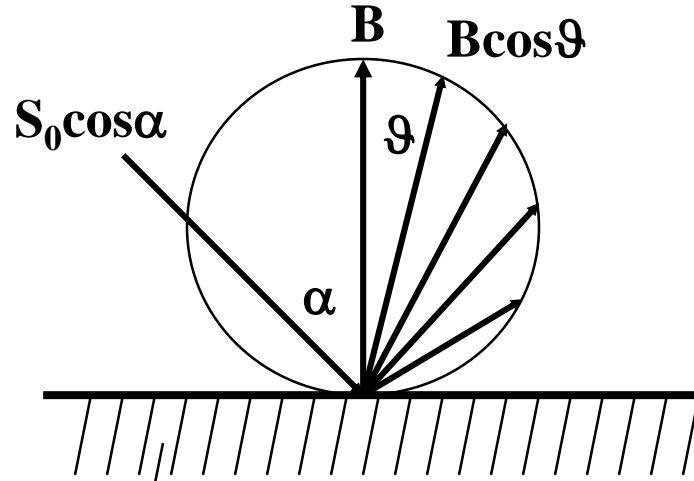
Example: metal, mirror, calm water surface



Diffuse reflection (irregular)

Reflection from a non-polished surface in all
directions of the half-space (matte surface)

Example: paper, powder, white wall

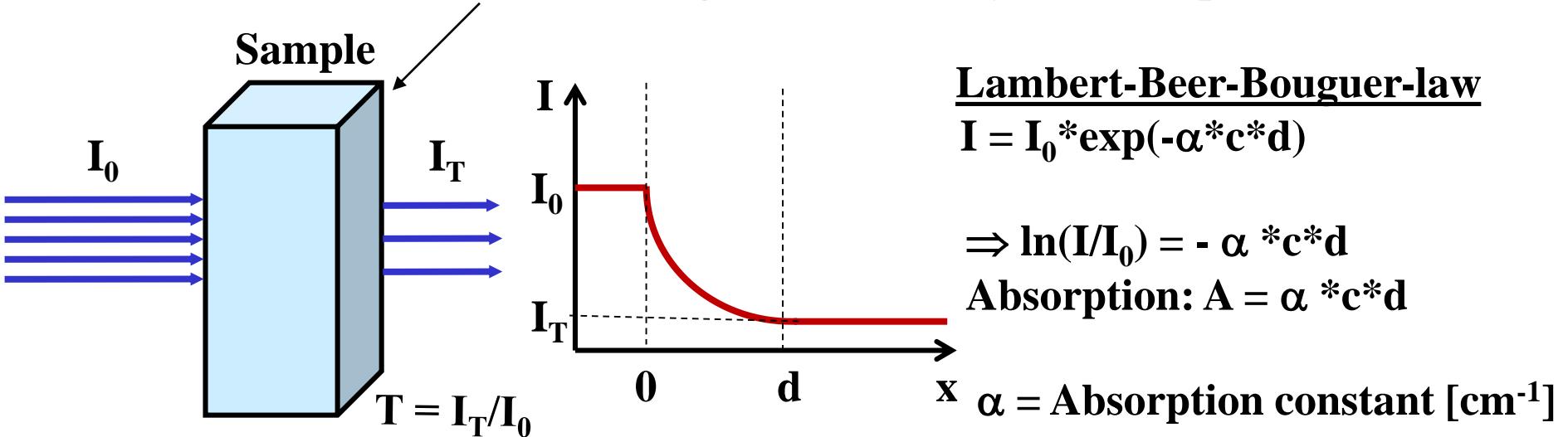


Angular distribution of a
Lambert radiator's

2.4 Interaction of Light and Matter

Absorption

Medium = gas, solution, crystal, transparent ceramics



Fate of the absorbed energy

- Vibrations (phonons)
- Charge storage
- Luminescence
- Charge separation

Physical effect

- Heat
- Color centers
- Light emission
- Current

Application areas

- bolometer
- detectors, imaging plate
- phosphors, laser gain media
- solar cells

2.4 Interaction of Light and Matter

Quantitative description of absorption

Energy conservation
(division by I_0 yield)

$$I_A + I_R + I_T = I_0$$
$$A + R + T = 1 \text{ (100%)}$$

Transmission

$$T = I_T/I_0 \leq 1 \text{ (100%)}$$

Absorption

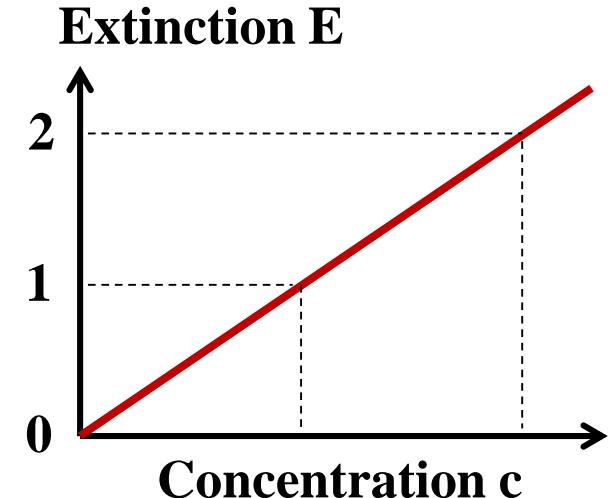
$$A = 1 - T = 1 - I_T/I_0 \leq 1 \text{ (100%)} \\ (\text{if } R = 0)$$

Extinction

$$E = -\lg(I_T/I_0) = \lg(I_0/I_T)$$
$$E = \epsilon \cdot c \cdot d$$

ϵ = molar extinction coefficient (material constant, dependent on λ, T, p)

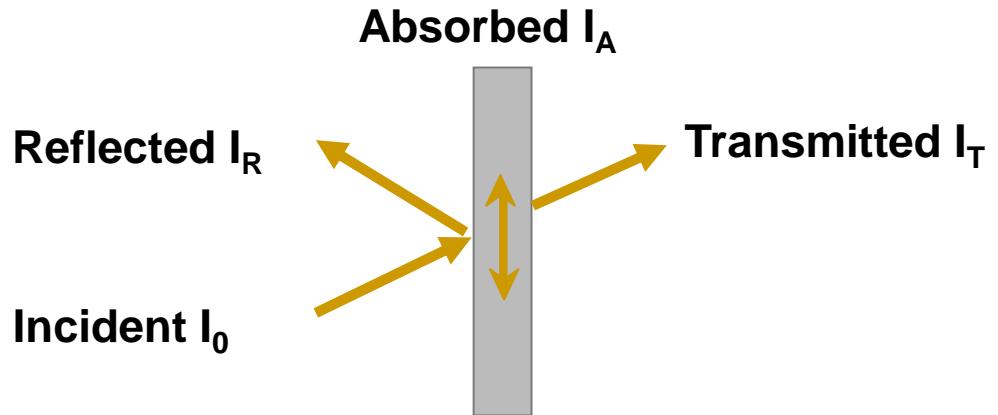
Interstellar extinction:
 $2 \text{ mag / kpc} = 0.83 / 3260 \text{ ly}$
since $1 \text{ pc} = 3.26 \text{ ly}$
(due to Rayleigh scattering at dust particles)



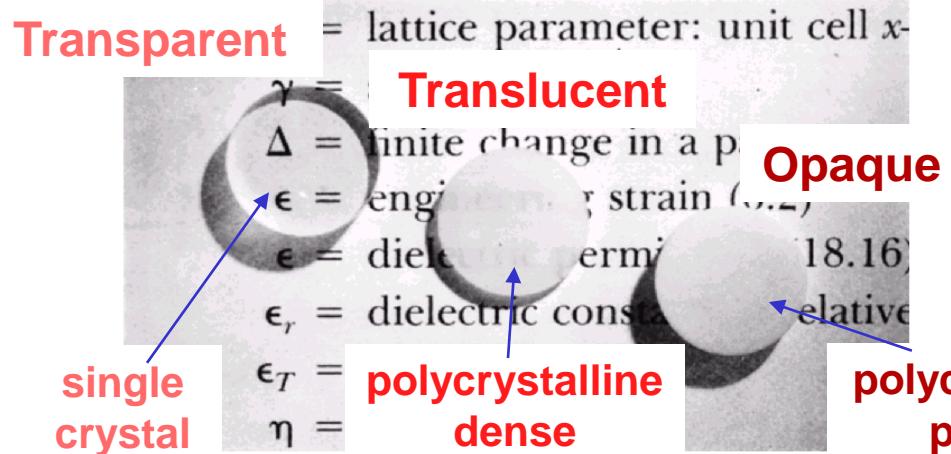
2.4 Interaction of Light and Matter

Transmission

$$I_T = I_0 \cdot I_A - I_R$$



Classification of Materials

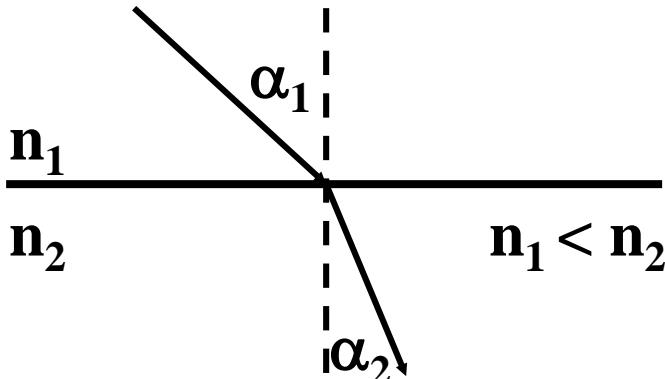


Adapted from Fig. 21.10,
Callister 6e. (Fig. 21.10 is by
J. Telford, with specimen
preparation by P.A. Lessing)

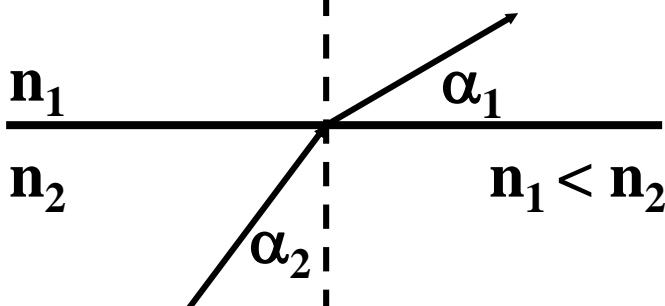
2.4 Interaction of Light and Matter

Refraction

Transition from an optically thin
into an optically denser medium



Transition from an optically dense
into an optically thinner medium



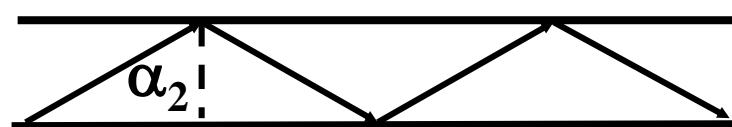
$$n_1 \cdot \sin \alpha_1 = n_2 \cdot \sin \alpha_2 \quad (\text{Snellius equation})$$

Total reflection occurs at the transition
from an optically dense to an optically
thinner medium, if the critical angle α_T is
exceeded

Critical angle $\alpha_2 = \arcsin(n_1/n_2)$

Vacuum/air $\alpha_T = \arcsin(1/n)$

Principle of light guiding in glass fibers
(light guides)



2.4 Interaction of Light and Matter

Refractive indices of some substances at 589.3 nm (center of Na double line)

Substance	Refractive index n
Vacuum	1.000
Air	1.0003
Water	1.333
Ethanol	1.360
CCl ₄	1.460
Benzene	1.501
CS ₂	1.628
Ice	1.309
Teflon	1.35
CaF ₂	1.434
Quartz glass	1.459
PMMA	1.49
NaCl, α -quartz	1.54
ZrSiO ₄ (Zircon)	1.923
Diamond	2.417
TiO ₂ (Rutil)	2.68

Critical angle for total reflection

Vacuum/air

$$\begin{aligned}\alpha_T &= \arcsin(1/n_{\text{air}}) \\ &= 88.6^\circ\end{aligned}$$

(air glimmering)

Quartz glass/air

$$\begin{aligned}\alpha_T &= \arcsin(n_{\text{air}}/n_{\text{glass}}) \\ &= 43.3^\circ\end{aligned}$$

(silica fibre)

Diamond/air

$$\begin{aligned}\alpha_T &= \arcsin(n_{\text{air}}/n_{\text{diamond}}) \\ &= 24.5^\circ\end{aligned}$$

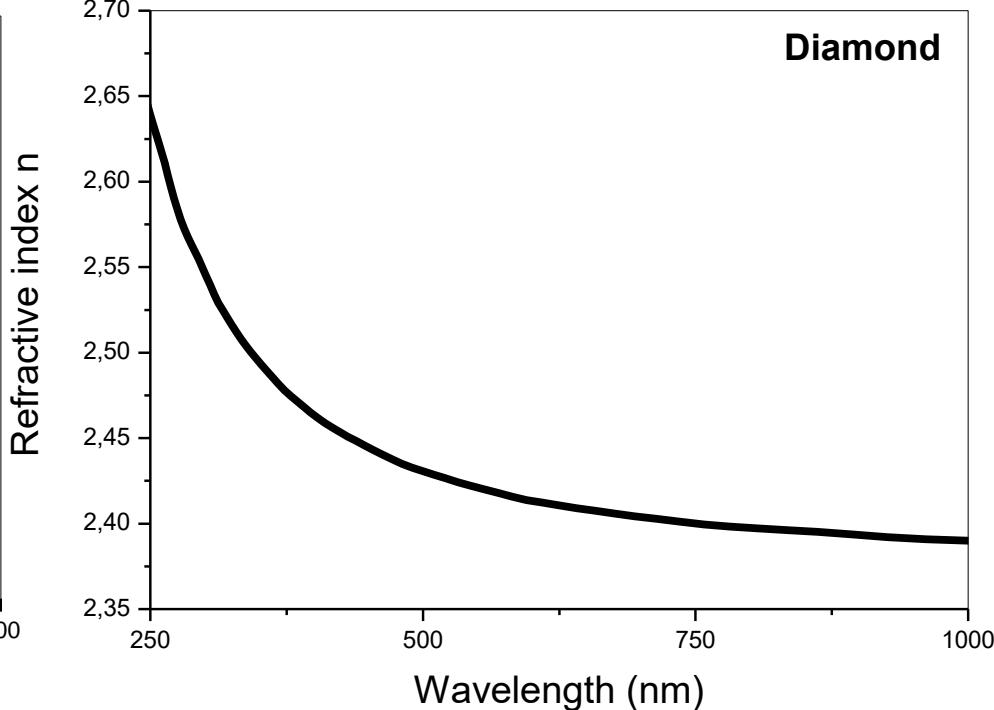
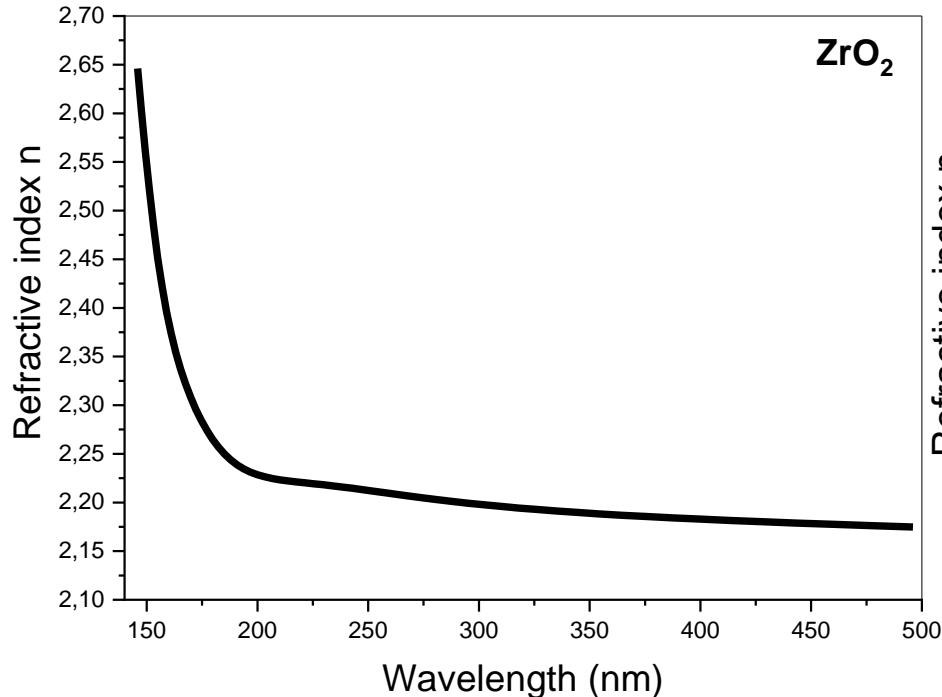
("Fire" of diamonds, since n
strongly depends on wavelength)

2.4 Interaction of Light and Matter

Refractive indices as a function of wavelength (dispersion)

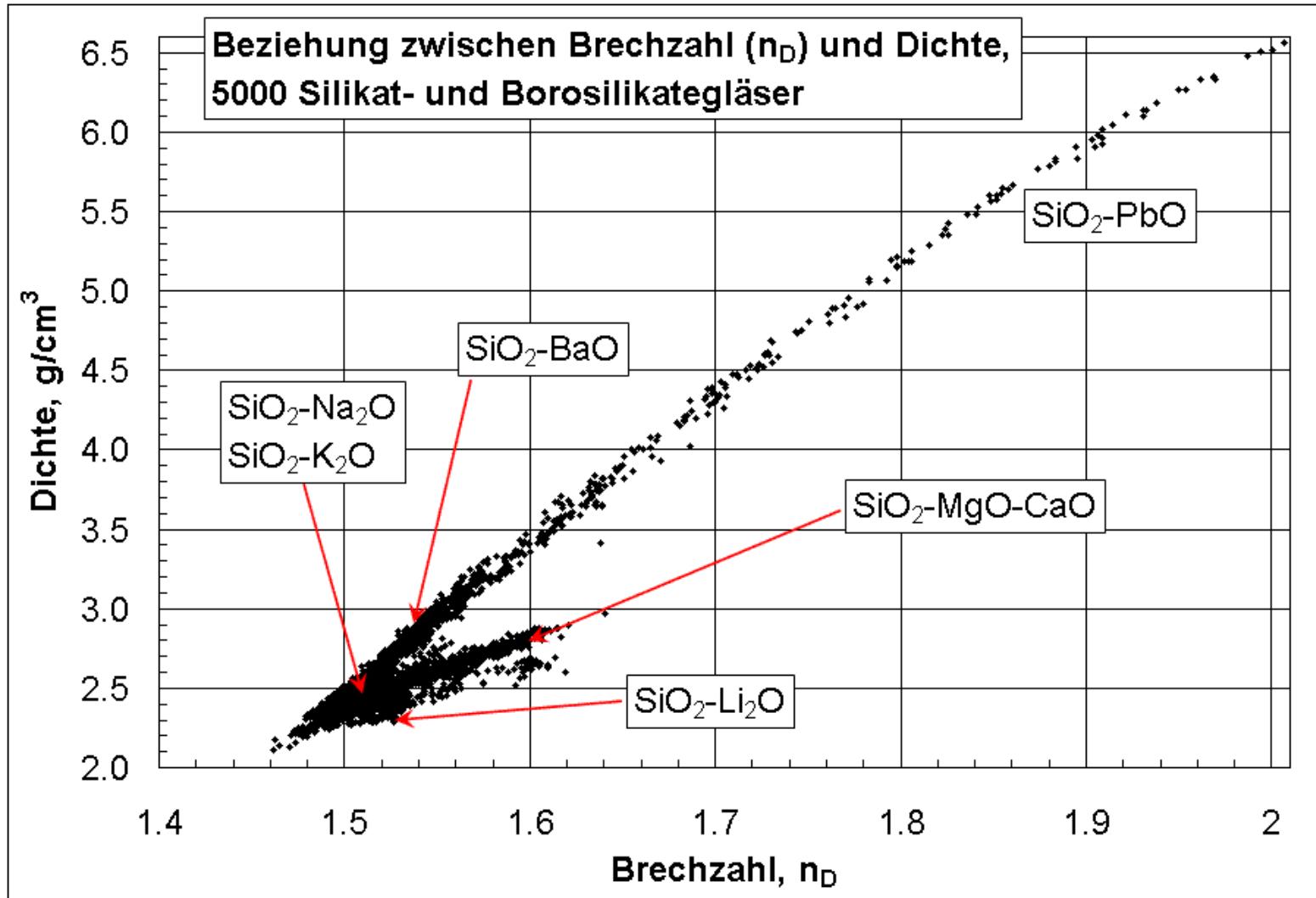
→ Causes chromatic aberration

weak dispersion
in the visible range
strong dispersion



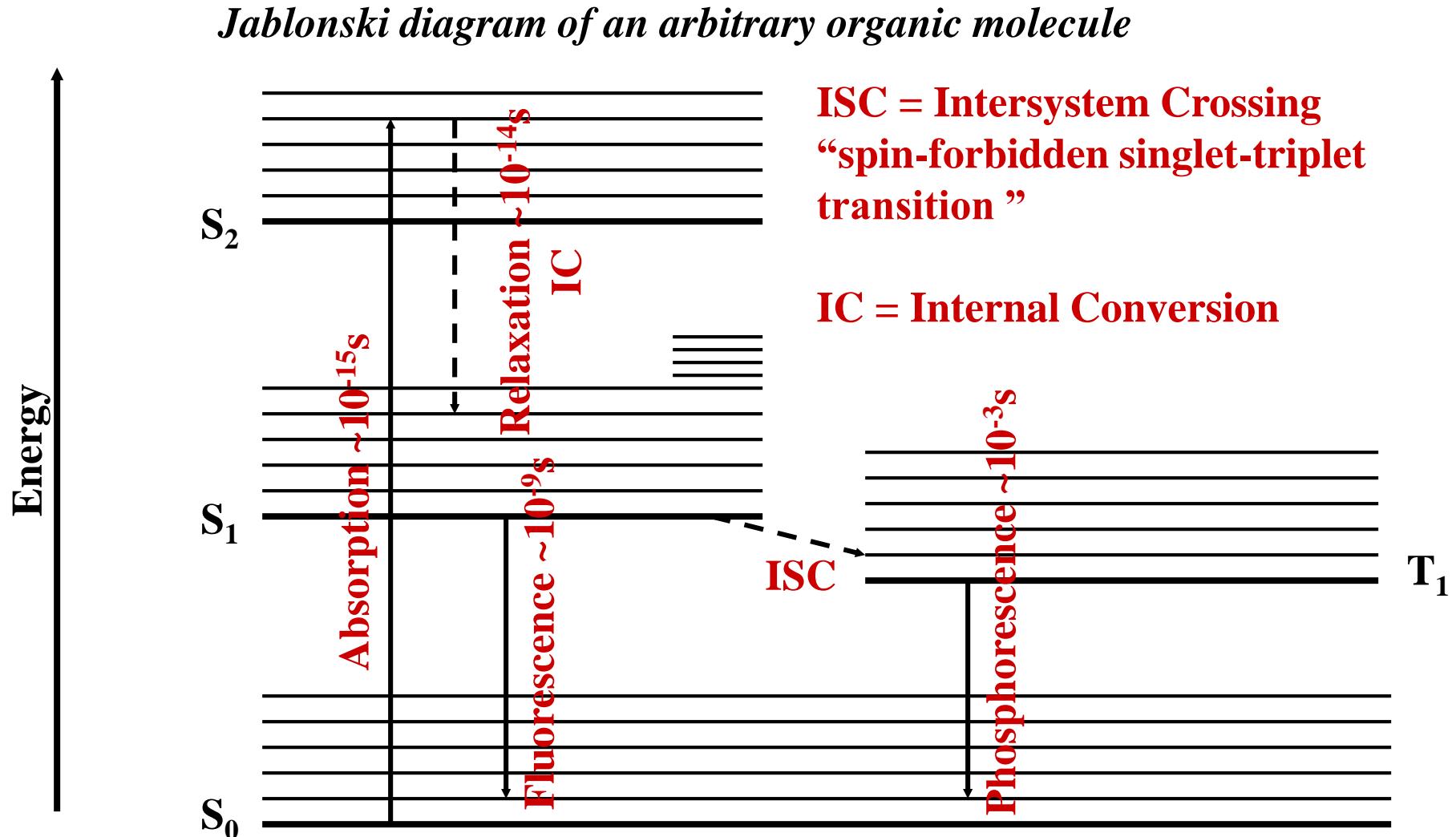
2.4 Interaction of Light and Matter

Refractive indices as a function of the density of glasses



2.4 Interaction of Light and Matter

Luminescence



ISC = Intersystem Crossing
“spin-forbidden singlet-triplet transition ”

IC = Internal Conversion

2.5 Radiometric Quantities

Measured variables for describing the performance (proportional to the number of photons per unit of time) = energetic parameters

Measured variables at the detector (photomultiplier tube, photodiode, eye)

(Radiant) intensity I = number of photons /area*time [N_{hv}/m^2s]

Irradiance E_e = number of photons /area*time [$J/m^2s = W/m^2$]

⇒ These parameter are proportional to the detector count rate [Counts/s]

⇒ Technically, intensity is measured with a PMT or a photodiode

(Integration over a given area, i.e. area of the window of the detector)

Radiant exposure = number of photons /area [J/m^2]

2.5 Radiometric Quantities

Physical variable	Symbol	Definition	Unit
Radiant power/flux	Φ_e	$= dW/dt$	[W] bzw. [J/s]
Spectral radiant power/flux	$\Phi_e(\lambda)$	$= d\Phi_e/d \lambda$	[W/nm]
Radiant emittance, irradiance	E_e	$= d\Phi_e/dA$	[W/m ²]
Spectral irradiance	L_e	$= dD_e/d\lambda$	[W/m ² nm]
Irradiance of the Earth: $E_e = 1.376 \cdot 10^3 \text{ J/m}^2\text{s} = 1.376 \text{ kW/m}^2$ (= solar constant)			
Number of photons:	$E = h\nu = hc/\lambda$ and $h\nu_{550} = 4 \cdot 10^{-19} \text{ J}$ $\Rightarrow 1 \text{ W} = 1 \text{ J/s} = 2.5 \cdot 10^{18} \text{ photons per second}$ with wavelength 550 nm $\Rightarrow 3.375 \cdot 10^{21} \text{ Photons/m}^2$ for 1.376 kW/m^2		

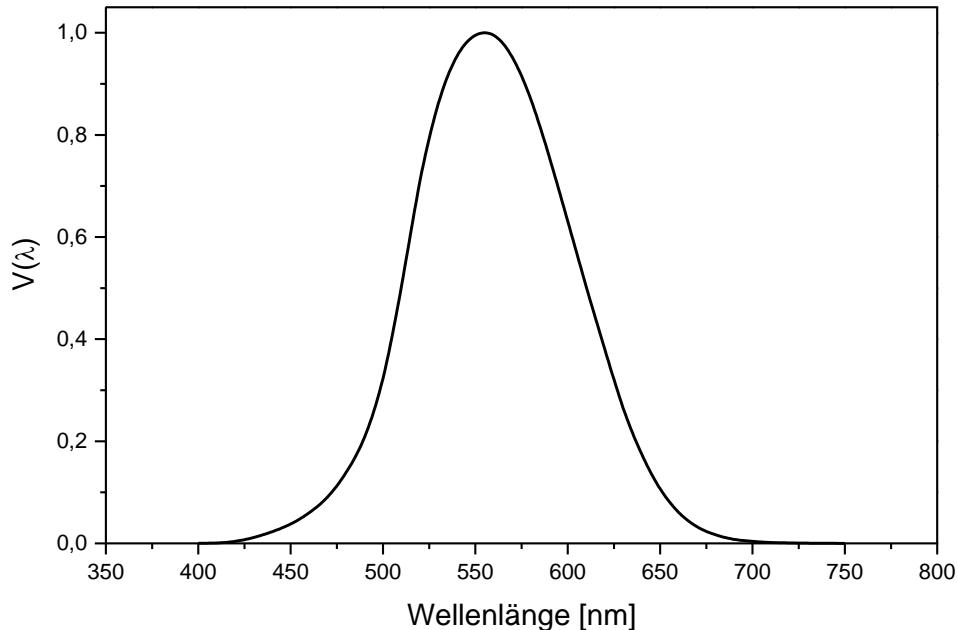
2.6 Photometric Quantities

Measured variables involving the spectral sensitivity of the observer (human eye)

Luminous flux $\Phi_v = \Phi_e/M_0$ [lm]

M_0 = Energetic light equivalent

$$= 0.00146 \text{ W/lm (1/683 W/sr)}$$



$$K_{\max} = 683 \text{ lm/W (bei } 555 \text{ nm)}$$

$$K(\lambda) = K_{\max} V(\lambda)$$

$$\Phi_v = K_{\max} \int_{380}^{780} V(\lambda) \Phi_e(\lambda) d\lambda$$

Luminous intensity $I = d\Phi_v/d\Omega$ [lm/sr = cd]

Ω = Solid angle [sr]

(surface of a sphere = $4\pi r^2 = 4\pi \text{ sr}$)

A 1 cd (1 lm / sr) light source emits 4π lm, i.e. 12.57 lm, isotropically in all spatial directions \Rightarrow Luminance spectrometer

2.6 Photometric Quantities

Origin of the quantity of the energetic light equivalent $M_0 = 0.00146 \text{ W/lm}$

Reference light source is the candle

Emission spectrum of a dinner light (~ 16 g)

$P \sim 40 \text{ W} = 40 \text{ J/s}$ (fuel: paraffine ~ 42 MJ/kg)

$P_{\text{optical}} = 0.0184 \text{ W}$ (integrating sphere)

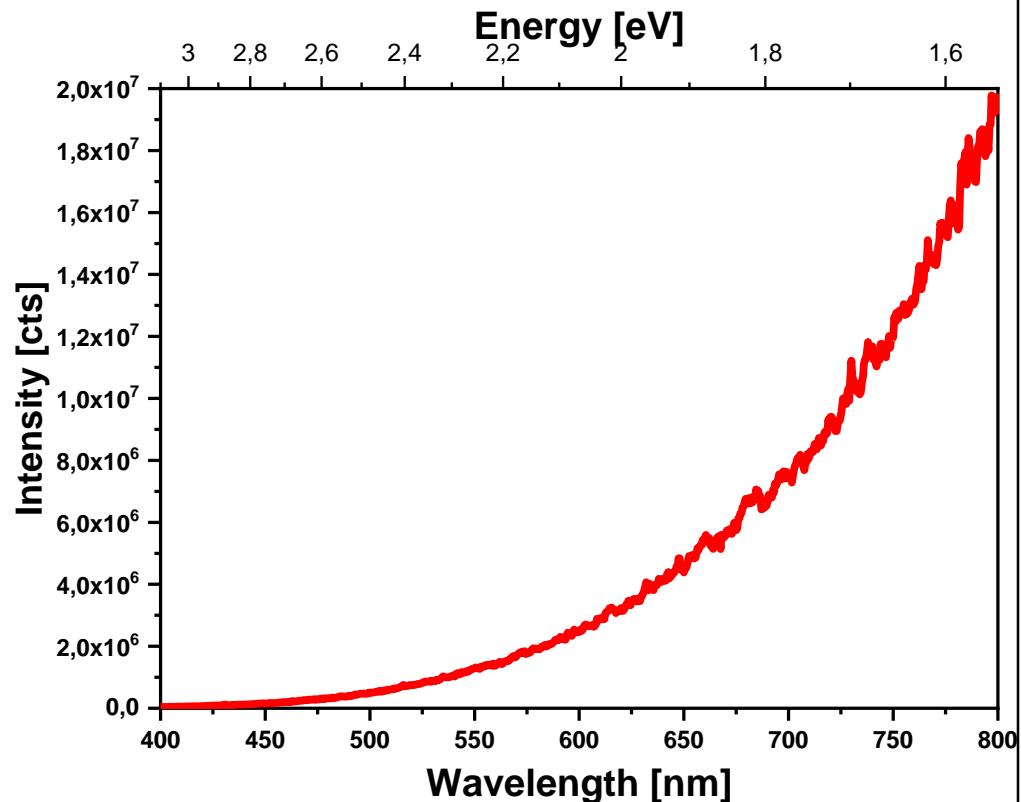
$\eta = 0.00046$ or 0.046%

$I_v = 1 \text{ cd} = 4\pi \text{ lm} \sim 12.566 \text{ lm}$

$\Rightarrow \eta_v = 0.31415 \text{ lm/W} \Rightarrow 3.1832 \text{ W/lm}$

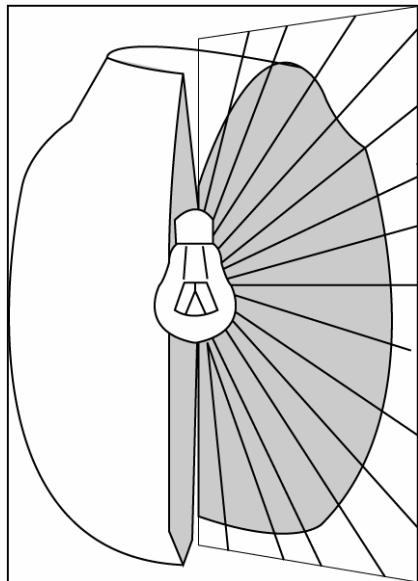
Efficiency $\eta = 0.00046$

$\Rightarrow 0.001464 \text{ W/lm}$ or 683 lm/W (q.e.d.)



2.6 Photometric Quantities

- Luminous flux** Total radiation power emitted from a light source in all spatial directions, which is measured via the sensitivity of the human eye [lm]. It is determined by measurement with an integrating sphere.
- Luminous int.** Luminous flux emitted per unit solid angle with the value of one steradian [cd]. For this purpose, the luminous flux detected only in a certain solid angle, typically in the range of the maximum beam intensity.



$$\text{Luminous flux: } \Phi_v = 683 \int_{380}^{780} V_{rel}(\lambda) \Phi_e(\lambda) d\lambda$$

2.6 Photometric Quantities

Illuminance

Ratio of luminous flux to the illuminated surface

$$[\text{lux} = \text{lm}/\text{m}^2]$$

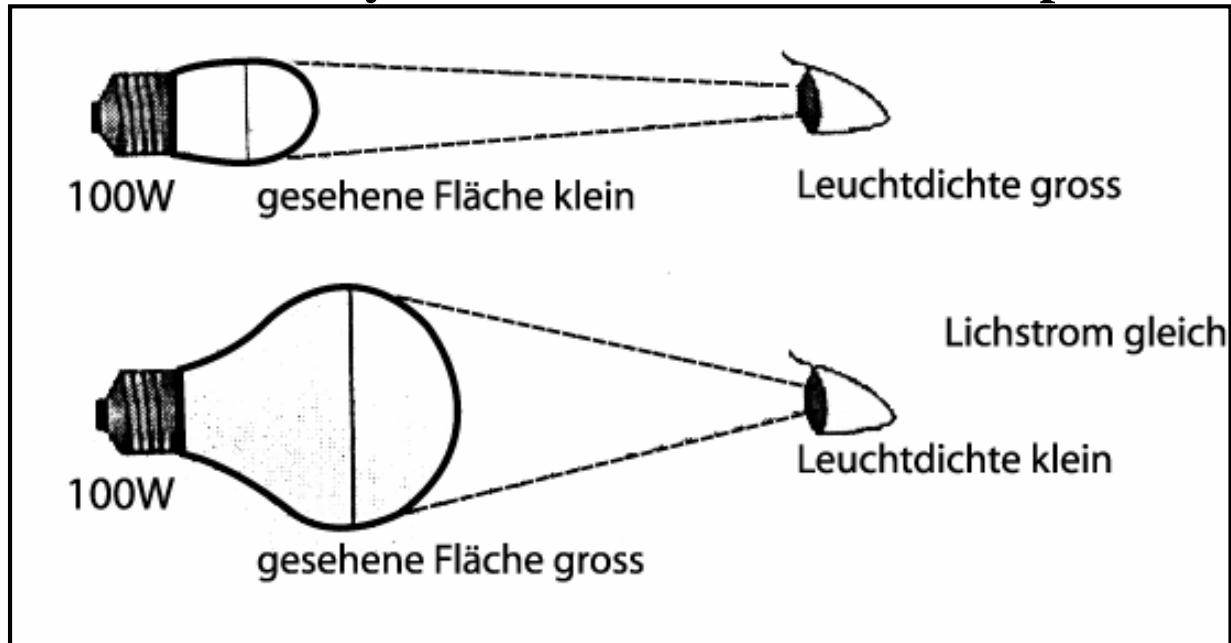
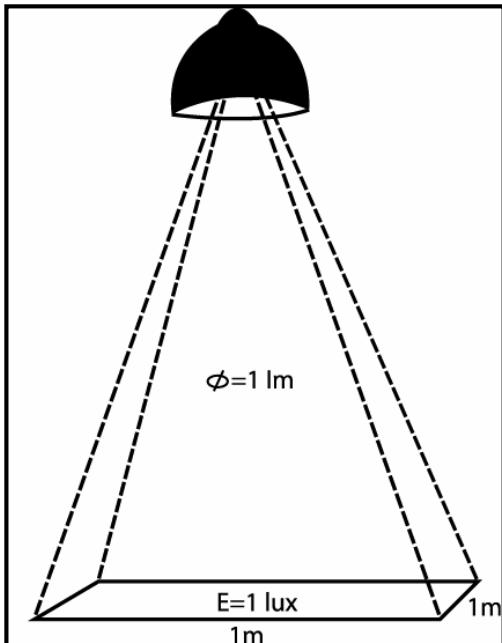
It is determined by measurement with a lux meter.

Luminance

Perceived brightness of a light source $[\text{cd}/\text{m}^2]$.

It is measured with a CCD camera.

Vision limit human eye $\sim 0.000003 \text{ cd}/\text{m}^2 \sim 250000 \text{ photons}$



2.6 Photometric Quantities

Integral Quantities

Luminous flux $\Phi_v = \Phi_e/M_0$ [lm]

Solid angle related parameters (for 1 sr)

Luminous intensity $I = d\Phi_v/d\Omega$ [cd]

Intensity of illumination $E = d\Phi_v/dA$ [lux = lm/m²] Light density $L = dI/dA \cos\gamma$
[cd/m²] = [nit] (luminance)

Light source	Light density / Luminance [cd/cm ²]
Sun	150000
Arc discharge	20000 - 100000
Bulb (clear)	200 – 2000
Bulb (matte)	5 – 50
Fluorescent lamp	0.4 – 1.4
Candle	0.75
Blue sky	0.3 – 0.5
Full moon	0.25
TV	0.05

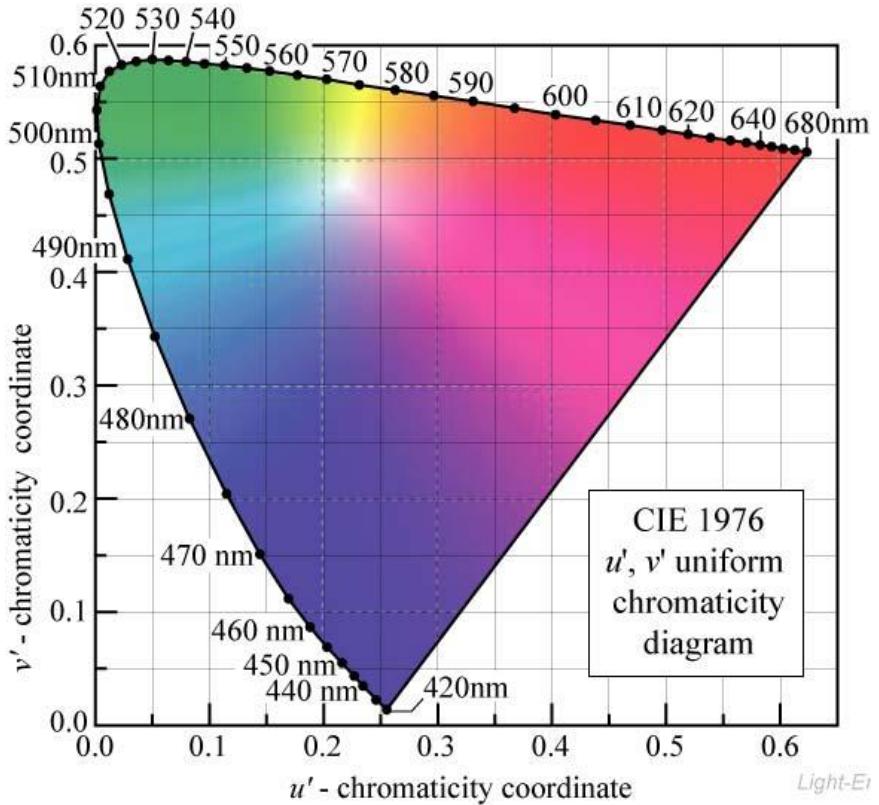
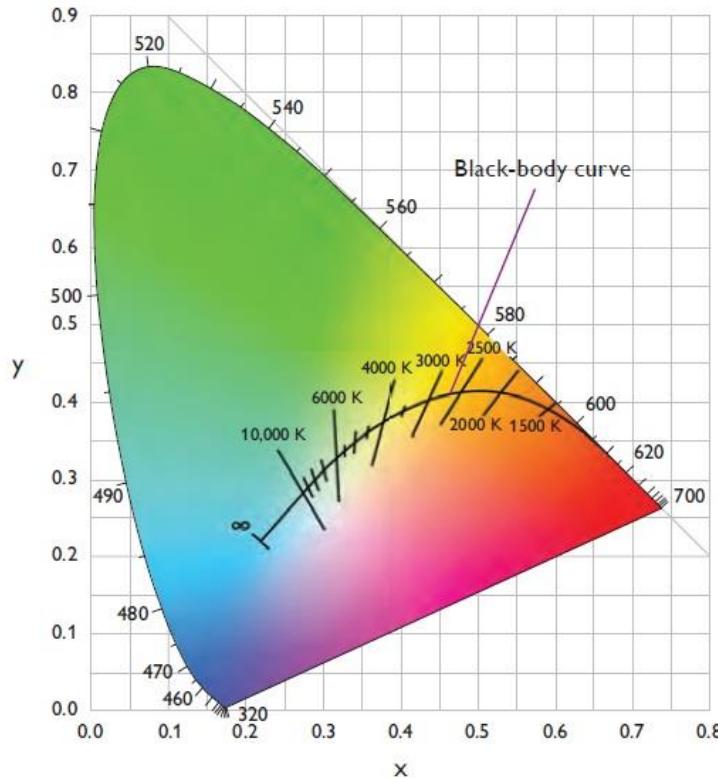
2.6 Photometric Quantities

Color Coordinates

Quantifies the perceived color of an emission spectrum
(additive color mixing)

Various systems exist to represent the color:

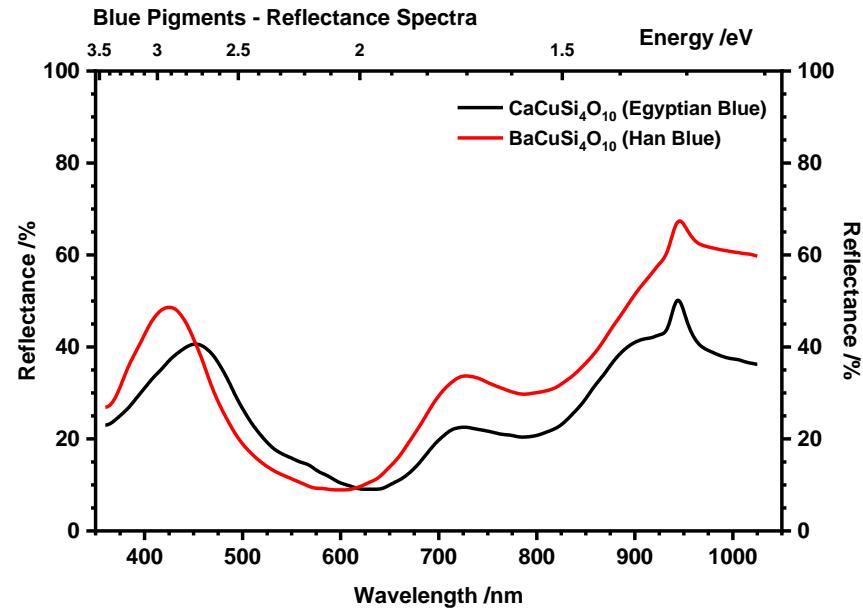
CIE1931, CIE1976, CIELab (includes luminance), ...



2.6 Photometric Quantities

Body Colour

- The perceived color of a subject. A result of subtractive color mixing
- Depends both on the absorption spectrum of the substance, as well as the emission spectrum of the light source
- A standard measurement is D65, where the light source mimicks solar light
- With very efficient phosphors the body color can be influenced by their luminescence. Absorbed UV radiation is converted into visible light



2.7 Photosynthetically Active Radiation

Photosynthetically Active Radiation (PAR)

Most autotrophic organism use photosynthesis for energy production and thus utilize photons of a specific wavelength range, mainly visible light

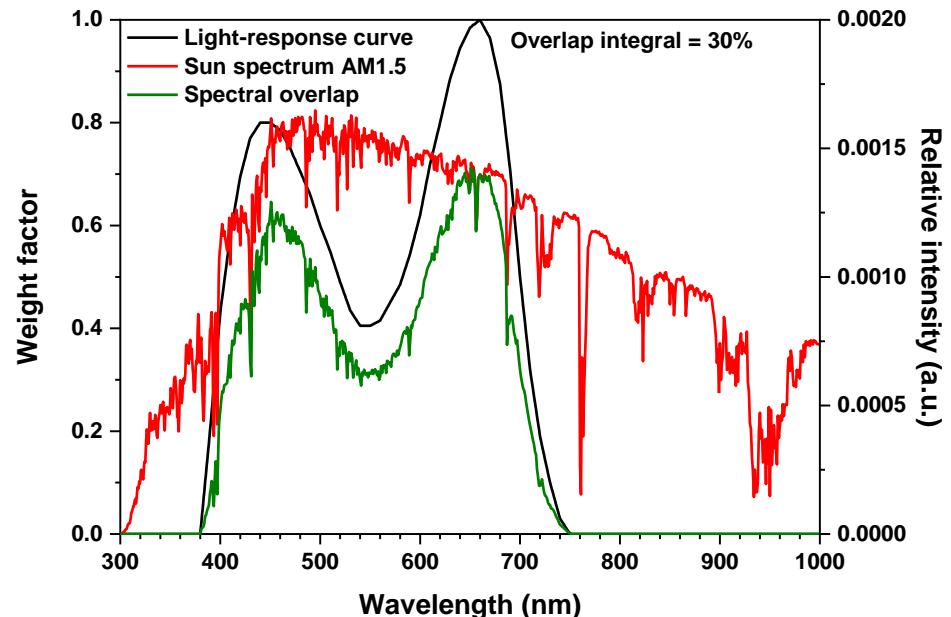
Photosynthetic photon flux (PPF)

- Number of 400-700 nm photons per time unit (treats all photons equally)
- Unit: $[\mu\text{mol} \cdot \text{s}^{-1}]$

Photosynthetic photon flux density (PPFD)

- Number of 400-700 nm photons per irradiated area and time
- Unit: $[\mu\text{mol} \cdot \text{s}^{-1} \cdot \text{m}^{-2} = \text{W} \cdot \text{m}^{-2}]$

→ Does not take the efficacy into account!



2.7 Photosynthetically Active Radiation

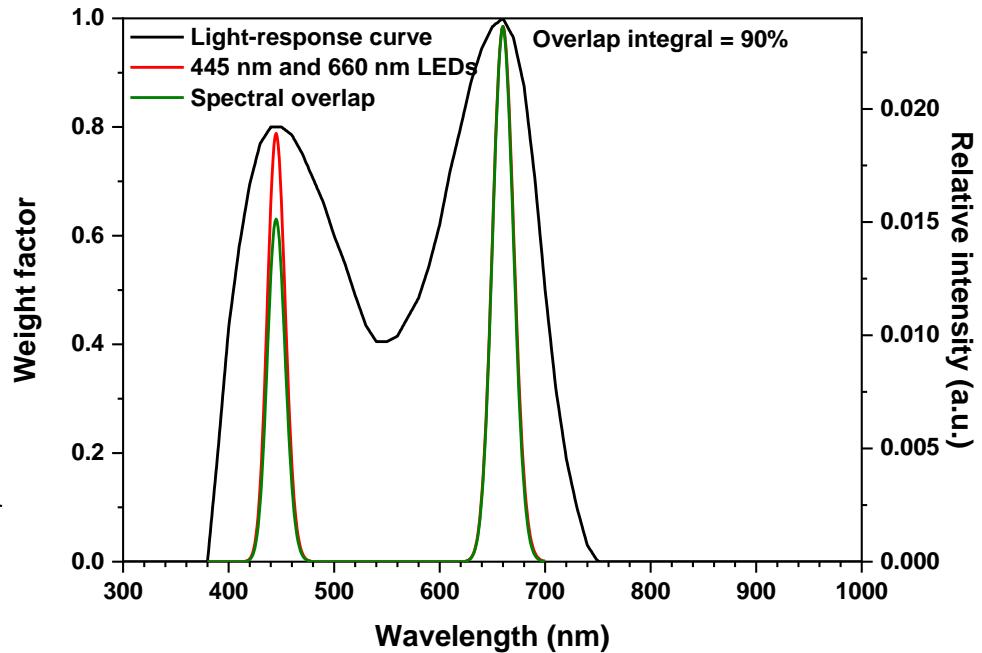
Photosynthetically Active Radiation (PAR)

Yield Photon Flux (YPF)

- Since the photosynthesis efficiency depends on the respective wavelength of the photon, a weight function is used
- Number of photons from 360-760 nm per time unit weighted by photosynthetic response
- Unit: $[\mu\text{mol} \cdot \text{s}^{-1}]$
- The spectral overlap integral (OI) quantifies the overall efficiency, thus $\text{YPF} = \text{OI} \cdot \text{PPF}$

Yield Photon Flux Density (YPFD)

- Number of photons from 360-760 nm per time unit and area weighted based on the photosynthetic response
- Unit: $[\mu\text{mol} \cdot \text{s}^{-1} \cdot \text{m}^{-2} = \text{W} \cdot \text{m}^{-2}]$



3. Experimental Setup

3. Experimental Setup

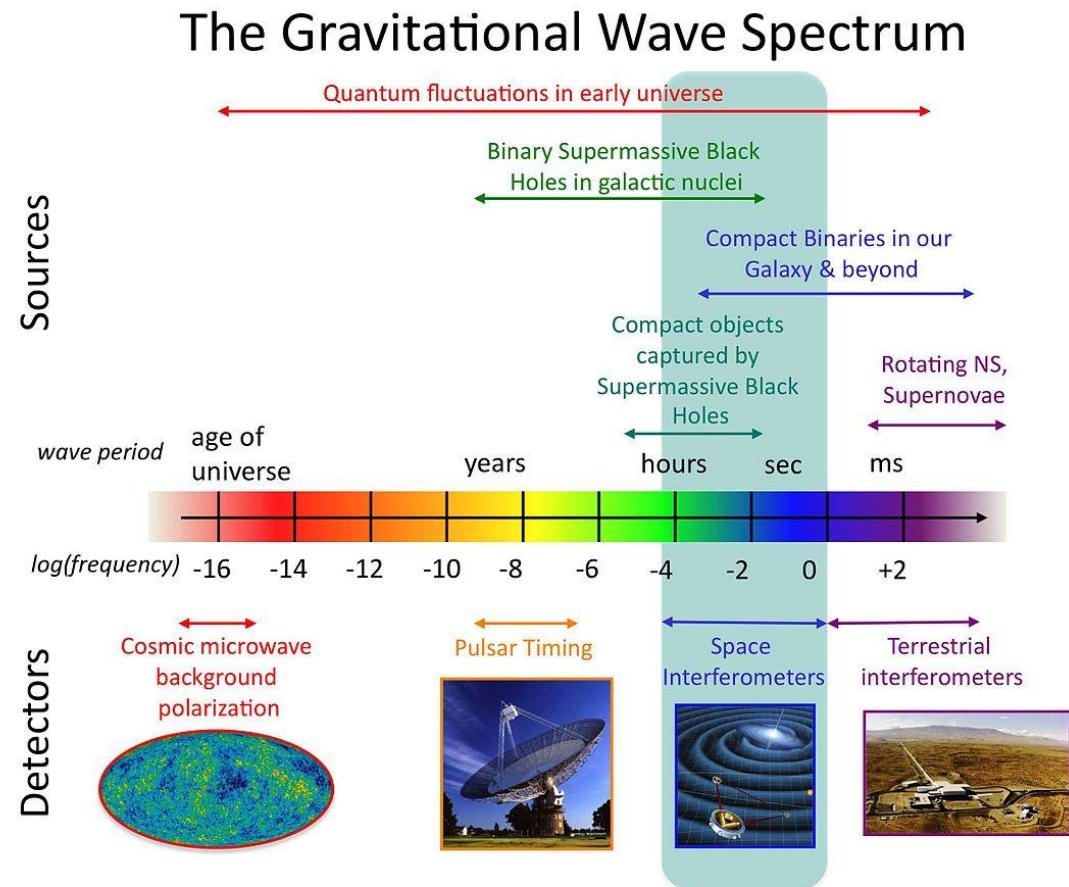
3.1 Construction of a Spectrometer

3.2 Excitation Sources

3.3 Dispersive Elements

3.4 Detectors

3.5 Sample Chamber



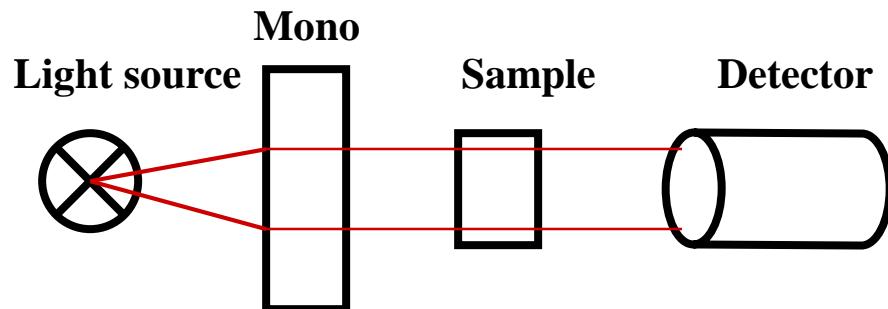
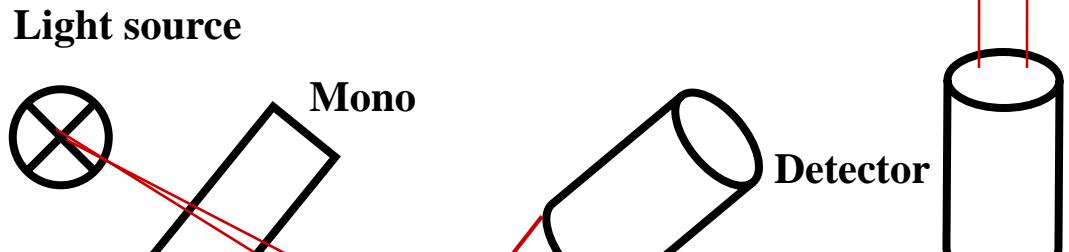
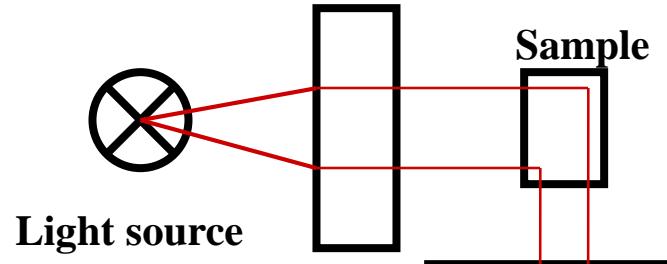
3.1 Construction of a Spectrometer

Geometries

Fluorescence, phosphorescence
of solutions
(rectangular geometry)

Reflection, scattering,
fluorescence, phosphorescence of
solids (rectangular geometry)

Absorption, transmission
of solutions, gases
(linear geometry)



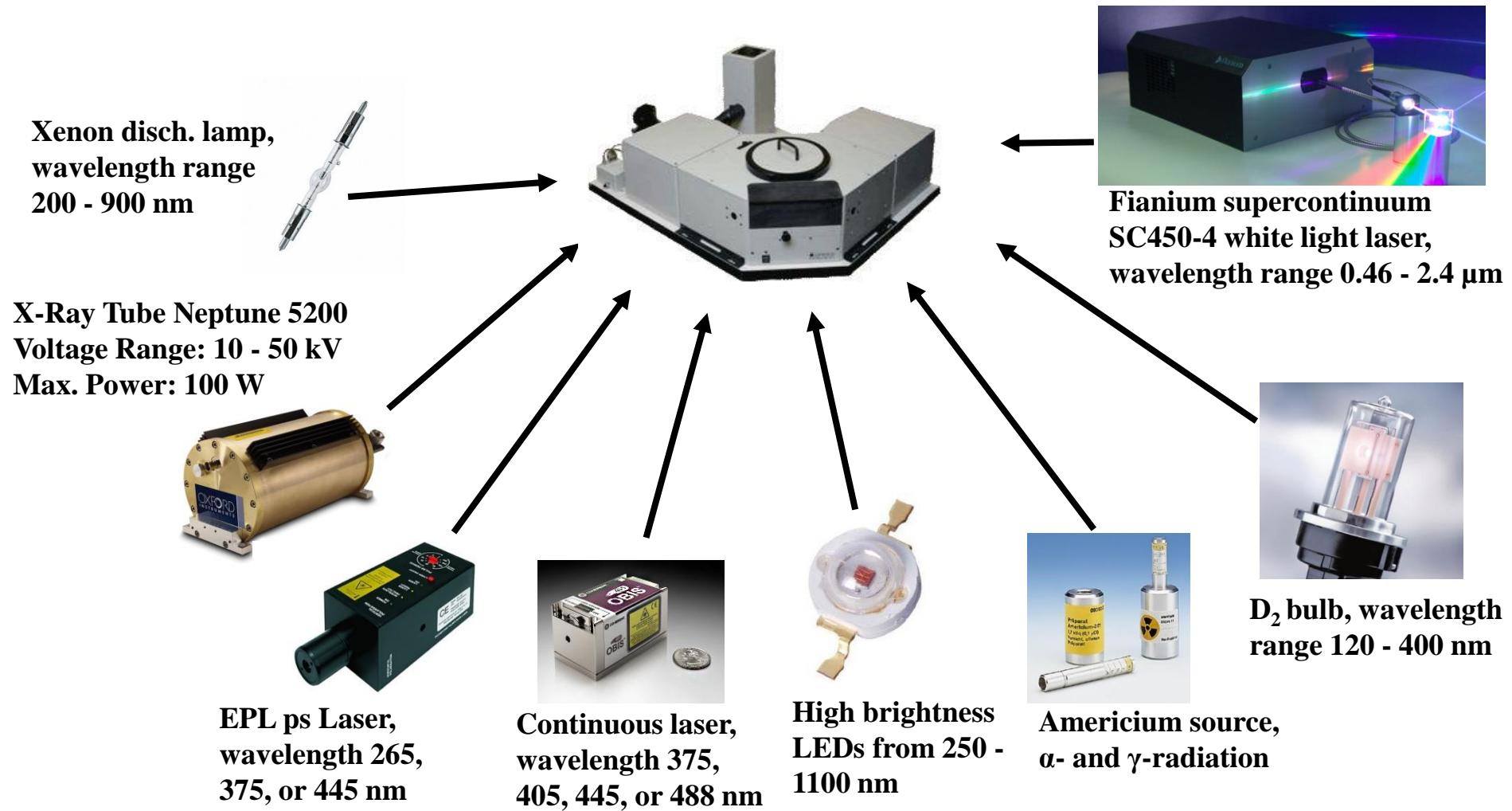
3.2 Excitation Sources

Typical radiation sources

- Celestial objects (only of interest for astrophysics) Continuous + Lines
- Halogen lamps Continuous
- Gas discharge lamps
 - Xe-medium pressure lamps Quasi-continuous
 - Low-pressure mercury lamps Lines
 - Medium pressure mercury lamps Quasi-continuous
 - High pressure mercury lamps Quasi-continuous
 - H₂/D₂-lamps Quasi-continuous
- LEDs narrow band
- Laser (less suitable for excitation spectra!) narrow line
 - Gas laser (CO₂, N₂, Kr, excimers: ArF*, XeF* etc.)
 - Solid-state laser
 - Dye laser
 - Laser diodes
- Synchrotrons (BESSY, DELTA, DESY, CERN, ...) White!

3.2 Excitation Sources

Radiation sources for fluorescence spectroscopy



3.2 Excitation Sources

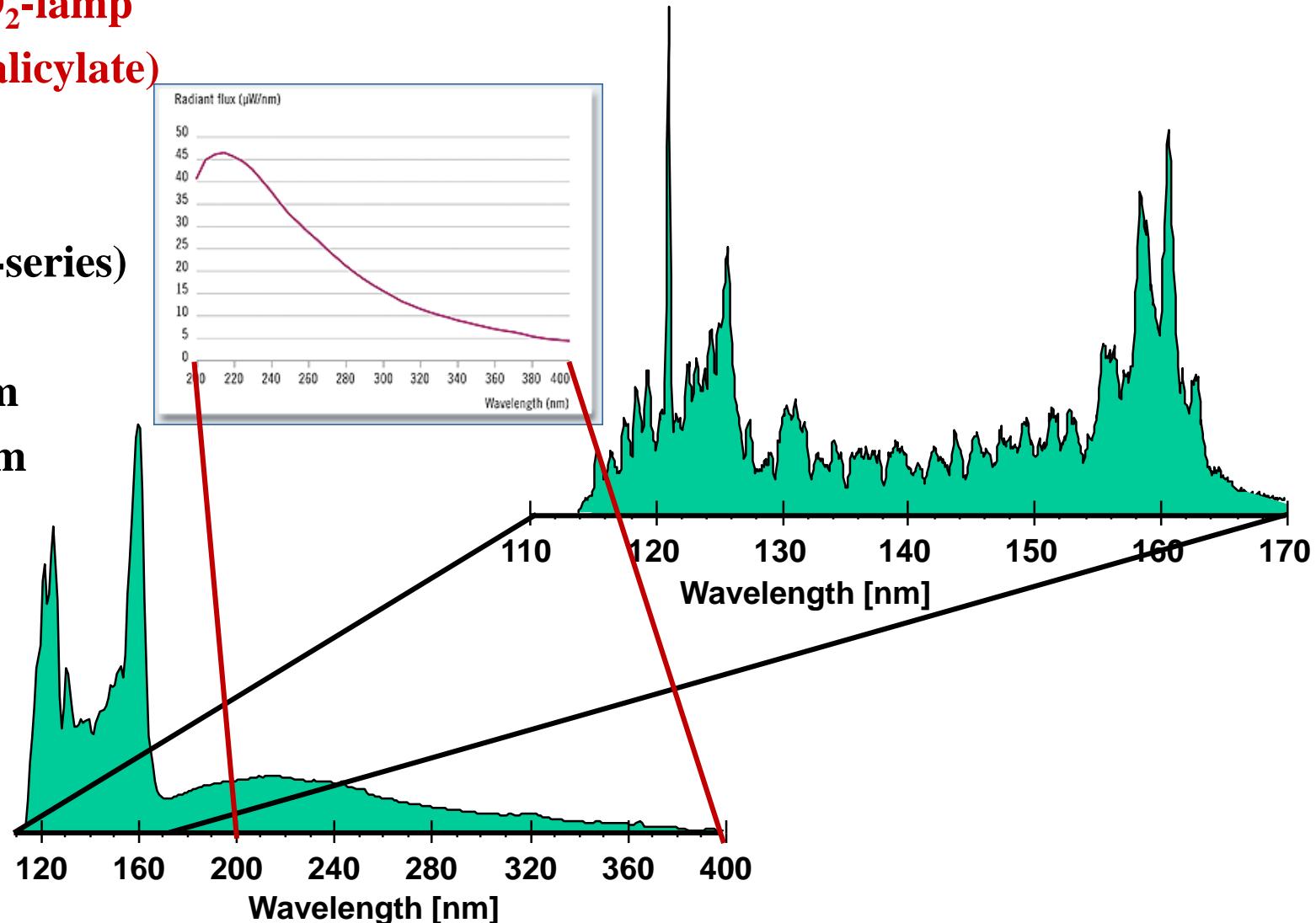
Spectrum of a D₂-lamp (Standard Na-salicylate)

110 – 170 nm

D-lines (Lyman-series)

170 nm – 400 nm

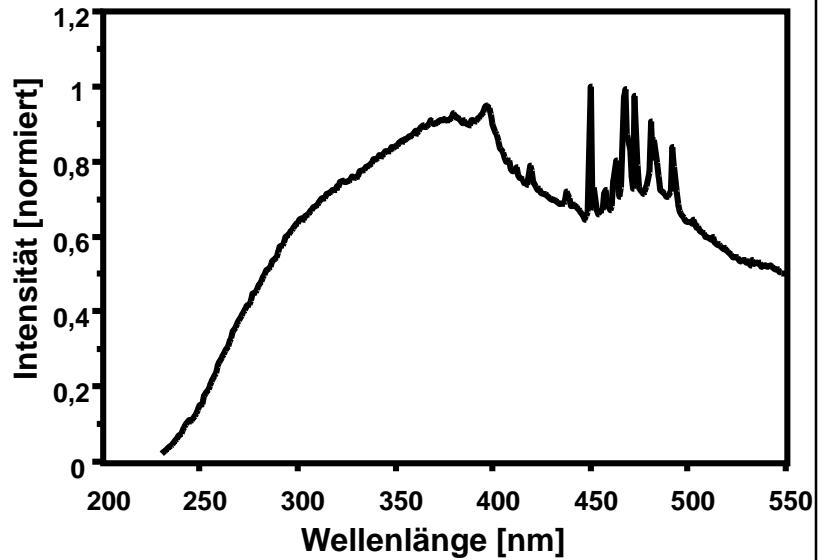
Quasi-continuum



3.2 Excitation Sources

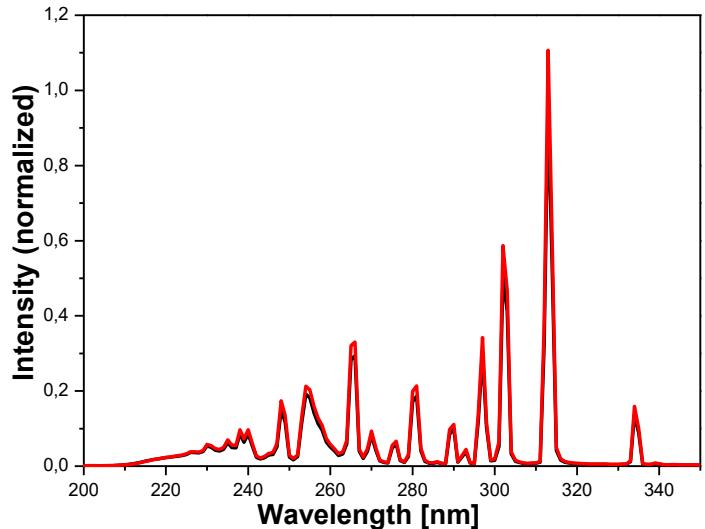
Spectrum of a Xe-medium pressure lamps

- 220 – 450 nm
Quasi-continuum
- > 450 nm
Quasi-continuum + Xe-lines



Spectrum of a high pressure mercury lamps

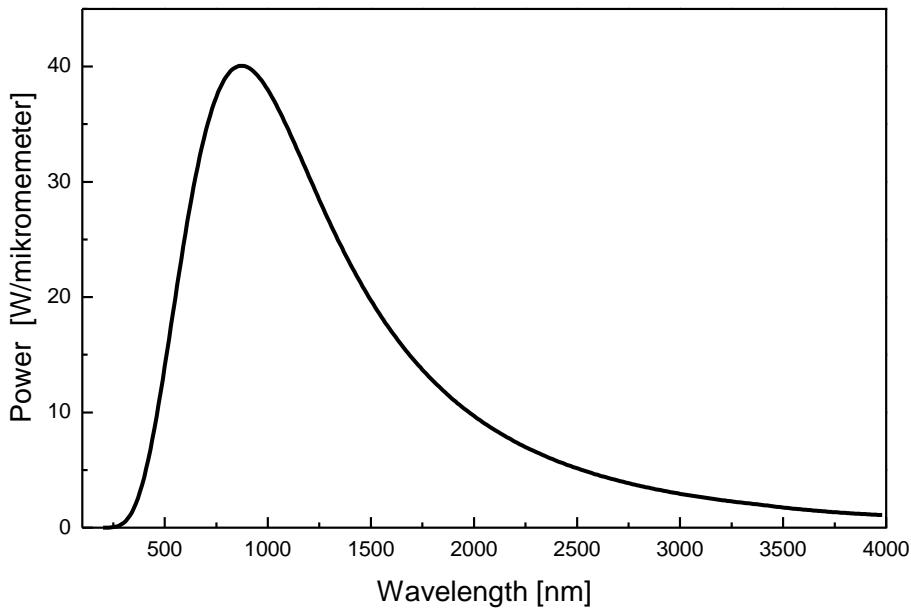
- 210 – approx. 300 nm
Quasi-continuum
- Hg-lines at 248, 254, 265, 303, 313, 365, 405, 435, 546 and 579 nm



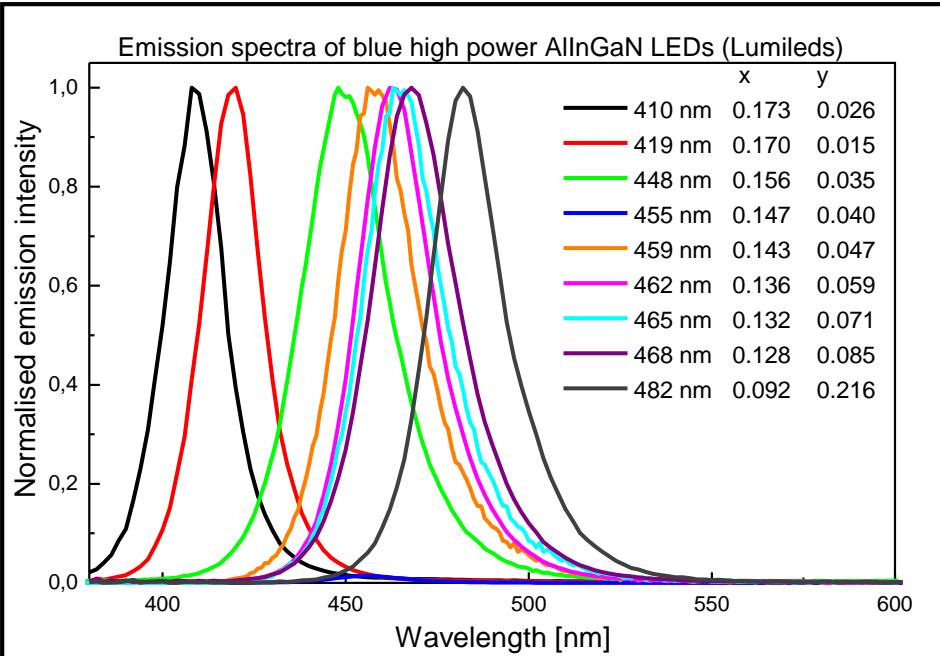
3.2 Excitation Sources

Spectra of halogen lamps and LEDs

Halogen lamp (broad band emitter)



LEDs (narrow band emitter)



LEDs are very flexible light sources (geometric and spectral)

(In,Ga)N **365 – 500 nm**

(Al,Ga)N **210 – 365 nm**

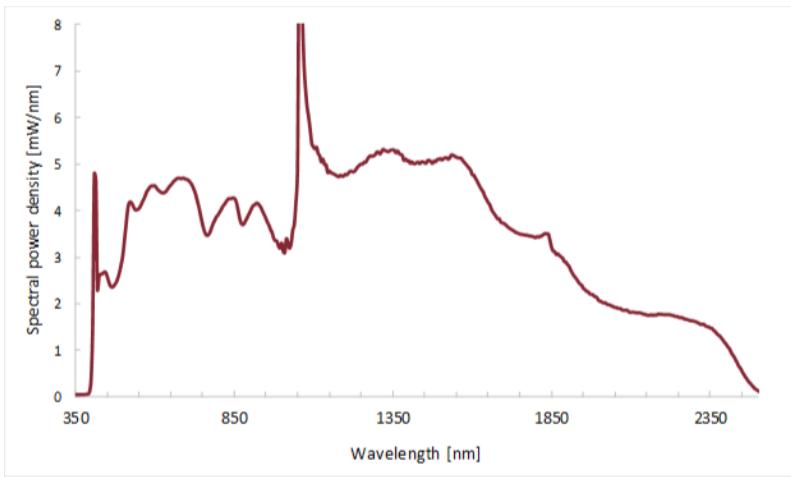
(Al,In,Ga)P **580 – 700 nm**

Ga(As,P) **> 650 nm**

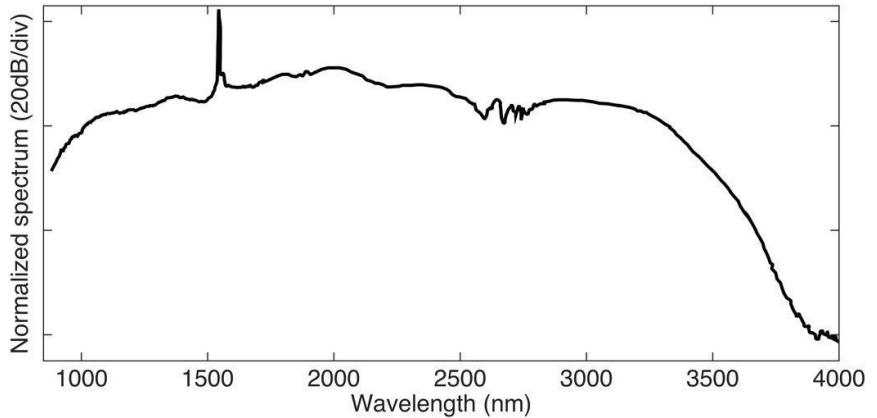
3.2 Excitation Source

Super continuum laser (white laser)

- Spectral broadening of a laser line via non-linear effects in a glass fiber
- In the spectrum the primary laser line can still be observed (e. g. 1064 nm of YAG:Nd)
- The optical power in W nm^{-1} is comparatively low, as the laser output power is spread over the whole spectrum



NKT Photonics „SuperK FIANIUM“
Optical power: 1.5 W

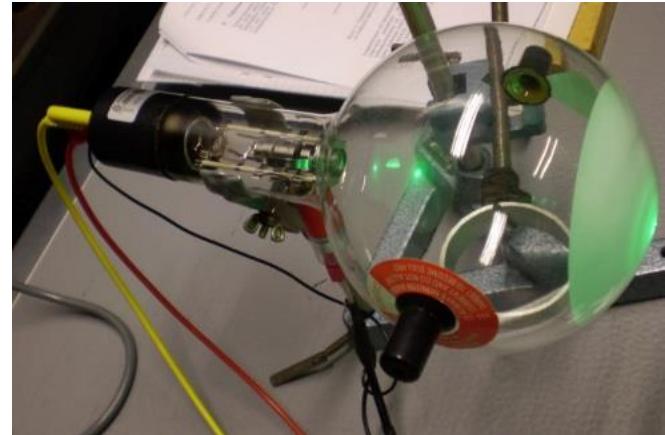
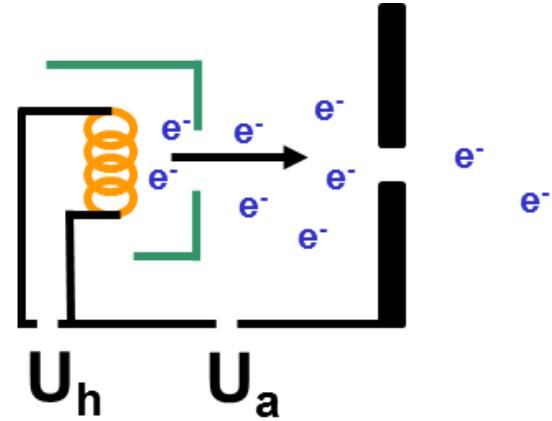


Experimental NIR SC-laser
Lit.: DOI 10.1117/2.1201510.006199

3.2 Excitation Sources

Electron gun

- By applying a high voltage between the **coil** and the anode plate, electrons are accelerated from the coil to the anode
- To allow electrons to be removed from the coil, it is electrically heated to 800 - 1000 °C
- The higher the temperature (heating current), the more electrons are emitted from the coil. The higher the voltage U_a , the higher is the energy of the electrons.
- The electron beam is focused by a positively charged **Wehnelt cylinder**

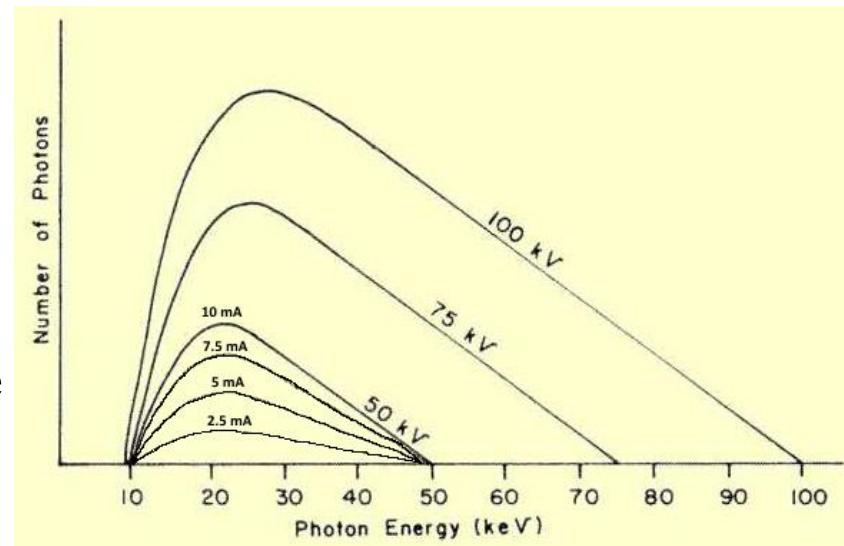
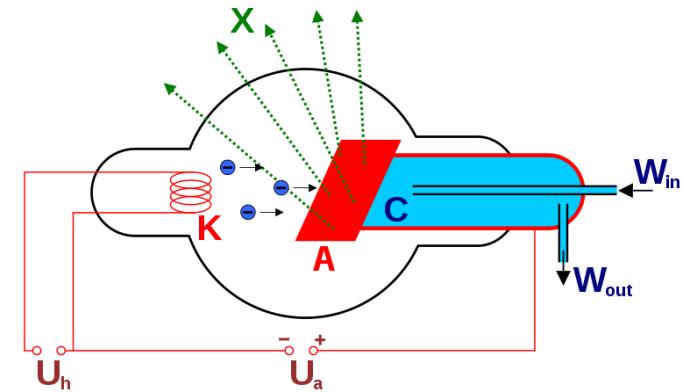


Source: Wikipedia

3.2 Excitation Sources

X-ray tubes

- Electrons emitted by a cathode (**K**) are accelerated and hit a target (**A**) (e. g. tungsten)
- The electrons are slowed down and emit the kinetic energy in the form of radiation (Bremsstrahlung) and heat
- The accelerating voltage **U_a** determines the intensity and spectral width of the x-rays
- The amount of electrons is controlled by the heating current and solely influences the intensity, not the shape, of the spectrum



$$\text{Efficiency (\%)} \approx U_a \cdot Z \cdot 10^{-4}$$

$$60 \text{ kV, tungsten: } 60 \cdot 74 \cdot 10^{-4} = 0.4\% (!)$$

Lit.: Perry Sprawls, "Physical Principles of Medical Imaging"

3.2 Excitation Sources

α -, β -, γ -emitter

α : He²⁺ cations

- typical α -emitters: ^{241}Am , ^{238}Pu , ^{239}Pu
- penetration depth: 10-15 μm (YAG, 5 MeV)

β : electrons, e⁻

- typical β -emitters: ^{137}Cs , ^{90}Sr , ^3H
- penetration depth: < 1 μm (10 keV), 8-10 μm (70 keV)

γ : high-energy photons

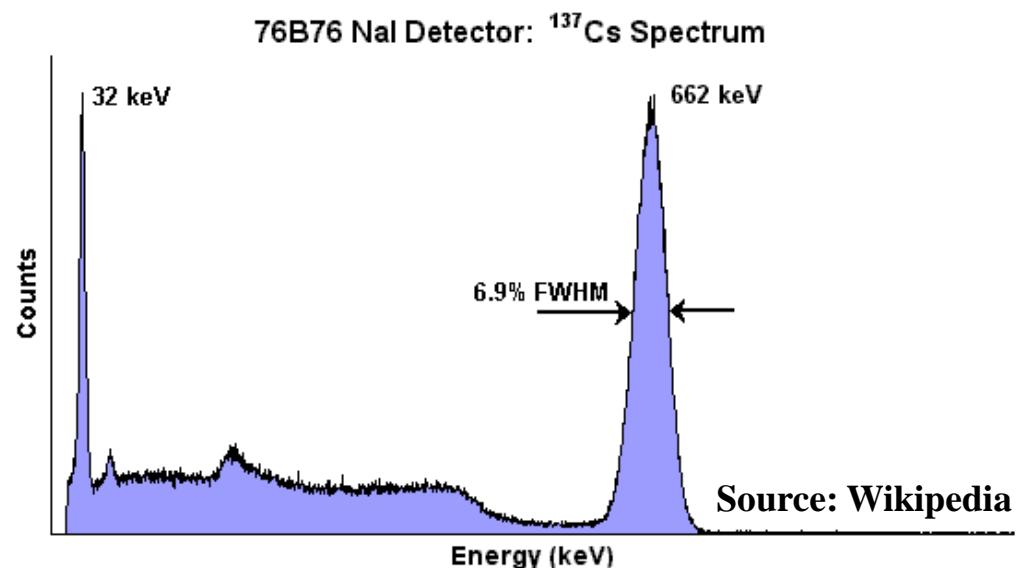
- typical γ -emitters: ^{137}Cs , ^{60}Co
- simultaneously emitted α or β particles are stopped by an thin metal absorber
- interacts with the whole particle



Tritium filled tube coated with $\text{Zn}_2\text{SiO}_4:\text{Mn}^{2+}$



Commercial β -source

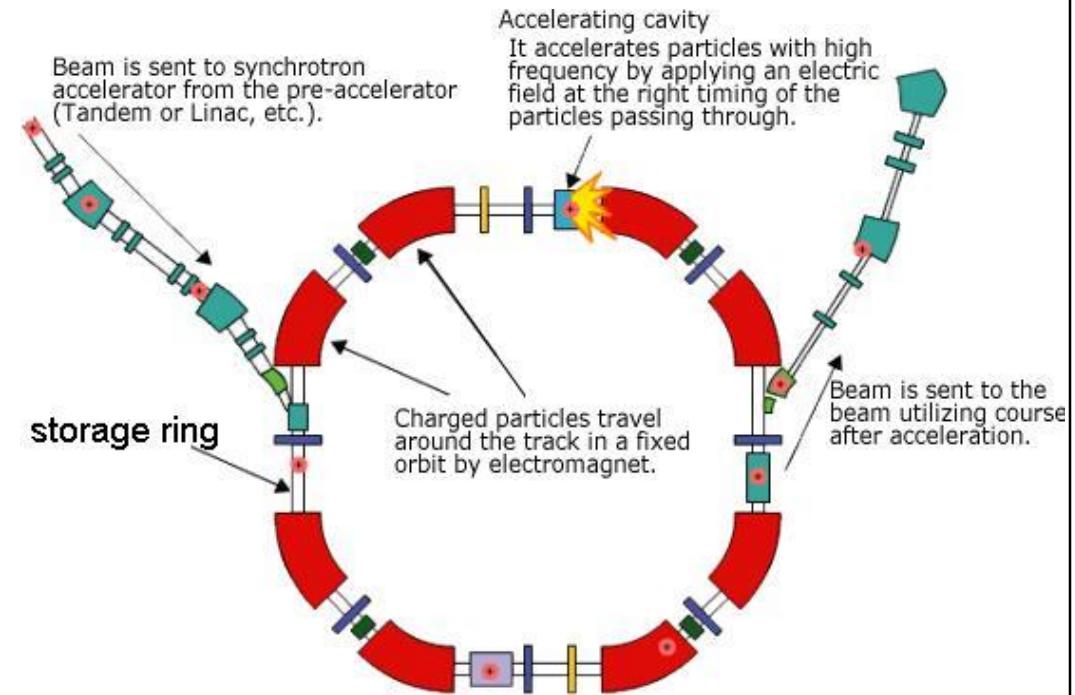


3.2 Excitation Sources

Synchrotron

Generates a continuous spectrum by accelerating electrons

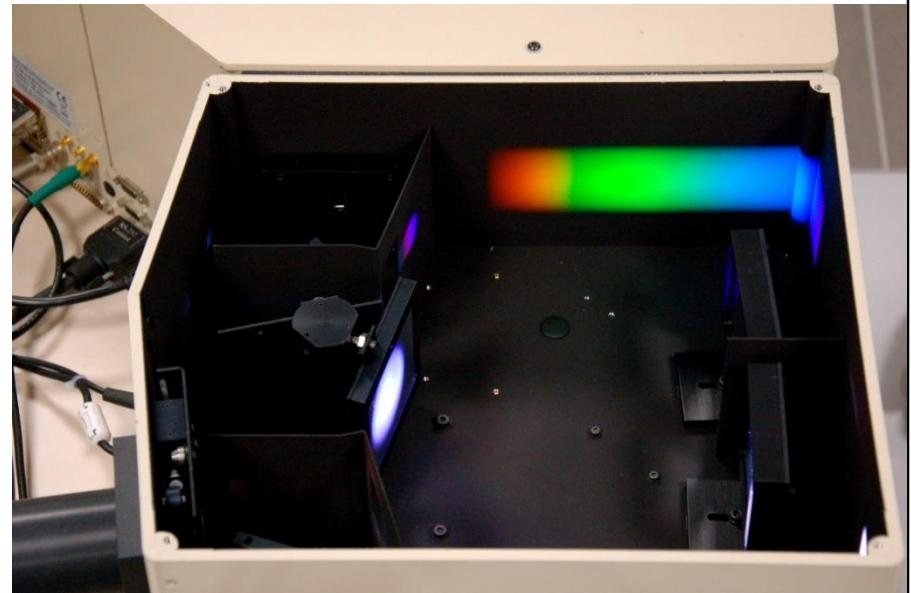
- Energy (wavelength) depends on the acceleration and can range from $< 1 \text{ eV}$ ($> 1 \mu\text{m}$, NIR) to $> 1 \text{ GeV}$ ($< 1 \text{ fm}$)
- Special magnet setups cause the spectrum to get narrower
- Very expensive and only used for specific experiments where high-intensity or short-pulsed radiation is required



Source: Wikipedia

3.3 Dispersive Elements

- Many excitation sources generate light over a broad spectral range. The dispersive element allows to select a specific wavelength
- The common PMT type detectors cannot distinguish between photons of different wavelength. Thus, the light emitted by the sample also has to be fed through a dispersive element
- Notch, short or long pass filters are be a simple method for wavelength selection, but they do not allow dynamic selection. For that reason prisms or gratings are used.



3.3 Dispersive Elements

- **Dispersion prisms**

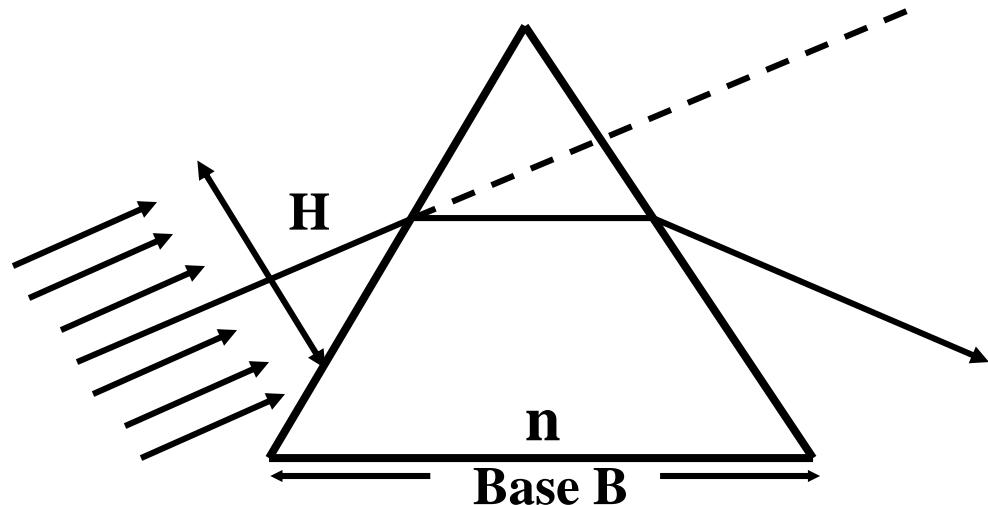
Angular dispersion:

$$d\delta/d\lambda = - B/H * dn/d\lambda$$

B = Prism base

H = Height of beam (slit width)

n = Refractive index



- **Dispersion gratings**

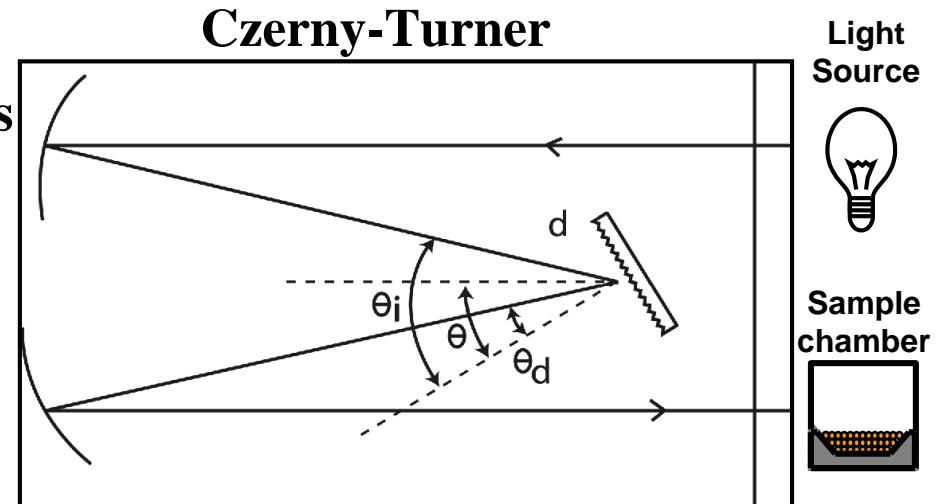
Grating constant (lines / mm) determines
the usable area \Rightarrow “Gratings”

e.g. 1200 or 2400 lines / mm

Resolution : $R = \lambda / d\lambda = m * N$

m = Order

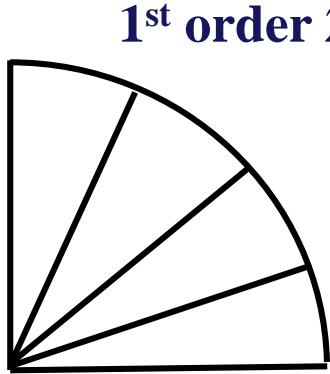
N = Slit number



Attention! Reflections of higher order \Rightarrow Filter

3.3 Dispersive Elements

Higher orders upon use of a grating



1st order 250 nm
1st order 500 nm
2nd order 250 nm (weak)
1st order 750 nm
2nd order 375 nm (weak)
3rd order 250 nm (very weak)

The second order has to be removed with a filter!

Example: 350 nm excitation, emission recorded from 450 to 800 nm

The grating will group 2nd order 350 nm photons from the lamp together with the 1st order 700 nm photons from the sample.....

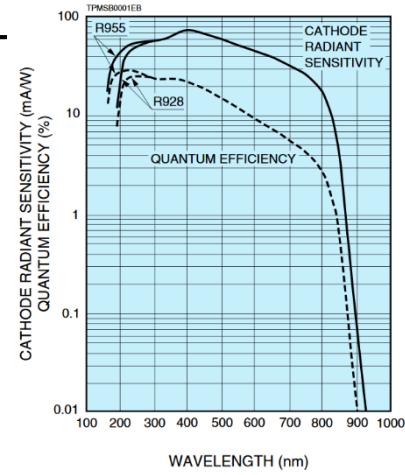
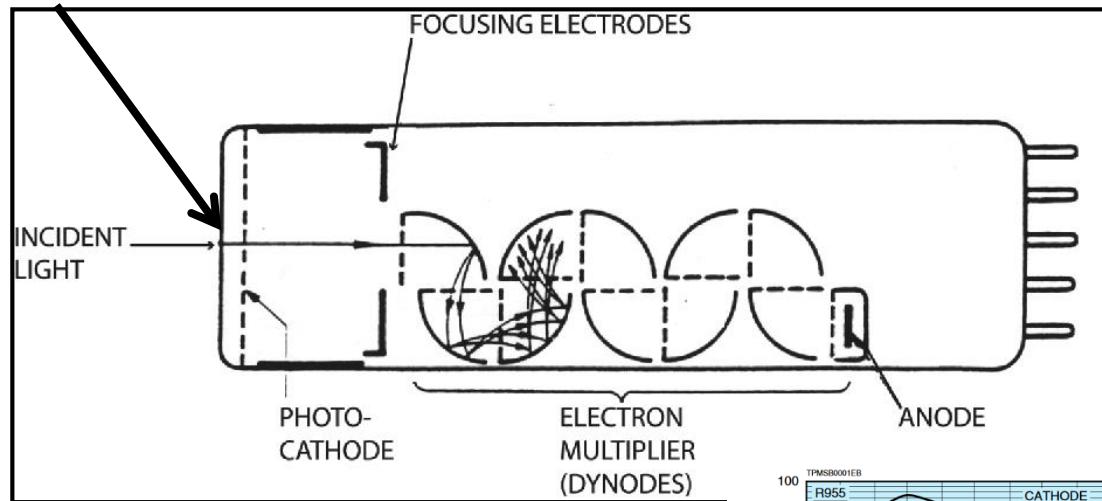
Solution: A 400 nm long pass filter behind the sample blocks the 350 nm photons

3.4 Detectors

Types of detectors

- With external photoeffect
 - Photocells
 - Photomultiplier (PMT) →
- With internal photoeffect
 - Photoresistors
 - Photodiodes
- Photochemical detectors
 - for absolute measurements (actinometry)
 - Photochemical reaction with a known quantum yield QY
 - $2 \text{ [Fe}^{\text{III}}(\text{C}_2\text{O}_4)_3]^{\text{3-}} \rightarrow 2 \text{ Fe}^{\text{II}} + 5 \text{ C}_2\text{O}_4^{\text{2-}} + 2 \text{ CO}_2 \uparrow$
 - Reaction of Fe^{2+} with phenanthroline + measuring the absorbance at 510 nm
 - $\text{Photons/s} = \text{number of formed product's molecules}/(\text{QY} * \text{irradiation time})$

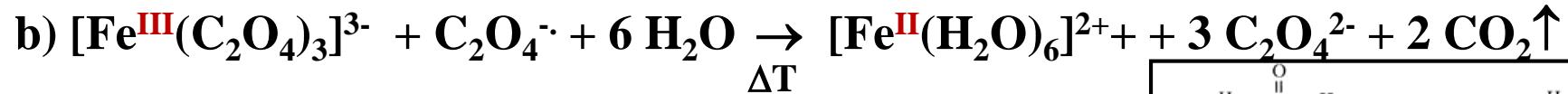
Window



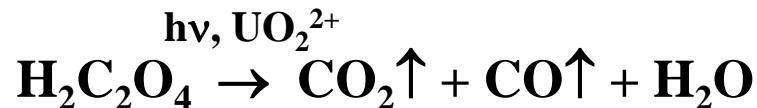
3.4 Detectors

Actinometer (Photochemical detectors)

- Ferrioxalate actinometer (UV – 500 nm)



- Uranyloxalat actinometer (208 – 426 nm)

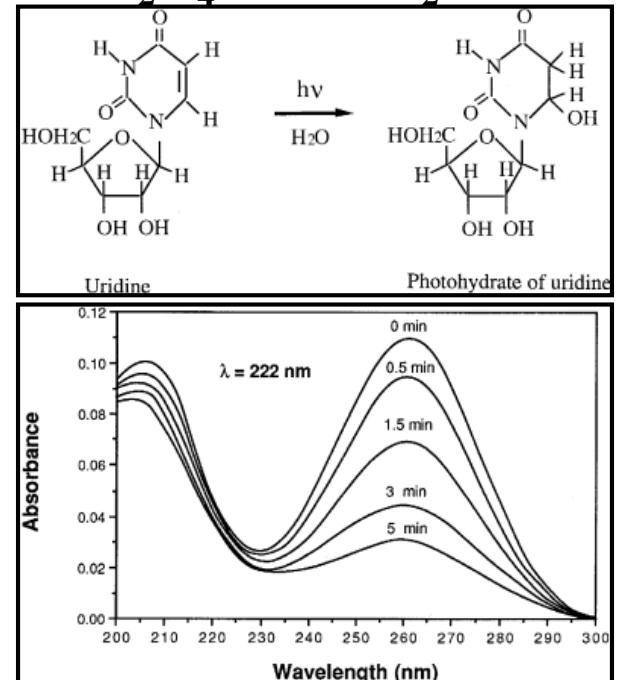


- KI/KIO₃ actinometer (UV/Vis \leftrightarrow Grätzel cell)



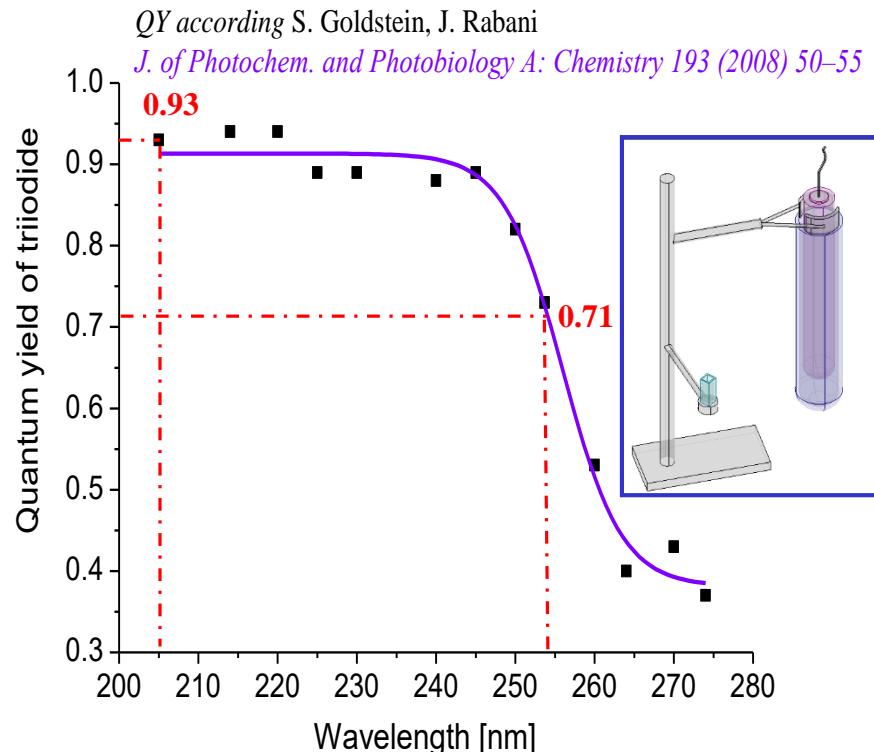
- Uridine actinometer (200 – 290 nm)

Lit.: J.-Y. Zhang et al., Applied Surface Science 109/110, 1997, 482



3.4 Detectors

Excursion: Triiodid Actinometer: $8 \text{ KI} + \text{KIO}_3 + 2 \text{ H}_2\text{O} \rightarrow 3 \text{ I}_3^- + 6 \text{ OH}^- + 9 \text{ K}^+$



- A = Absorbance at 352 nm (UV/Vis spectroscopy)
d = 1 cm cuvette
V = 3 cm³
S: Irradiated surface 3 cm²
F: Represent photon energy [mJ·mol⁻¹] at emission wavelength of radiation source
Φ = Quantum yield of actinometer at emission wavelength of radiation source
ε = 27600 M·cm: molar absorption coefficient of I₃⁻ at 352 nm
h = Planck constant = 6.626·10⁻³⁴ Js
c = Speed of light in vacuum = 3·10⁸ ms⁻¹
N_A = Avogadro constant = 6.022·10²³ mol⁻¹
λ = Wavelength of emitted radiation, e.g. 254 nm

$$\text{Energy of photons/mol F} = \frac{h \cdot c \cdot N_A}{\lambda}$$

$$\text{Irradiance H} = \frac{A \cdot V \cdot F}{\emptyset \cdot \epsilon \cdot d \cdot S} [\text{mW/cm}^2]$$

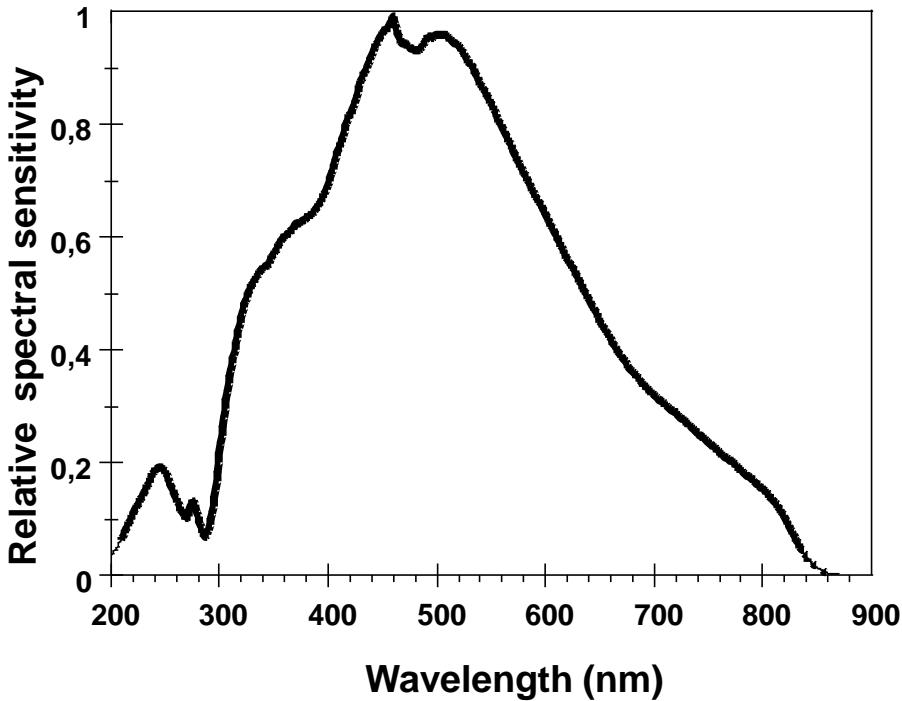
3.4 Detectors

Spectral sensitivity

Photomultiplier tube (PMT)

200 – 850 nm

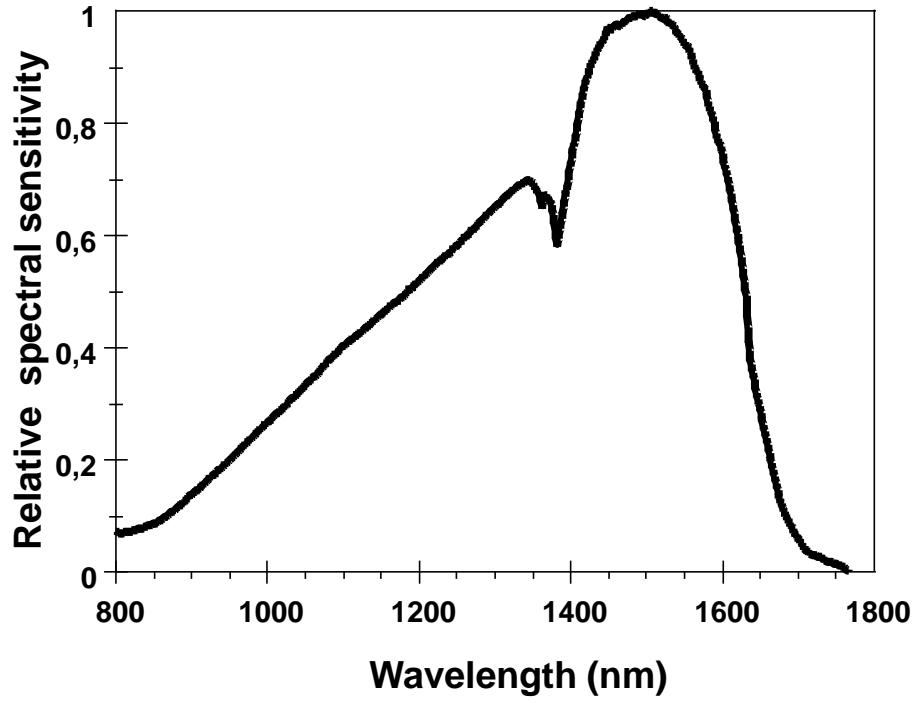
(Peltier element cooling ~ 253 K)



Ge semiconductor detector

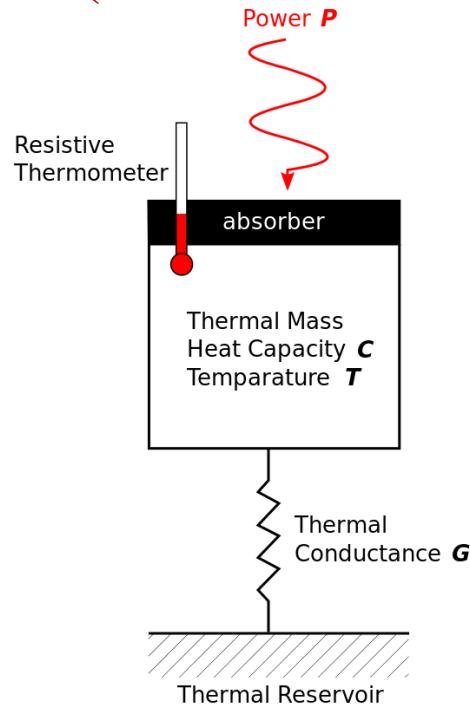
800 – 1750 nm

(Liquid N₂ cooling 77 K)



3.4 Detectors

Bolometer (micro and radio waves)



- Measurement of the spectral intensity via an increase in temperature caused by absorption of photons
- Slow, but very high sensitivity

Scintillation detector (ionizing radiation)

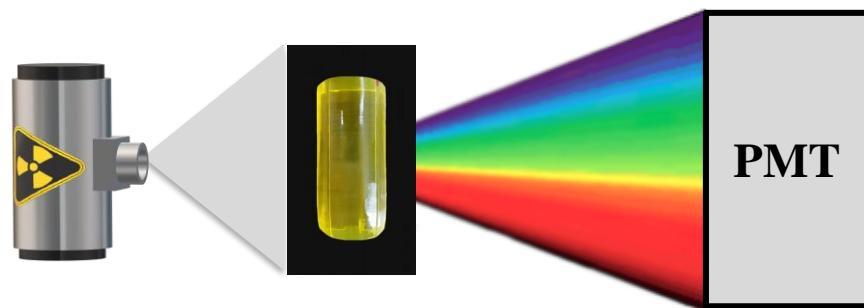


- Ionizing radiation is converted into visible radiation and detected by a common PMT or CCD
- Fast and very high sensitivity

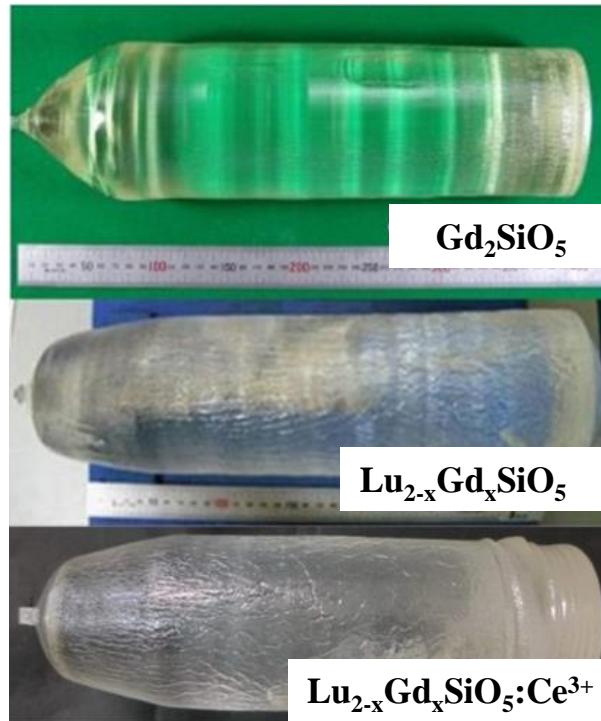
3.4 Detectors

Scintillators

Scintillators absorb ionizing radiation and convert it to optical radiation, especially UV or visible radiation → X-ray and gamma detectors or cameras, space telescopes



PMT = PhotoMultiplier Tube



Scintillator materials (selection)

- $\text{Lu}_3\text{Al}_5\text{O}_{12}:\text{Ce}$ 525 nm
- $\text{Bi}_4\text{Ge}_3\text{O}_{12}$ 500 nm
- PbWO_4 420 nm
- $\text{Lu}_{2-x}\text{Gd}_x\text{SiO}_5:\text{Ce}$ 420 nm
- $\text{NaI}:\text{Tl}$ 415 nm
- $\text{Lu}_2\text{Si}_2\text{O}_7:\text{Ce}$ 380 nm
- $\text{LuPO}_4:\text{Ce}$ 350 nm
- $\text{LuPO}_4:\text{Pr}$ 235 nm

Inorganic scintillators are the technological backbone of modern medical diagnostic techniques, viz. CT, SPECT, PET

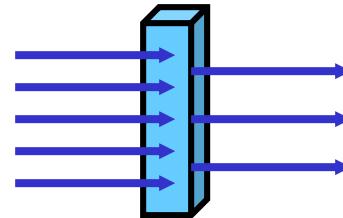
Source: <https://www.opt-oxide.com/en/singlecrystals> (03.12.2018)

3.5 Sample Chamber

Geometry

Absorption spectroscopy

- Linear arrangement: laser source - mono - sample - detector
- One-or two-beam arrangement



Fluorescence spectroscopy

- Rectangular arrangement: laser source - mono1 - sample - mono2 - detector
- Sample is lying (powder sample holder) or standing (quartz cells for solutions)



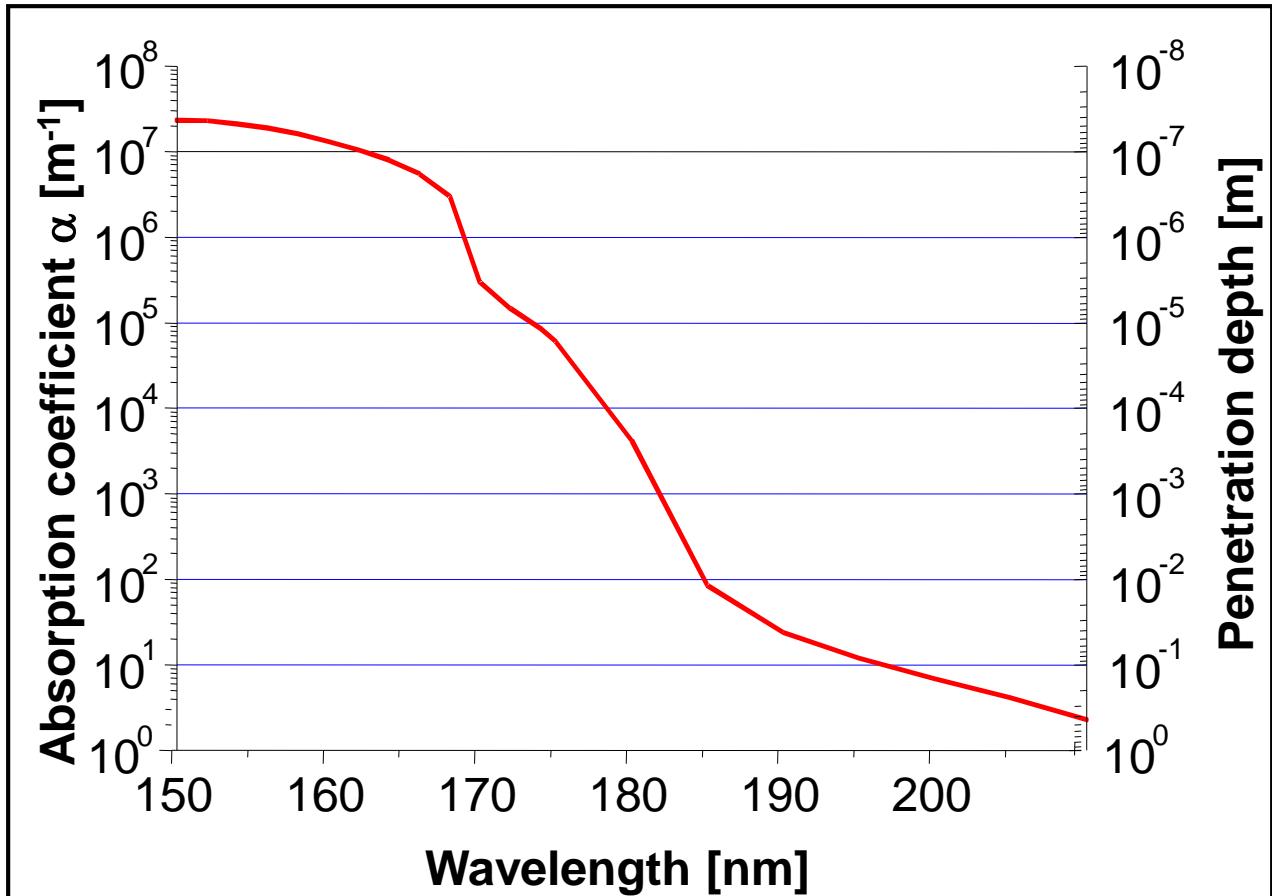
3.5 Sample Chamber

Atmosphere

H_2O and O_2 absorb radiation below approximately 200 nm

UV/Vis-measurements
(200 – 800 nm excitation)
⇒ air /water

VUV- measurements
(100 – 200 nm excitation)
⇒ Vacuum, Ar, or N_2

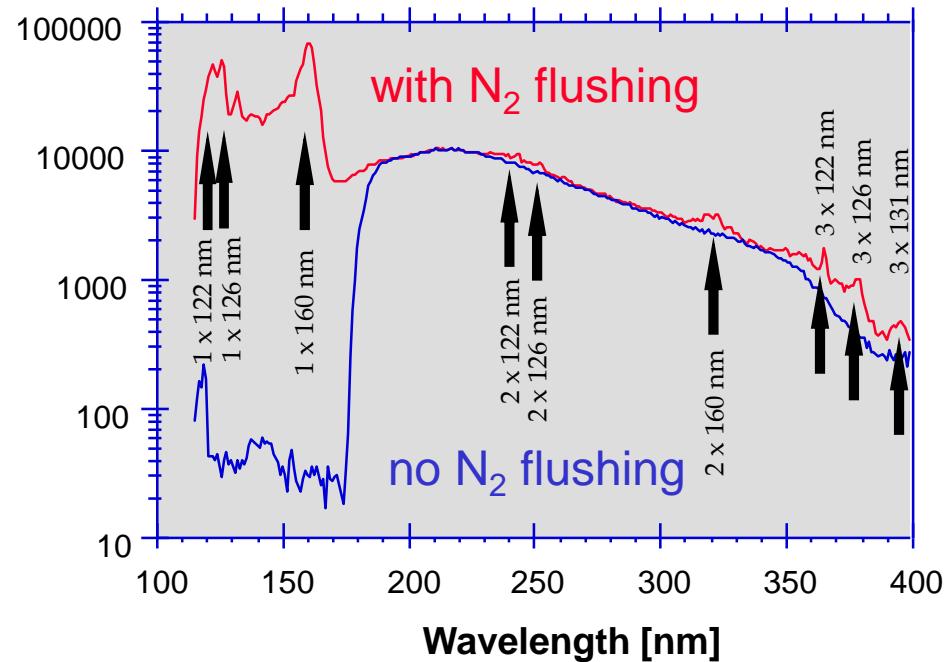
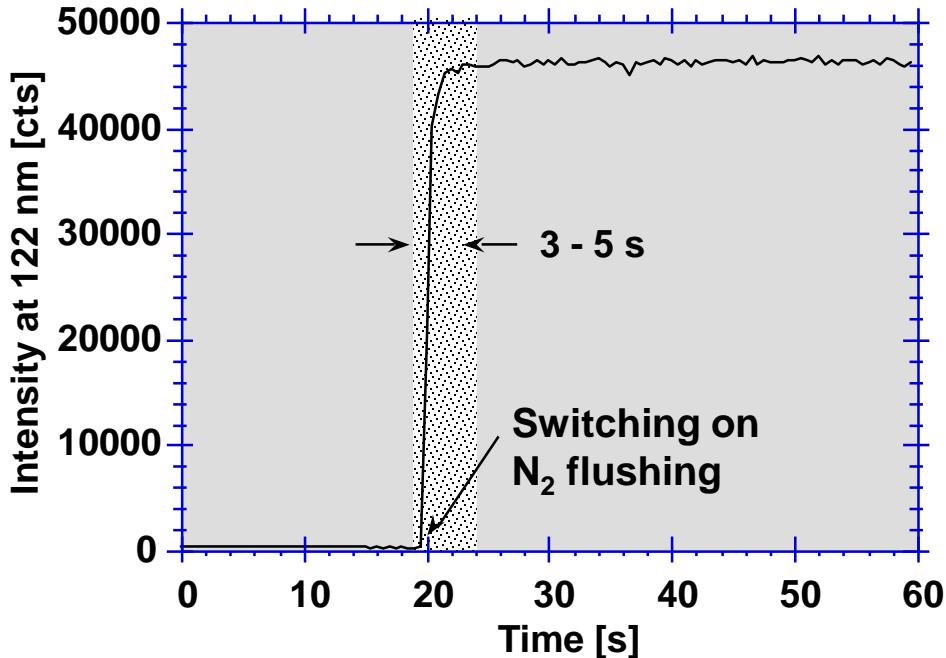


3.5 Sample Chamber

Atmosphere

Effect of flushing with an inert gas, e.g. N₂

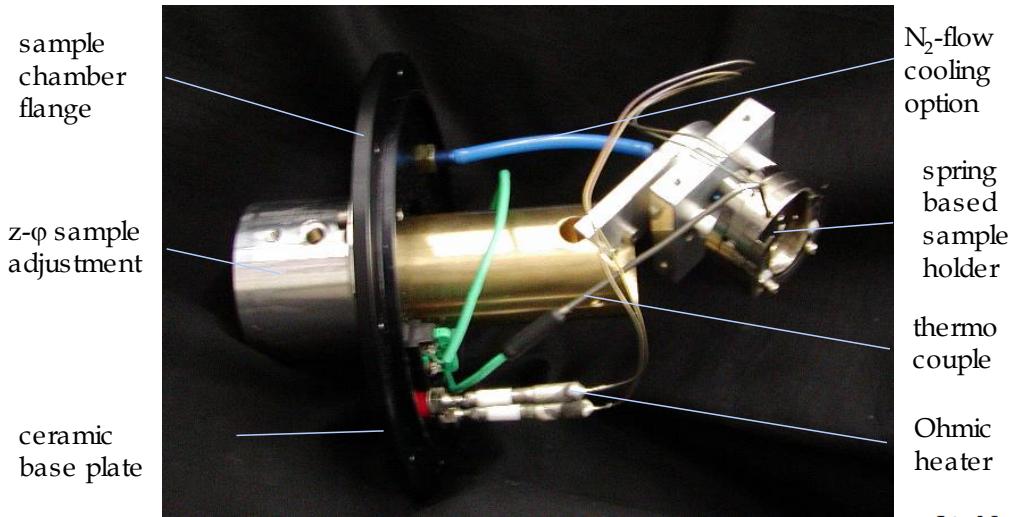
1. Transmission up to about 120 nm
2. Loss of the filtering effect for higher-order maxima (2nd & 3rd order)



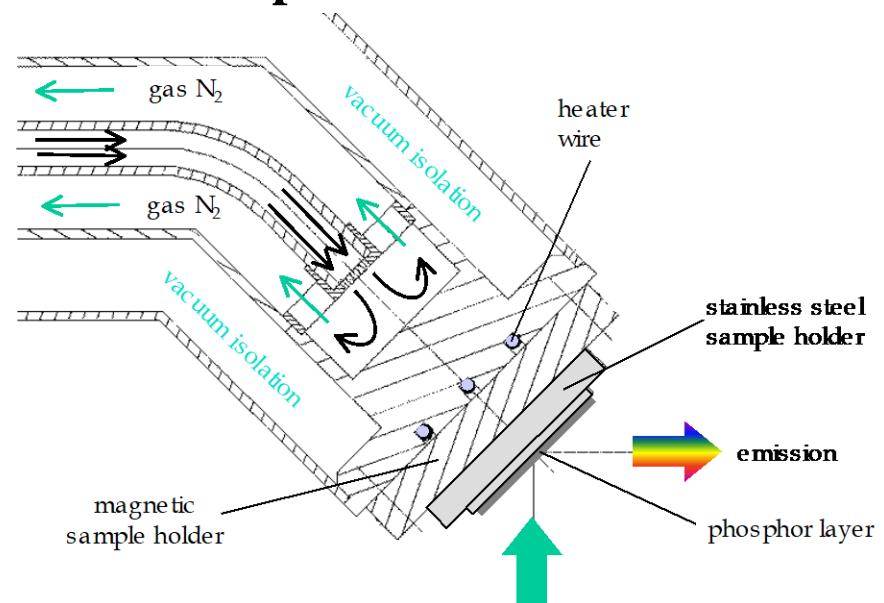
3.5 Sample Chamber

Depending on required conditions, the sample can be heated or cooled

Sample holder with heater



Schematics of a coolable sample holder



3.5 Sample Chamber

He cryostat

- Temperatures around 0.8 – 4.2 K
- Continuous flow cryostats have a high He consumption (cost ~ 7 €/l)
- Closed-cycle cryostats cool the He vapour have a high energy consumption
- Pulse tube cryostats do not consume He. They work similar to a Refrigerator. Aquisition is rather expensive ~ 50,000 €



Pulse tube cryocooler
(Image: Oxford Instruments)

4. Methods of Optical Spectroscopy

4.1 Absorption Spectroscopy

4.1.1 Bouguer-Lambert-Beer-Law

4.1.2 Validity of the Lambert-Beer-Law

4.1.3 Deviation from the Lambert-Beer-Law

4.1.4 Extinction of Biologically Relevant Molecules

4.1.5 Atomic Absorption Spectroscopy

4.1.1 Bouguer-Lambert-Beer-Law

Derivation

The weakening of the intensity is proportional to the intensity and the layer thickness (**Bouguer-Lambert**):

$$dI \sim I \cdot dx \quad \text{or} \quad dI = -\alpha(\lambda) \cdot I \cdot dx$$

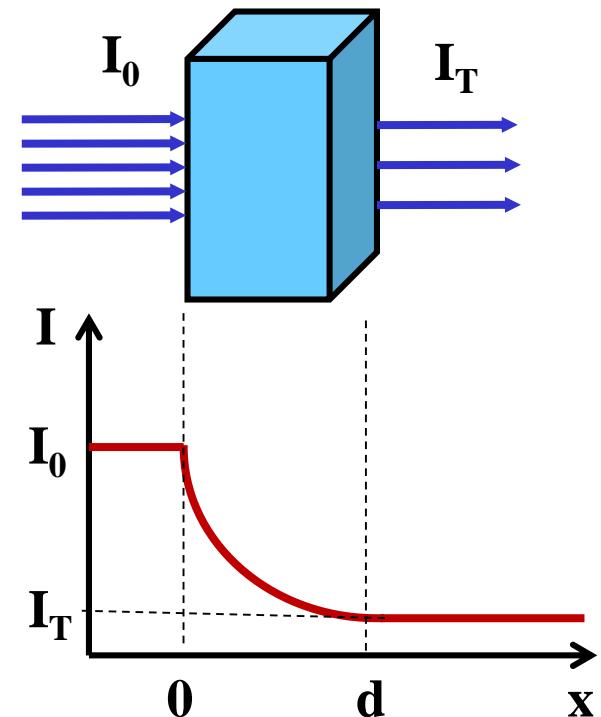
The proportionality factor $\alpha(\lambda)$ is proportional to concentration c (**Beer 1852**):

$$dI = -\alpha(\lambda) \cdot c \cdot I \cdot dx \Rightarrow \frac{dI}{I} = -\alpha(\lambda) \cdot c \cdot dx$$

Integration yields $\ln(I_0/I_T) = \alpha(\lambda) \cdot c \cdot x$ or $\log(I_0/I_T) = \alpha(\lambda) \cdot c \cdot x / \ln 10$

Combining gives $\log(I_0/I_T) = \varepsilon(\lambda) \cdot c \cdot x = A$ with $\varepsilon(\lambda) = \alpha(\lambda) / \ln 10$

$\ln(x) = \log(x) \cdot \ln(10)$ or $\log(x) = \ln(x) / \ln(10)$

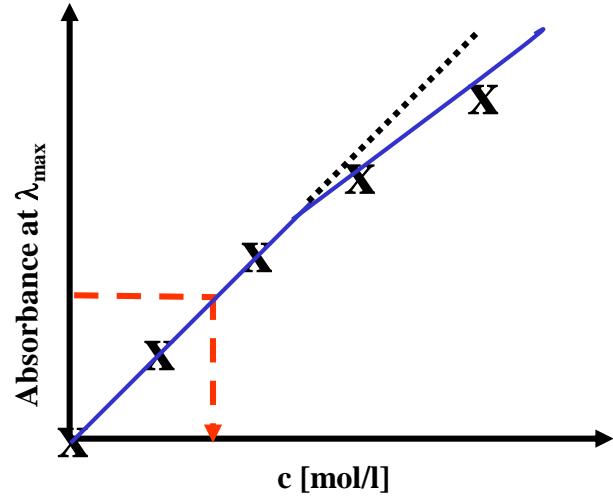


4.1.2 Validity of the Lambert-Beer-Law

The formulation and the validity depends on several boundary conditions

Conditions

1. The incident light must be collimated and monochromatic
2. The absorbing molecules must be highly diluted that a mutual interaction of the chromophoric groups of different molecules does not occur
3. The pure solvent does not absorb radiation of measuring wavelength
4. Radiation losses due to reflection from the plane-parallel walls of the cuvette and by scattering by particles does not occur



4.1.3 Deviation from the Lambert-Beer-Law

Apparent and real deviations

Apparent: Technical measurement inadequateness

- Sieve effects: scattering centers in the solution
- Dissociation or association equilibria, e.g.



- Lack of monochromaticity of the applied radiation
(false light I_s) leads to a background signal that reduces the dynamic
Maximum absorption dynamics $A = 1.04$ for 10% false light
Maximum absorption dynamics $A = 4.00$ for 0.01% false light
(typical of single monochromators)

$$A = \log \frac{I_0 + I_s}{I + I_s}$$

Real: Intermolecular interaction

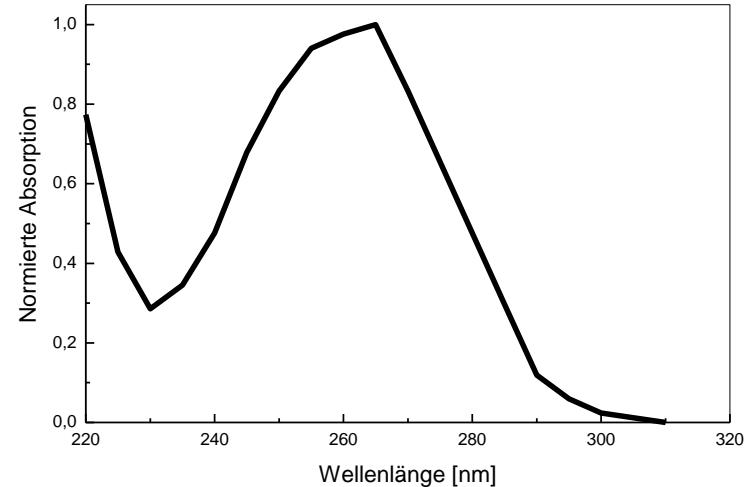
- Too high concentration $> 10^{-3}$ mol/l
⇒ Mutual interaction of the chromophores in the solution

4.1.4 Extinction of Biologically Relevant Molecules

Analysis of biomolecules

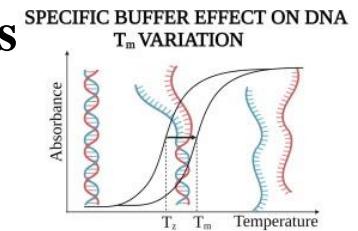
Species	λ_{\max} [nm]	ϵ [$\text{Imol}^{-1}\text{cm}^{-1}$]	Transition
Cytochrome c Fe ²⁺	420	120000	$\pi-\pi^*$
Cytochrome c Fe ³⁺	410	110000	$\pi-\pi^*$
Hemoglobin	560	12000	$\pi-\pi^*$
Oxyhemoglobin	550, 574	13000, 14000	$\pi-\pi^*$
Adenine	260	13400	n- π^* , $\pi-\pi^*$
Guanine	275	8100	n- π^* , $\pi-\pi^*$
Cytosine	267	6100	n- π^* , $\pi-\pi^*$
Thymine	264	7900	n- π^* , $\pi-\pi^*$
AMP	260	15500	n- π^* , $\pi-\pi^*$
ss-poly-AMP	260	10600	n- π^* , $\pi-\pi^*$
ds-poly-ATMP	258	9600	n- π^* , $\pi-\pi^*$

Absorption spectrum of dTMP



Applications areas

- Fusing of DNA
- Oxidation state of hemoglobin and Cytochromes



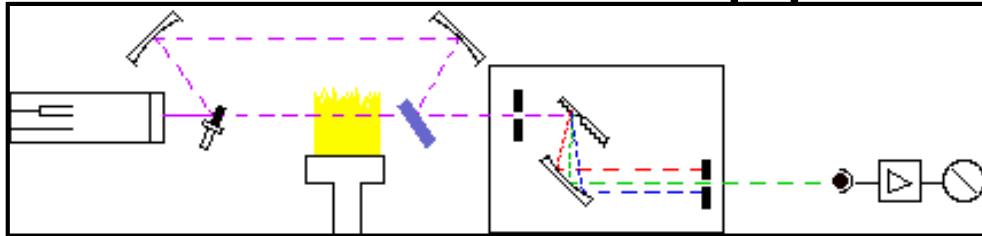
DNA stabilization buffer series:
Tris > Cacodylate > Phosphate > Citrate

4.1.5 Atomic Absorption Spectroscopy

Principle and application of Atomic Absorption Spectroscopy (AAS)

Kirchoff's law: Any material can also absorb the radiation, which was emitted by itself (each element has a characteristic line spectrum).

Basic construction of an AAS equipment



The weakening of the emission (absorption) is determined by a calibration process

Some application areas of AAS

- Metal analysis in alloys
- Lead in fuel
- Pollutants in environment
- Mg in plants

Hollow cathode lamps

Burners (air / acetylene: 2300 °C)

Monochromator or filter

Detector

4. Methods of Optical Spectroscopy

4.2 Luminescence Spectroscopy

4.2.1 Definition and Applications

4.2.2 Construction of a Fluorescence Spectrometer

4.2.3 Emission Spectroscopy

4.2.4 Excitation Spectroscopy

4.2.5 Quantum Yield

4.2.6 Light output

4.2.7 Temperature Resolved Spectroscopy

4.2.8 Time Resolved Spectroscopy

4.2.1 Definition and Applications

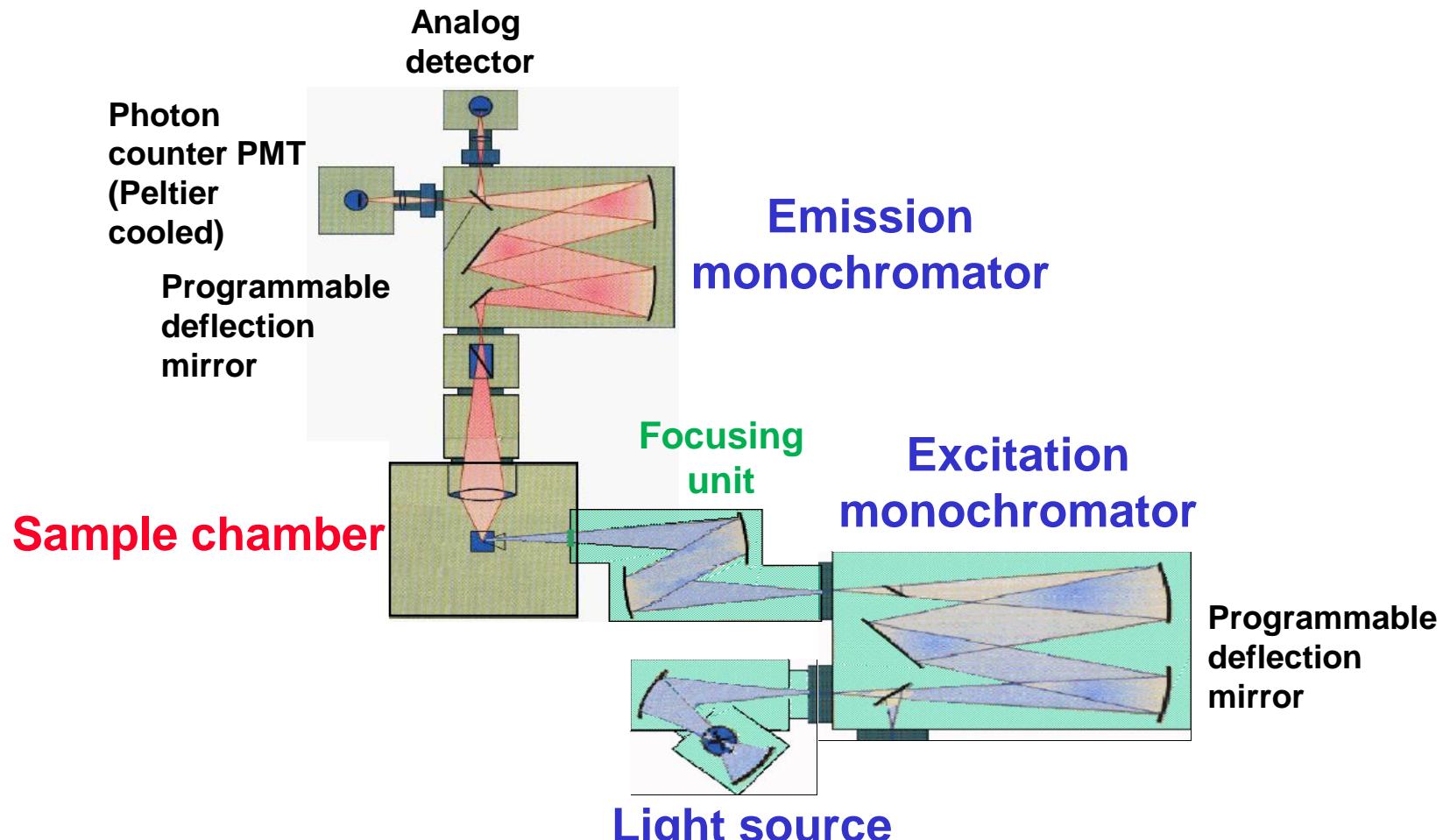
Definition

Luminescence is the emission of radiation by a system (solid, liquid or molecule) in a non-equilibrium state after energy has been supplied to the system

Applications

- Characterization of the spectral energy distribution of the emission of liquid or solid samples (glass, ceramics, solutions, or powders)
⇒ Emission spectra
- As a function of excitation energy/wavelength
⇒ Excitation spectra
- As a function of temperature
⇒ Thermal quenching
⇒ Glow curves and thermo luminescence
- As a function of time after the excitation pulse
⇒ Decay curves and constants

4.2.2 Construction of a Fluorescence Spectrometer



Typical excitation wavelength from 100 to 1000 nm

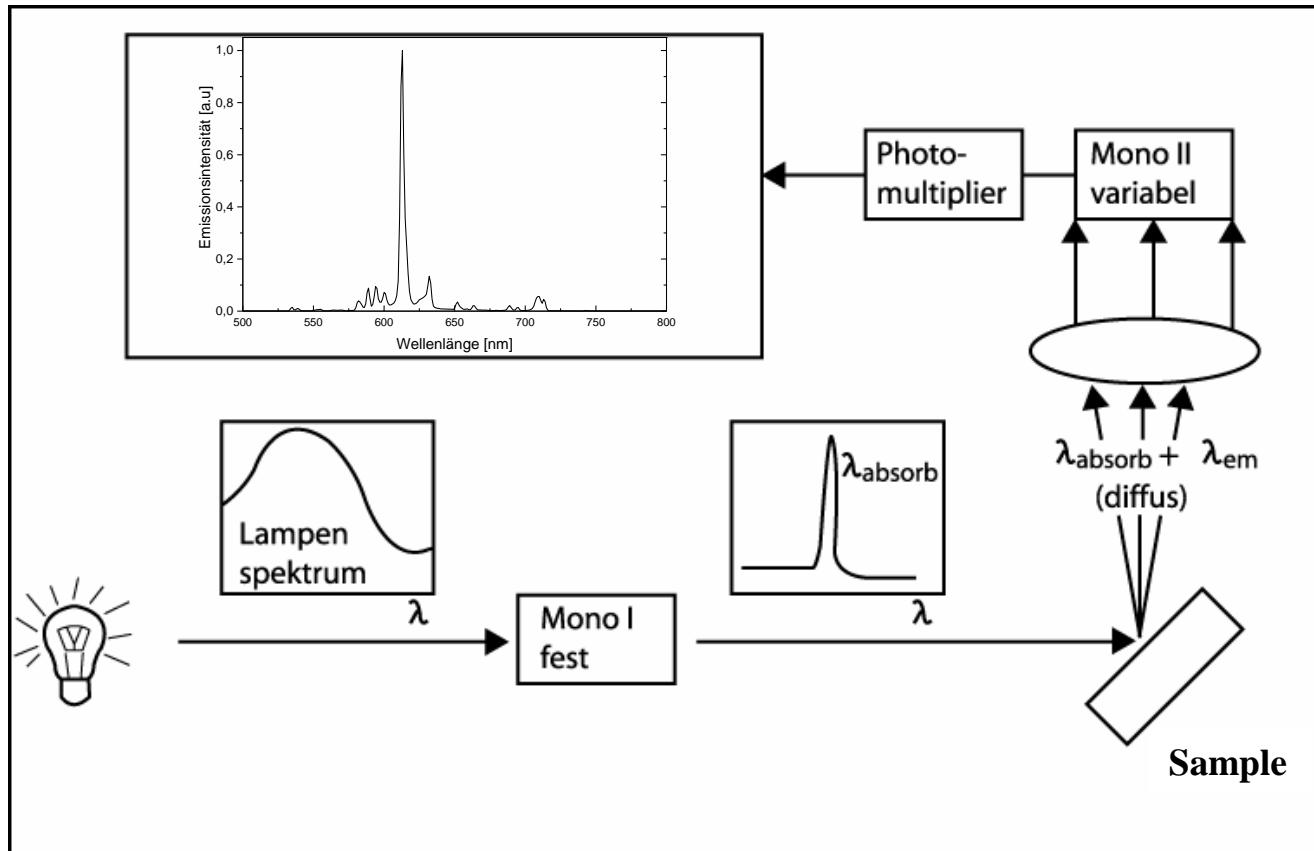
4.2.3 Emission Spectroscopy

Measurement of the intensity as a function of emission wavelength

Mono 1:
Constant, e.g. @ 254 nm

Mono 2:
variable, e.g. from 500
to 800 nm

Slit width determines
the maximum attainable
optical resolution

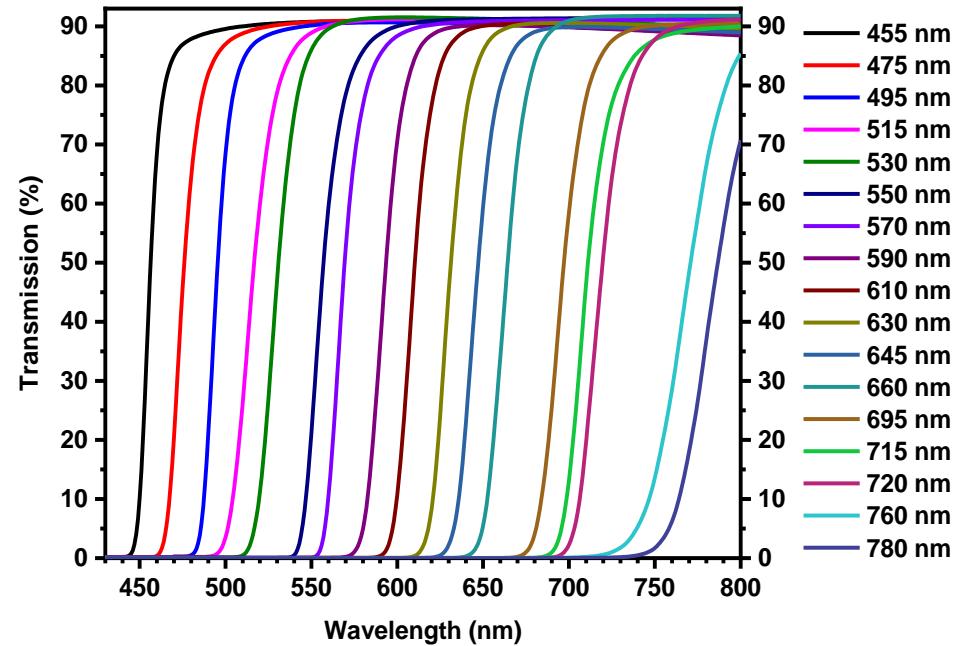
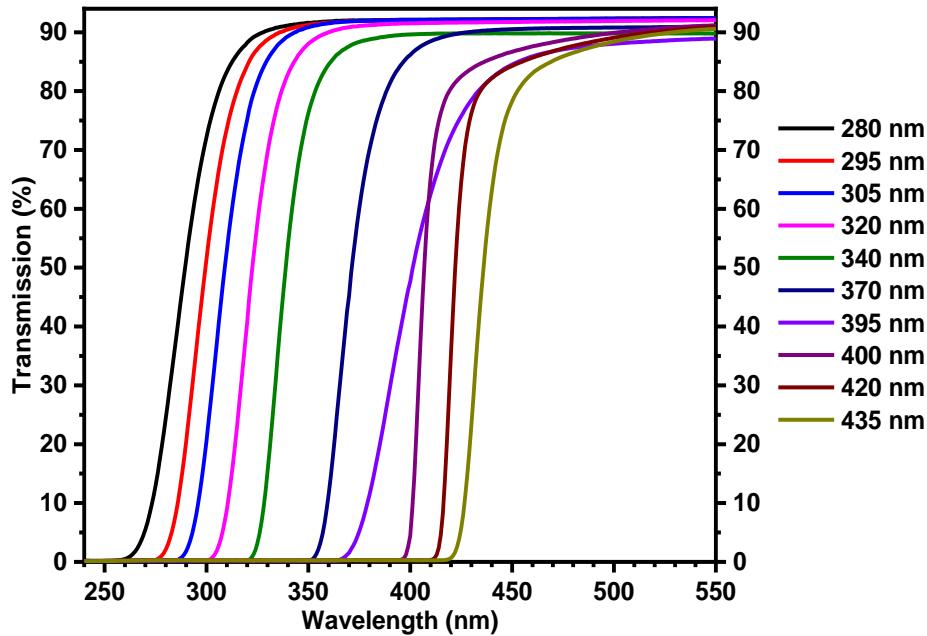


$I(\lambda)$ must be compensated for the fluctuations of the light source, for example, by a quantum counter, reflected monochromatic light of excitation beam must be filtered

4.2.3 Emission Spectroscopy

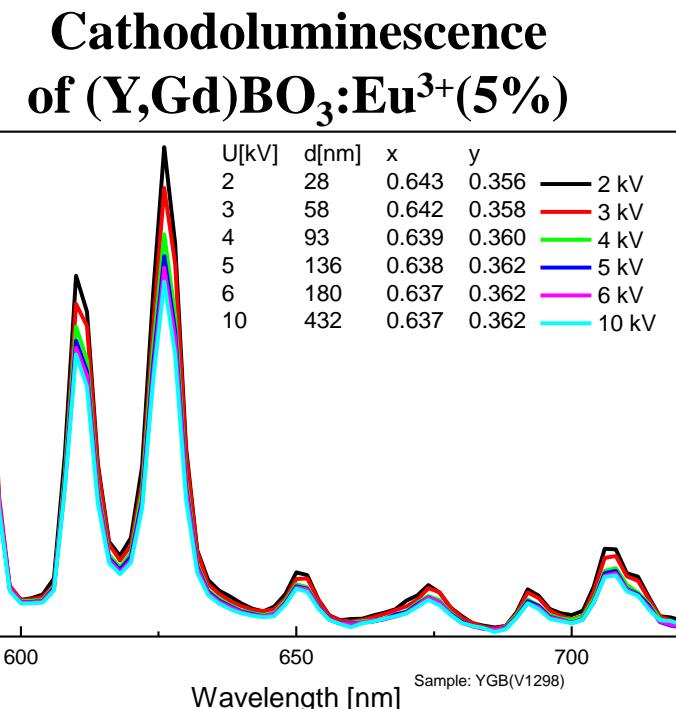
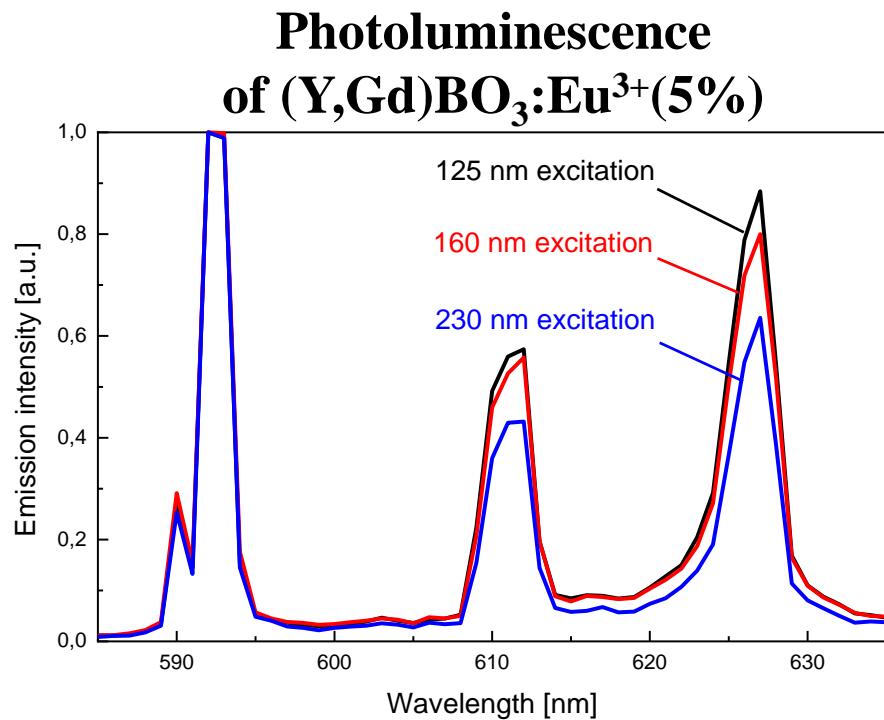
Measurement of the intensity as a function of emission wavelength

Removal of reflected monochromatic light of excitation beam by a long-pass filter



4.2.3 Emission Spectroscopy

The emission spectrum of a sample depends on the excitation energy



The excitation energy or wavelength determines the penetration depth of the radiation or of the electrons in the sample, whereby the emission spectrum might be modulated → change of CIE colour coordinates x, y

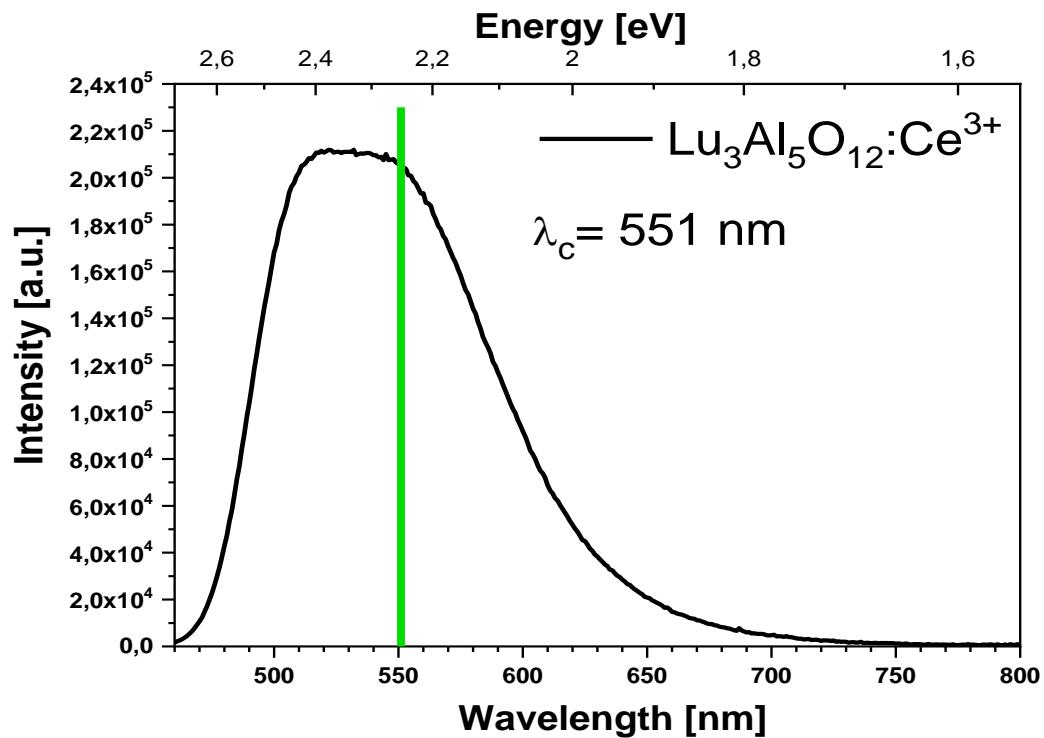
Penetration depth of electrons: $r \sim 0.046 * U^{5/3} / \rho$ [μm] → 2 kV: 30 nm, 10 kV: 400 nm

4.2.3 Emission Spectroscopy

From the emission spectrum the emission maxima is obtained and the centroid wavelength can be calculated

The centroid wavelength λ_c is the wavelength, at which the integral of an emission spectrum is subdivided into two equal fractions. It thus corresponds to the median photon energy of the emission spectrum.

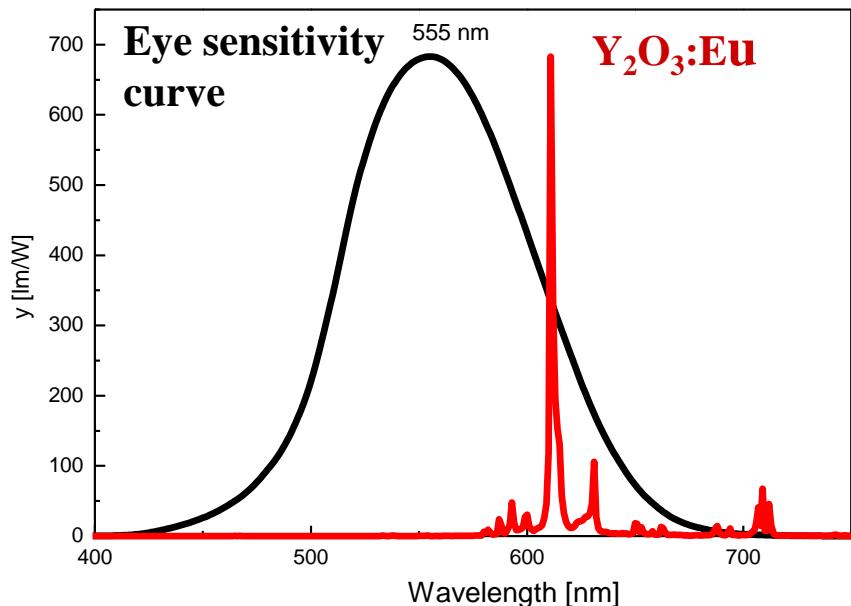
$$\lambda_c = \frac{\int_{\lambda_1}^{\lambda_2} \lambda \cdot S(\lambda) \cdot d\lambda}{\int_{\lambda_1}^{\lambda_2} S(\lambda) \cdot d\lambda}$$



4.2.3 Emission Spectroscopy

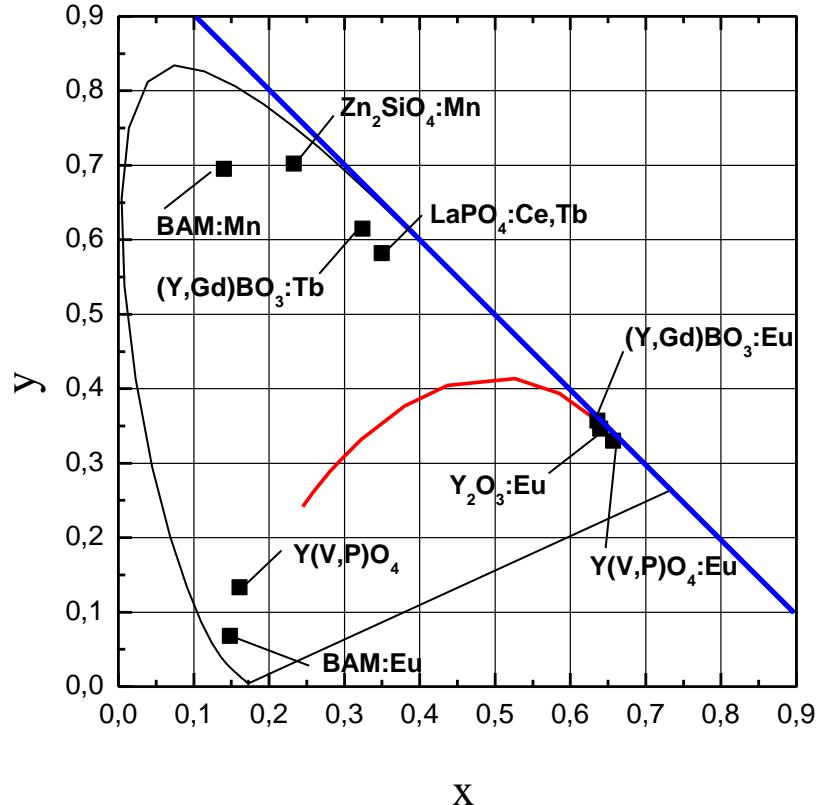
From the emission spectrum of light physiological figures can be calculated

Lumen equivalent LE [lm/W]



These quantities characterize inter alia
the emission spectrum of light sources and
phosphors

Chromaticity coordinate x, y (CIE 1931)
or coordinate u' , v' (CIE 1976)



4.2.3 Emission Spectroscopy

VUV Spectroscopy

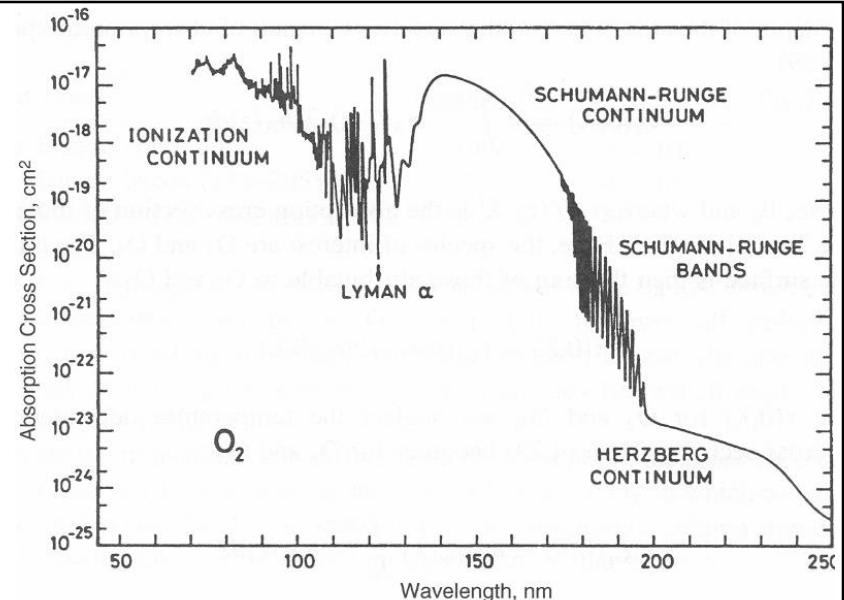
VUV spectroscopy requires a special instrumental setup, as VUV radiation is absorbed by O₂ and H₂O. CO₂ will be split photochemically into C and O₂ (dark layer on mirrors)

- Lamp (deuterium), monochromator and detector in high vacuum (10^{-6} mbar)
- The sample chamber can be flooded with N₂. This facilitates changing of the sample. The sample chamber is connected to the rest of the system via LiF or MgF₂ windows

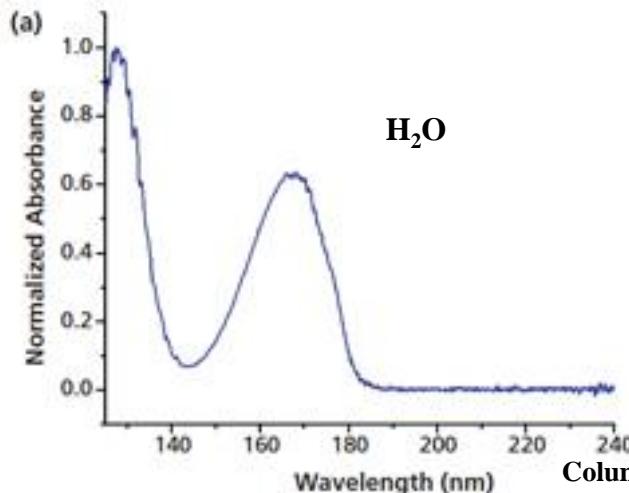
MgF₂: $E_g = 10.8$ eV (115 nm)

LiF: $E_g = 13 - 14$ eV (89 - 95 nm)

Problem: Fluorides degrade due to OH take-up



J. Seinfeld, S. Pandis, *Atmospheric Chemistry and Physics*



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4.2.3 Emission Spectroscopy

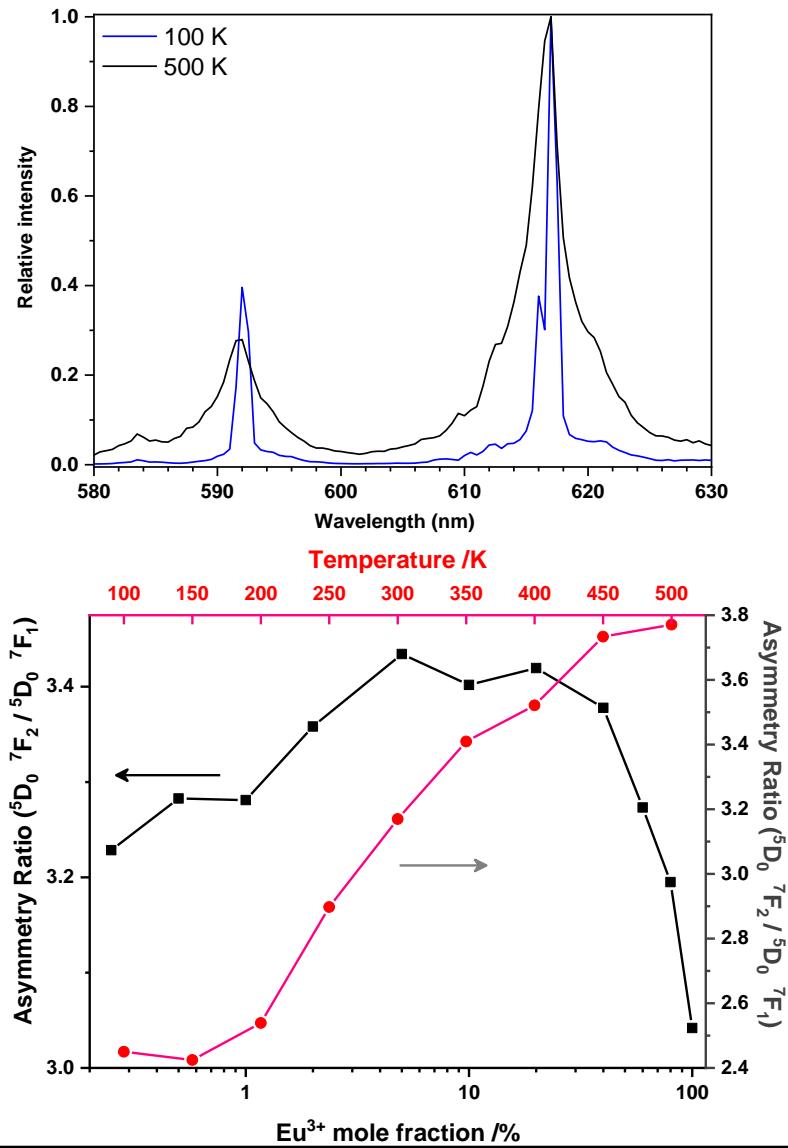
Impact of temperature and site symmetry

Temperature increase results in line broadening due to coupling to vibronic levels

Some Eu^{3+} related optical transitions, e. g. at about 615 nm, are the less probable, the closer the local site symmetry is to an inversion symmetry.

The intensity ratio of I_{615}/I_{590} is called **asymmetry ratio** – the larger the value, the less symmetrical the dopant site.

The asymmetry ratio can change depending on temperature, pressure or other parameters

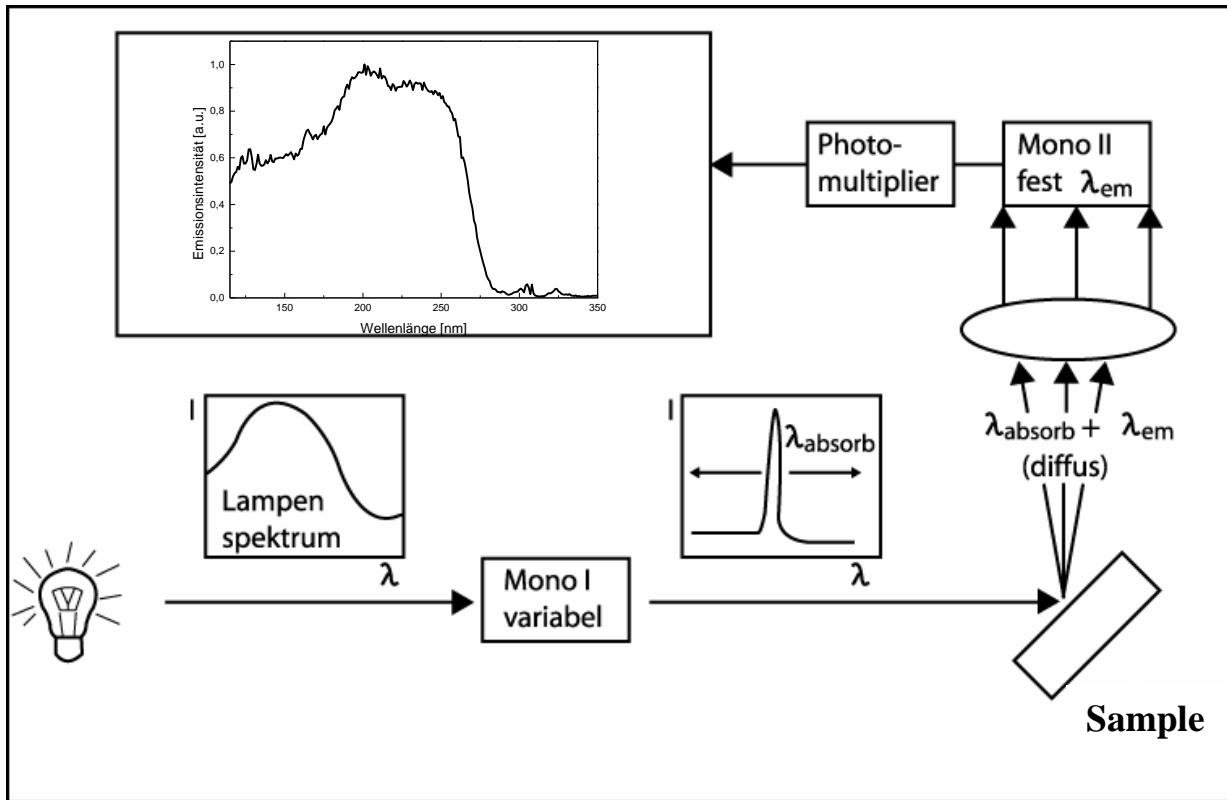


4.2.4 Excitation Spectroscopy

Measurement of the intensity as a function of excitation wavelength

Mono 1:
variable, e.g. from 120
to 400 nm

Mono 2:
constant, e.g. 611 nm
($\text{Y}_2\text{O}_3:\text{Eu}$)



Correction of the excitation spectrum for the spectrometer transfer function (set)
by use of rhodamine B (constant quantum yield below approximately 500 nm)

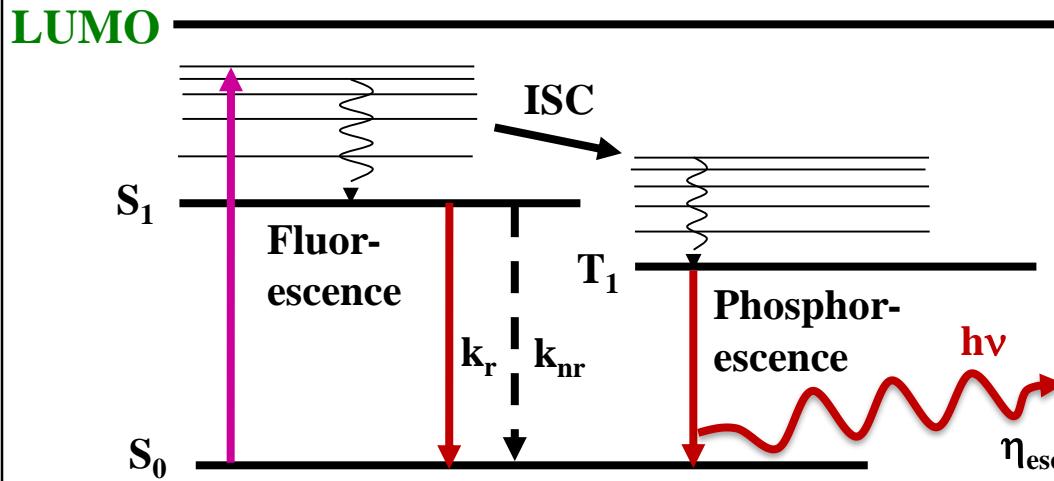
$$I(\lambda_{\text{exc}}) = I^{\text{sample}}(\lambda_{\text{exc}}) / I^{\text{set}}(\lambda_{\text{exc}})$$

4.2.5 Quantum Yield

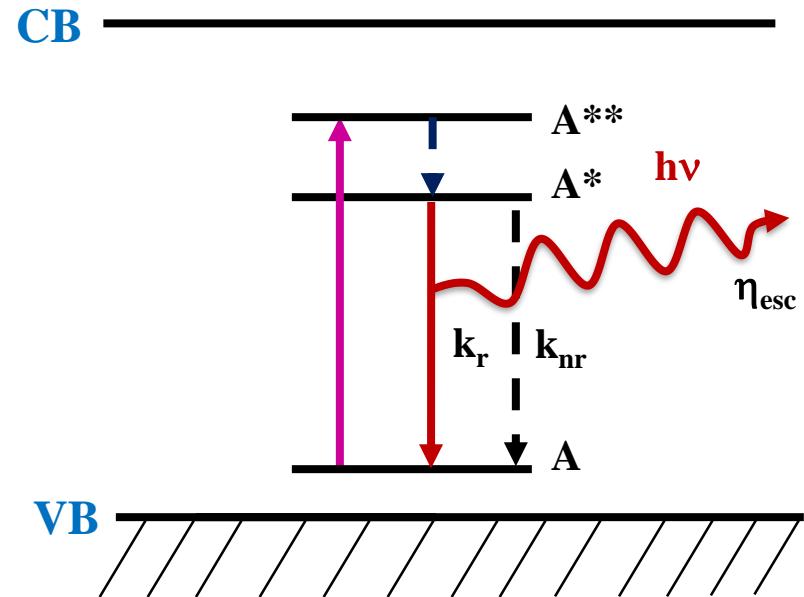
Absorption of radiation yields excited states, which then return to the ground state

Excitation energy $< E_G$ of solvent or host crystal $\sim \Delta E$ of optical centres

Organic materials (luminophores, dyes)



Inorganic materials (phosphors)



$$QY = \frac{\text{number of photons emitted}}{\text{number of photons absorbed}}$$

Such QY has to be corrected by
- absorption strength
- refractive indices

4.2.5 Quantum Yield

Definitions

- The PL quantum efficiency QY is defined as the ratio of the number of emitted photons by the number of absorbed photons, while this figure can be regarded as the external quantum yield (EQY), which is the macroscopic view on a material
- In practice, QY is often determined by comparative measurements with a reference compound for which it has been determined with a high degree of accuracy, e.g. from US NIST or German PTB
- Ideally, a reference compound should have
 - the same absorbance (A) as the compound of interest at given excitation wavelength
 - similar excitation-emission characteristics to compound of interest, otherwise instrument's transfer function should be taken into account
 - Same solvent/matrix, because intensity of emitted light is dependent on refractive index (n), otherwise a correction has to be applied
 - Yields similar fluorescence intensity to ensure measurements are taken within the range of linear instrument response

$$\frac{QY^u}{QY^s} = \frac{I^u}{I^s} \times \frac{n^2(u)}{n^2(s)}$$

4.2.5 Quantum Yield

Relative measurements in comparison to references

Typical reference materials (determined by PTB, Germany)

Spectral region	Material	QY at 253.7 nm (Hg line)
UV-B	LaPO ₄ :Ce	90%
UV-A	BaSi ₂ O ₅ :Pb	88%
Blue	BaMgAl ₁₀ O ₁₇ :Eu	89%
Green	Zn ₂ SiO ₄ :Mn	80%
Red	Y ₂ O ₃ :Eu	84%

For the determination of the quantum yield one determines the integral I of the emission spectrum (proportional to the number of emitted photons) and the reflectance R of the substance or the reference at the excitation wavelength λ_{exc}

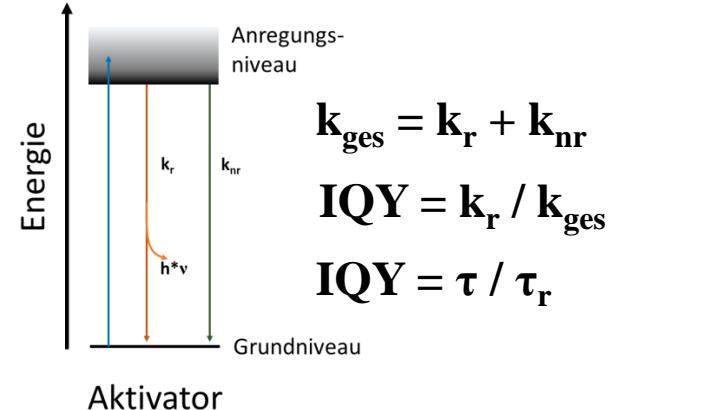
$$\Phi_{\text{sample}, \lambda_{\text{exc}}} = \Phi_{\text{reference}, \lambda_{\text{exc}}} \cdot \frac{\int I_{\text{sample}} d\lambda - \int I_{\text{black}} d\lambda}{\int I_{\text{reference}} d\lambda - \int I_{\text{black}} d\lambda} \cdot \frac{1 - R_{\text{reference}, \lambda_{\text{exc}}}}{1 - R_{\text{sample}, \lambda_{\text{exc}}}}$$

4.2.5 Quantum Yield

Absolute methods for the determination of QY

- Internal quantum yield:

decay time



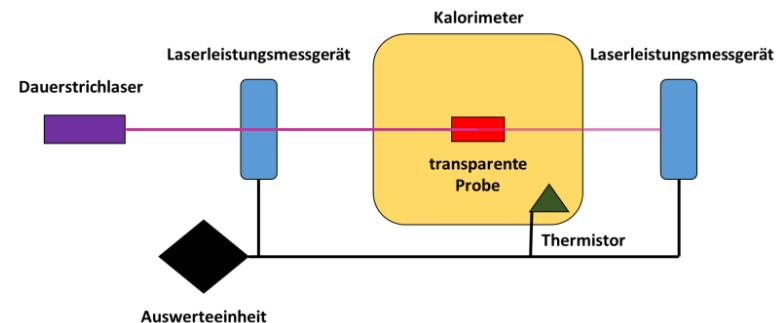
$$k_{\text{ges}} = k_r + k_{\text{nr}}$$

$$\text{IQY} = k_r / k_{\text{ges}}$$

$$\text{IQY} = \tau / \tau_r$$

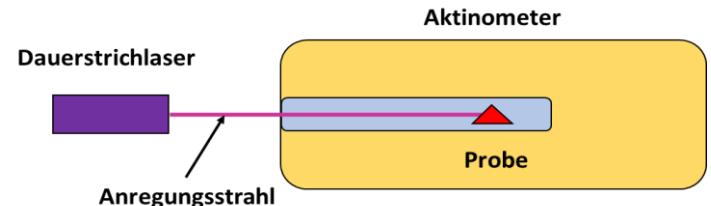
- External quantum yield:

calorimetric



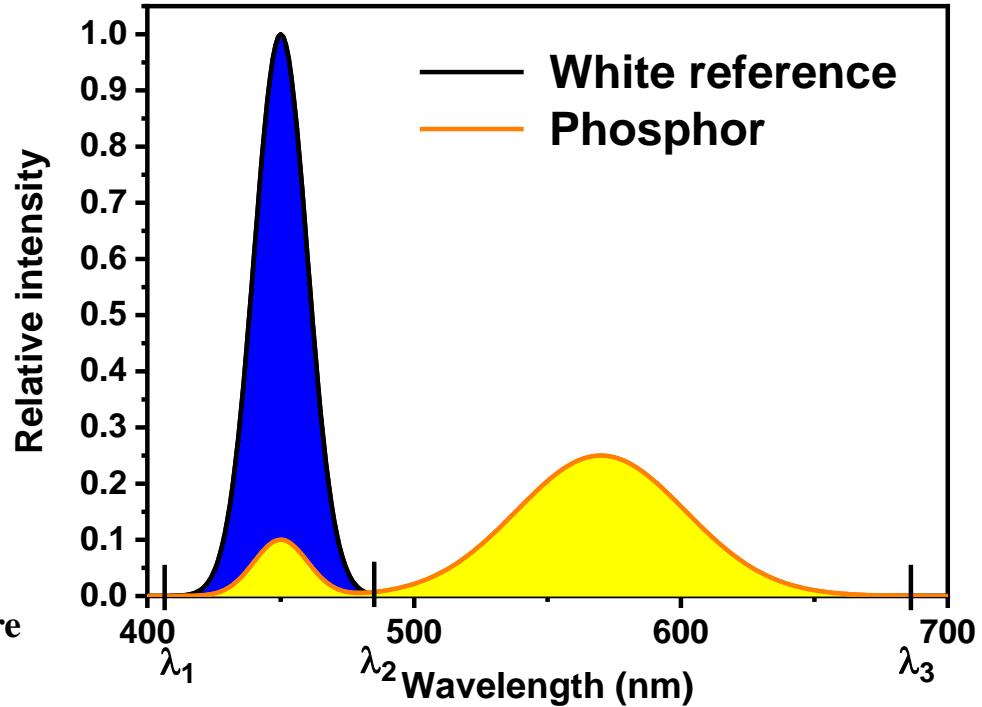
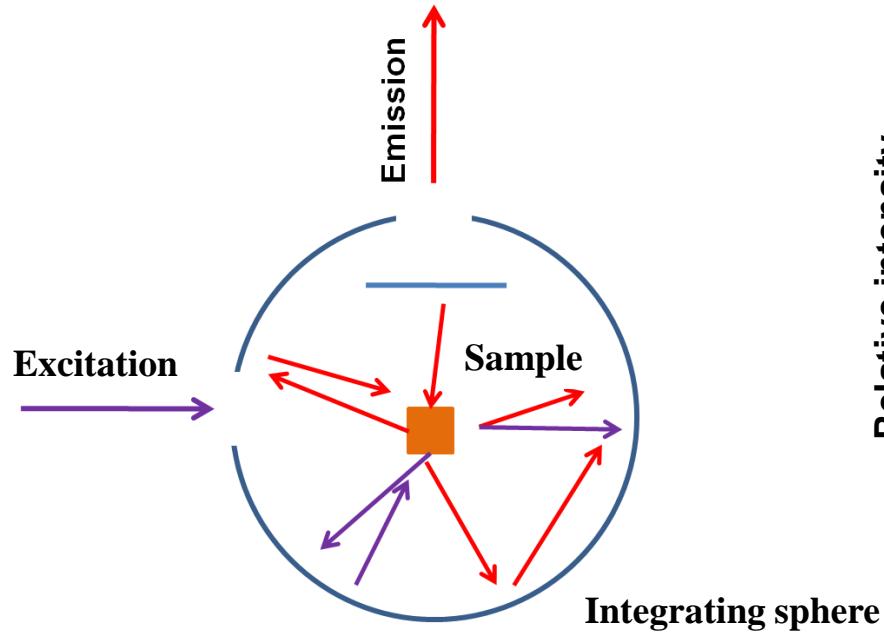
- External quantum yield:

actinometric



4.2.5 Quantum Yield

External quantum yield: absolute methods by using an integration sphere
 (Lit.: Jap. J. Applied Physics 43 (2004) 7729)



Material	λ_{PL} (nm)	η_{PL}	τ_{PL} ($\times 10^{-9}$ s)	k_r (s $^{-1}$)
Alq ₃	520	0.20 ± 0.01	15	1.3×10^6
BAlq ₂	492	0.42 ± 0.02	29	1.4×10^6
TPD ⁹⁾	424	0.41 ± 0.02	0.6	6.8×10^8
α -NPD ⁹⁾	445	0.29 ± 0.02	3.5	8.0×10^7
CBP ⁹⁾	393	0.61 ± 0.04	0.7	8.7×10^8
Ir(ppy) ₃ ^{a)}	512	0.97 ± 0.02	1.3×10^3	7.5×10^5

$$EQY_{sample} = \frac{n_{emission}}{n_{absorption}} = \frac{\int_{\lambda_2}^{\lambda_3} \frac{\lambda}{hc} [I_{sample}^{emission}(\lambda) - I_{reference}^{emission}(\lambda)] d\lambda}{\int_{\lambda_1}^{\lambda_2} \frac{\lambda}{hc} [I_{reference}^{excitation}(\lambda) - I_{sample}^{excitation}(\lambda)] d\lambda}$$

4.2.5 Quantum Yield

Definition by the lifetime of the excited state of the luminescent center

- Another definition for QY is given by wherein k_r is the radiative rate constant and Σk is the sum of the rate constants for all processes that depopulate the excited state
- Radiative lifetime τ_r is related to k_r
- Measured lifetime τ depends on the sum of the rate constants
- The observed fluorescence lifetime τ , is the average time the molecule spends in the excited state
⇒ Measured lifetime is proportional to quantum yield

$$QY = \frac{k_r}{\sum_i k_i}$$

$$\tau_r = \frac{1}{k_r}$$

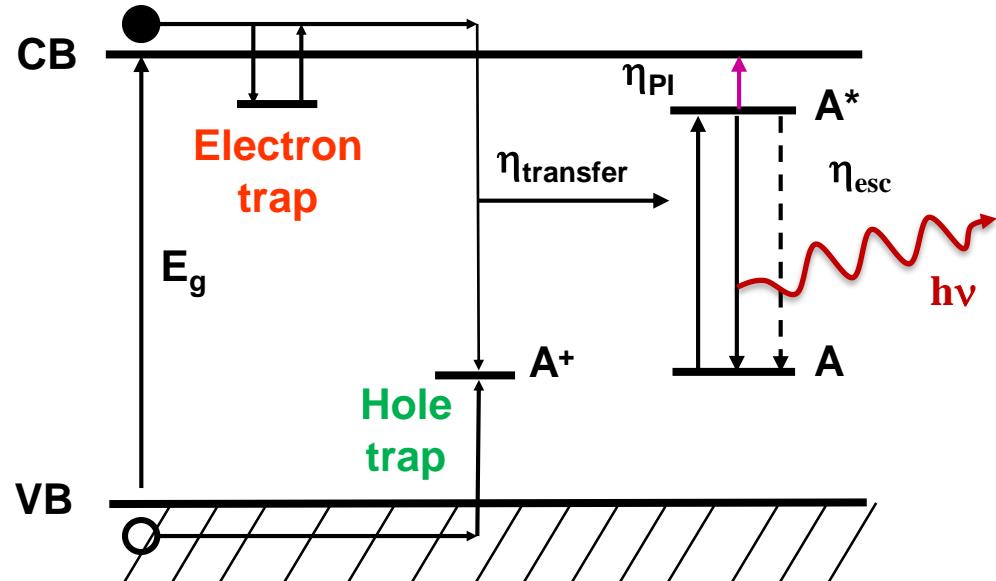
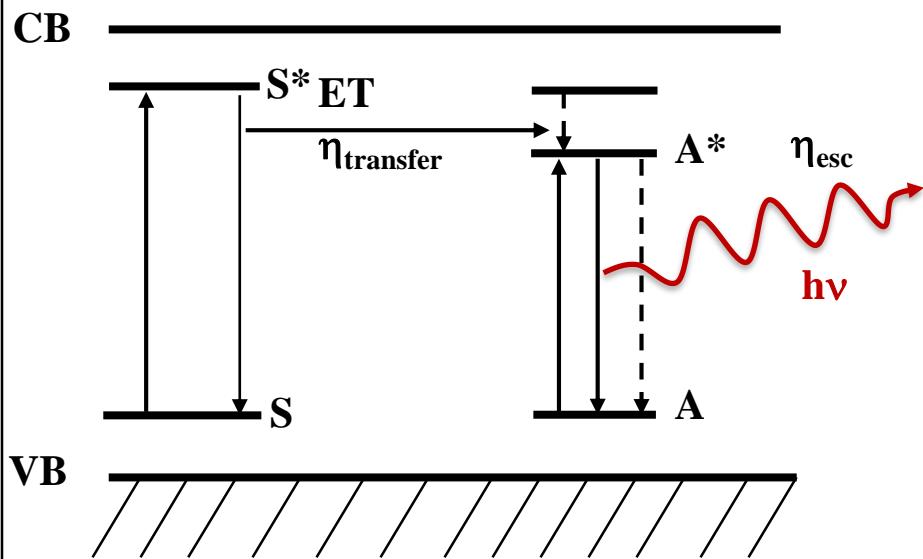
$$\tau = \frac{1}{\sum_i k_i}$$

$$QY = \frac{\tau}{\tau_r}$$

4.2.5 Quantum Yield

Excitation by VUV or EUV radiation

Excitation energy $\sim \Delta E$ of sensitizers or E_g of the host matrix/solvent



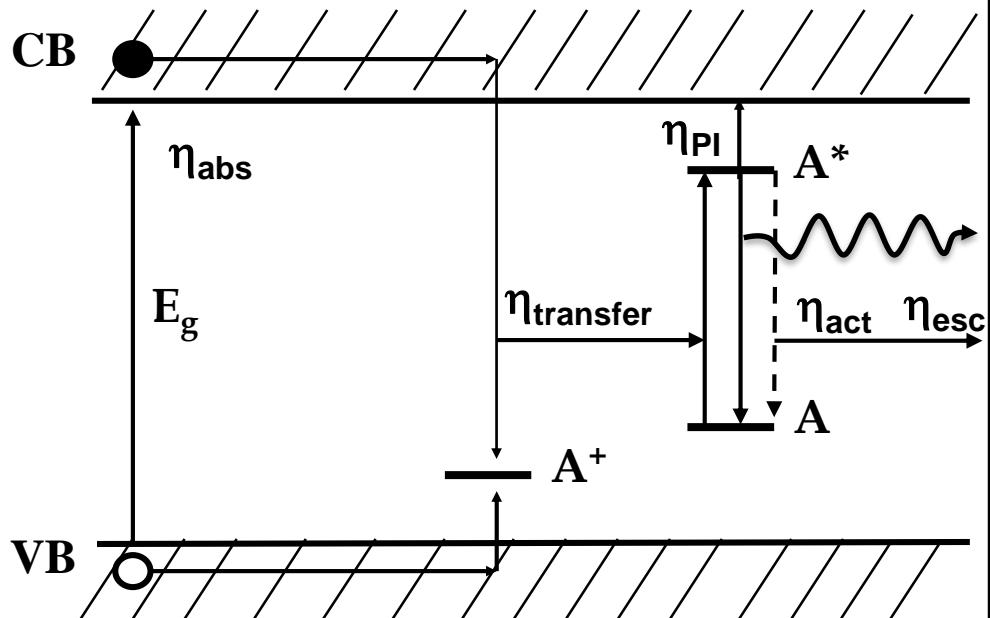
$$\eta_{\text{total}} = \eta_{\text{act}} * \eta_{\text{transfer}} * \eta_{\text{esc}} * (1 - \eta_{\text{PI}}) \text{ with PI = Photoionisation}$$

η_{total} = External Quantum Yield (EQY) or External Quantum Efficiency (EQE)

η_{act} = Internal Quantum Yield (IQY) or Internal Quantum Efficiency (IQE)

4.2.5 Quantum Yield

Quenching is caused by the luminescent center, e.g. the activator ion, ($\text{IQY} \downarrow$) and/or the host matrix ($\text{EQY} \downarrow$)



Internal Quantum Yield (Efficiency)

$$\begin{aligned}\text{IQY} &= \eta_{\text{act}} \\ &= k_r / (k_r + k_{\text{nr}}) \\ &= \tau / \tau_r\end{aligned}$$

with $1 / (k_r + k_{\text{nr}}) = \tau$ and $k_r = 1 / \tau_r$
(anti proportional to decay time)

Decay time (Fluorescence spectrometer)

External Quantum Yield (Efficiency)

$$\begin{aligned}\text{EQY} &= N_{hv(\text{emitted})} / N_{hv(\text{absorbed})} \\ &= \eta_{\text{transfer}} * \eta_{\text{act}} * \eta_{\text{esc}} * (1 - \eta_{\text{PI}})\end{aligned}$$

(No correlation to decay time!)

Emission spectrum (Ulbricht sphere)

Light Yield (Fluorescence Intensity)

$$\text{LY} = \text{EQY} * \eta_{\text{abs}}$$

(No correlation to decay time!)

Absorption via reflection spectrum and Kubelka-Munk-Function (Ulbricht sphere)

4.2.5 Quantum Yield

Comparison: Internal Quantum Yield (IQY) vs. External Quantum Yield (EQY)

Internal Quantum Yield (Efficiency)

Efficiency of a luminescent center

Microscopic view

Decay time measurement

Fluorescence spectrometer

$$IQY = \frac{\tau}{\tau_r}$$

External Quantum Yield (Efficiency)

Efficiency of a luminescent powder or solution

Macroscopic view

Emission spectrum measurement

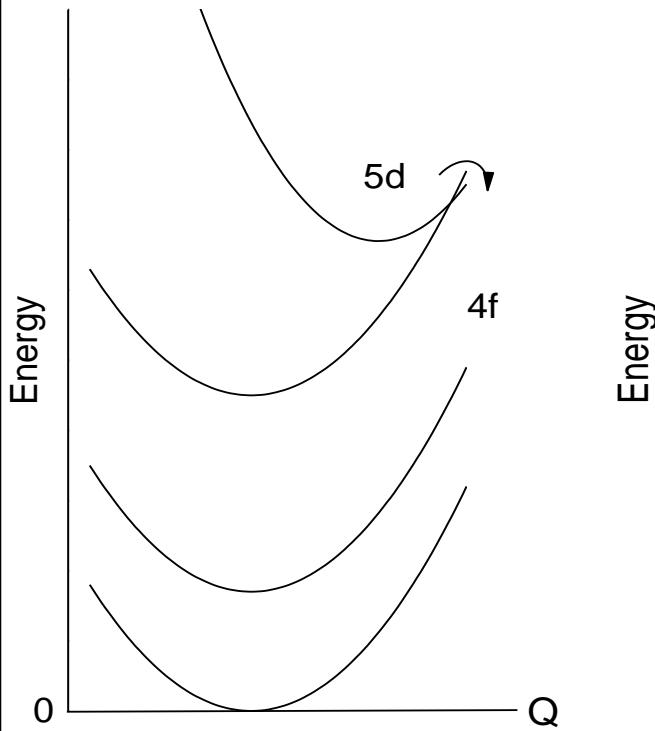
Ulbricht (integrating) sphere

$$EQY = \frac{\text{number of photons emitted}}{\text{number of photons absorbed}}$$

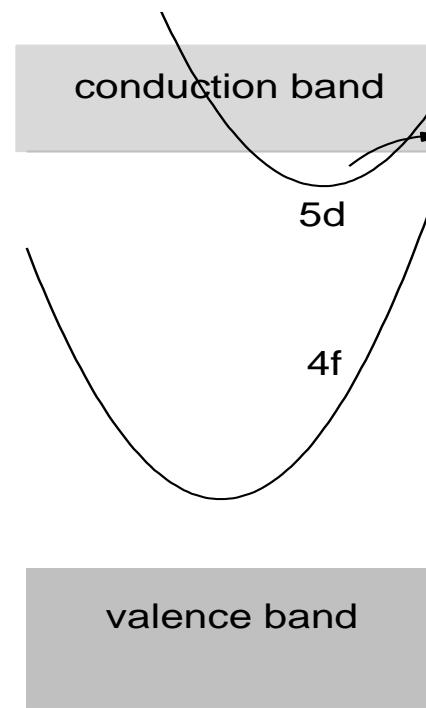
with τ = experimental lifetime and τ_r = radiative lifetime

4.2.5 Quantum Yield

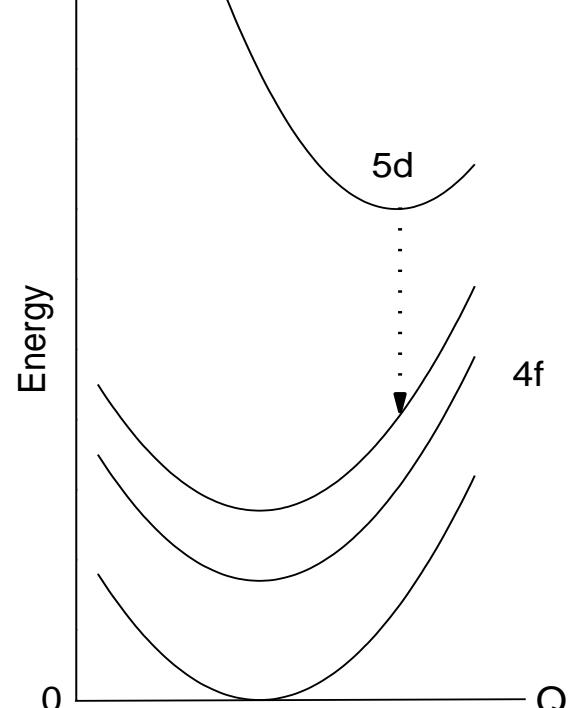
Several mechanisms for the quenching of luminescent materials results in the reduction of the photoluminescence quantum yield, e.g.



**Thermally activated
InterSystem Crossing (ISC)**



**Thermally activated
PhotoIonization (PI)**

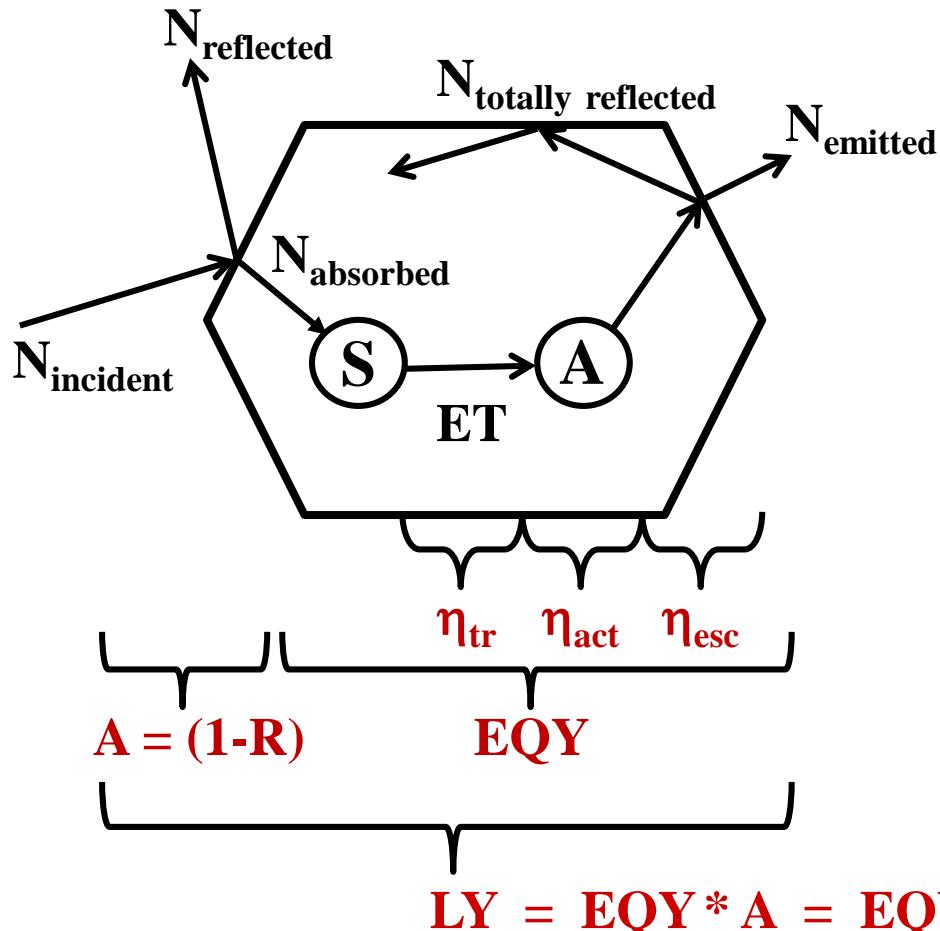


**Multi Phonon Relaxation
(MPR)**

4.2.6 Light Yield

Light Yield LY = Fluorescence Intensity ~ $I(\lambda_{\text{exc}})$ (or integral lamp output)

Light conversion at a phosphor particle



External Quantum Yield (EQY)

$$\text{EQY} = N_{\text{emitted}} / N_{\text{absorbed}}$$

Absorption A

$$A = N_{\text{absorbed}} / N_{\text{incident}}$$

Reflectance R

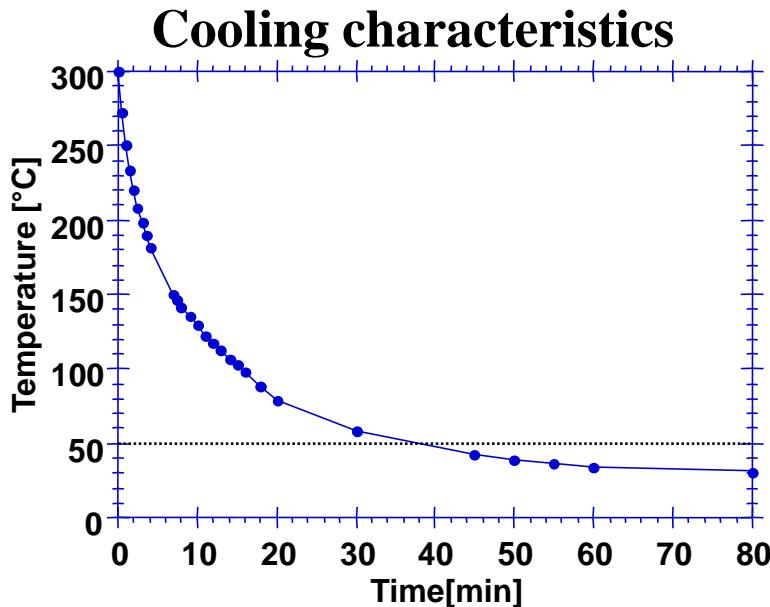
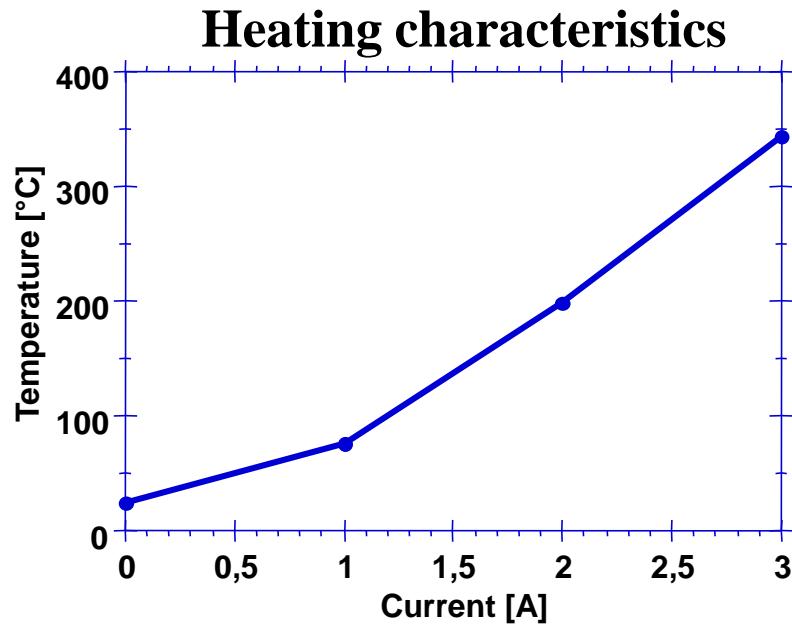
$$R = 1 - A \quad (\text{for transmission } T = 0)$$

Light Yield LY

$$\text{LY} = \frac{N_{\text{emitted}}}{N_{\text{incident}}} = \frac{N_{\text{emitted}}}{N_{\text{absorbed}}} * \frac{N_{\text{absorbed}}}{N_{\text{incident}}}$$

4.2.7 Temperature Resolved Spectroscopy

Heating and cooling behavior of a sample holder for temp. resolved spectroscopy



Heating the sample

N_2 -flushing cools the sample

⇒ Actual temperature of the sample is
5-10 °C below the measured value

Cooling the sample

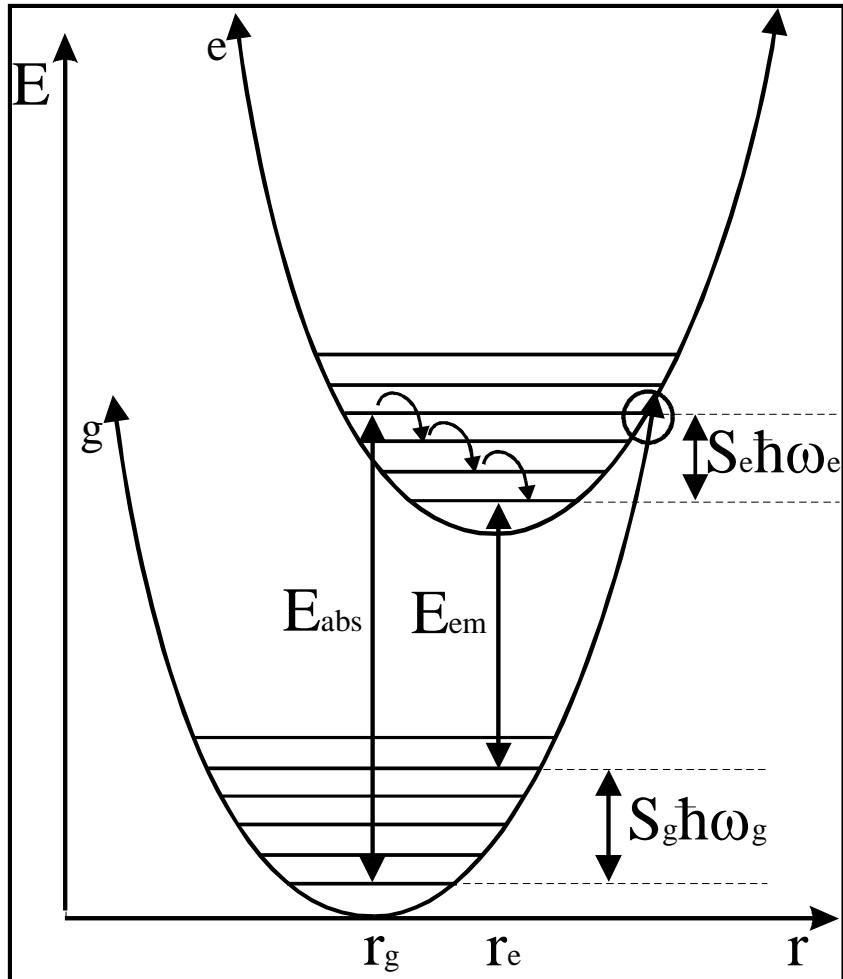
< 0 °C: H_2O Resublimation

< -78 °C: CO_2 Resublimation

⇒ Additional absorption in
the VUV range

4.2.7 Temperature Resolved Spectroscopy

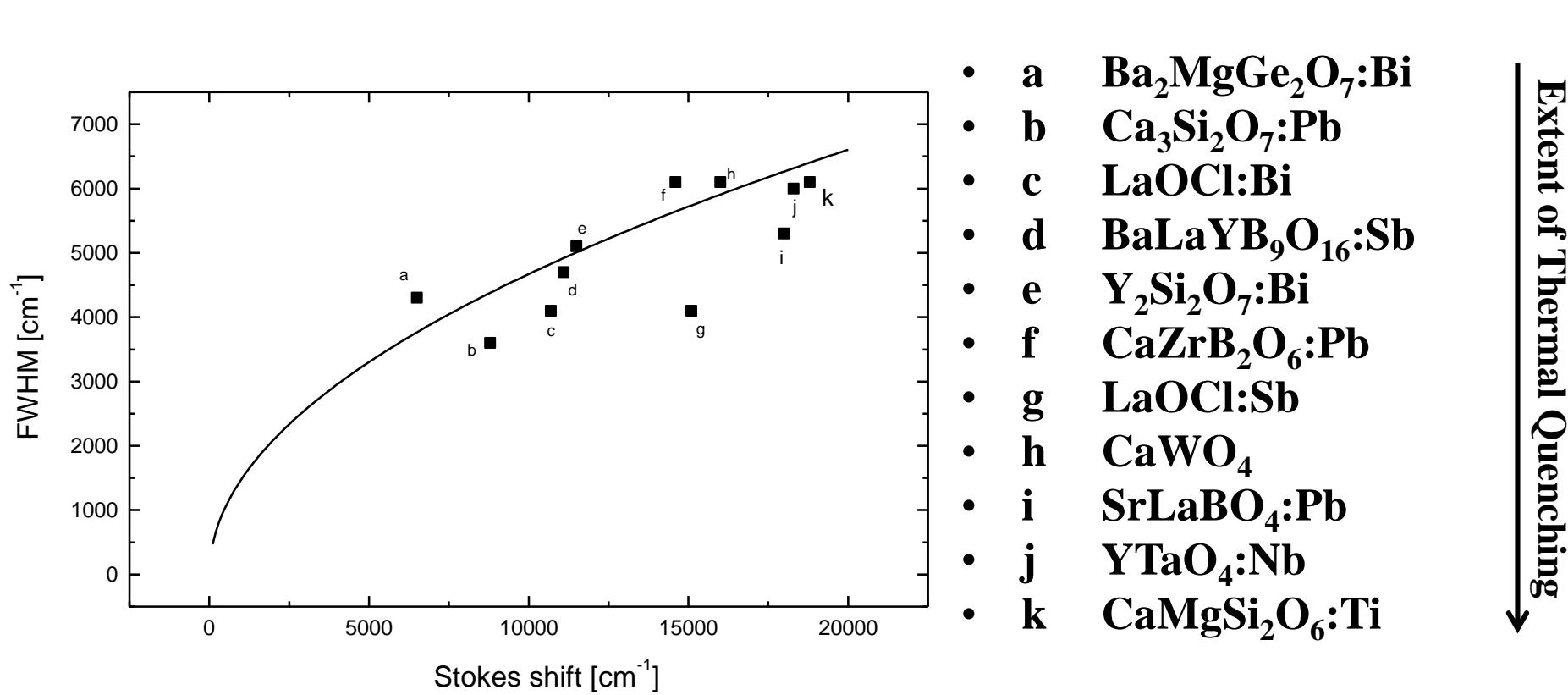
Thermal quenching of luminescence processes



- Stokes shift
Energy gap between absorption and emission band
$$S = S_e \hbar \omega_e + S_g \hbar \omega_g$$
- Full width at half maximum of the emission band
$$\text{FWHM} \sim \sqrt{S}$$
- Thermal quenching decreases with increasing $\Delta R = r_e - r_g$
- ΔR depends on the activator-host lattice interaction

4.2.7 Temperature Resolved Spectroscopy

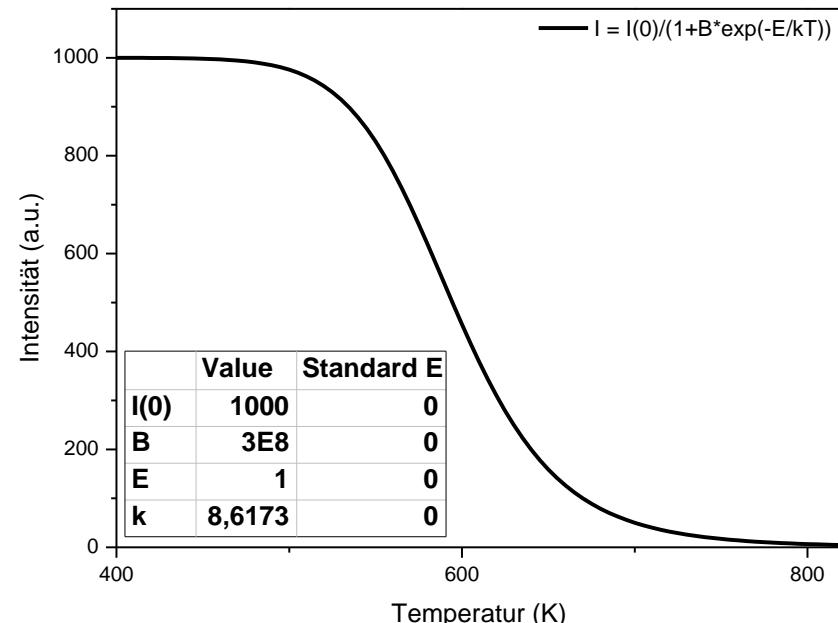
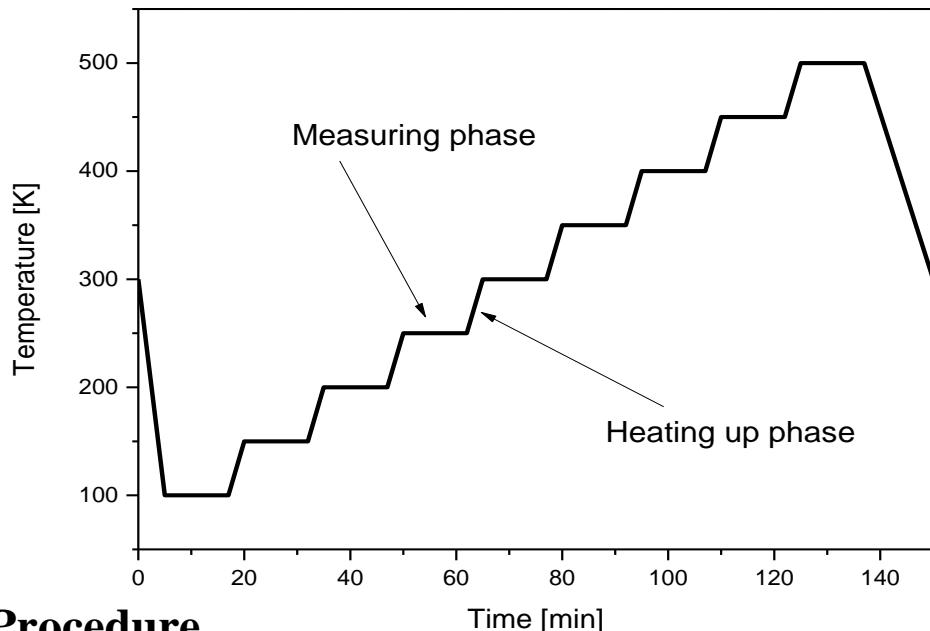
Thermal quenching of luminescence processes: FWHM = f(Stokes Shift)



$$\text{FWHM} \sim \sqrt{S}$$

4.2.7 Temperature Resolved Spectroscopy

Thermal quenching experiments \Rightarrow Intensity/integral as f(T)



Procedure

- Heating up to temperature T (step wisely)
- Measurement of the emission spectrum at λ_{max} upon excitation of the sample at $\lambda_{\text{exc.}}$
- Plotting of the integral and/or of the peak intensity as function of T
- Fitting of the obtained quenching curve by the following function

$$I(T) = A_0 + I_0/(1 + B \exp(-\Delta E/k_B T)) \quad \text{Fitting by the „Struck-Fonger-Model“}$$

4.2.7 Temperature Resolved Spectroscopy

Curve fitting

$N_{NR}(T)$ is the non-radiative transition rate at temperature T (s⁻¹)

$N_{NR,max}$ is the maximal non-radiative transition rate (s⁻¹)

E is the activation energy for thermal quenching (eV, cm⁻¹, ...)

k is the Boltzmann constant (eV, cm⁻¹, ...)

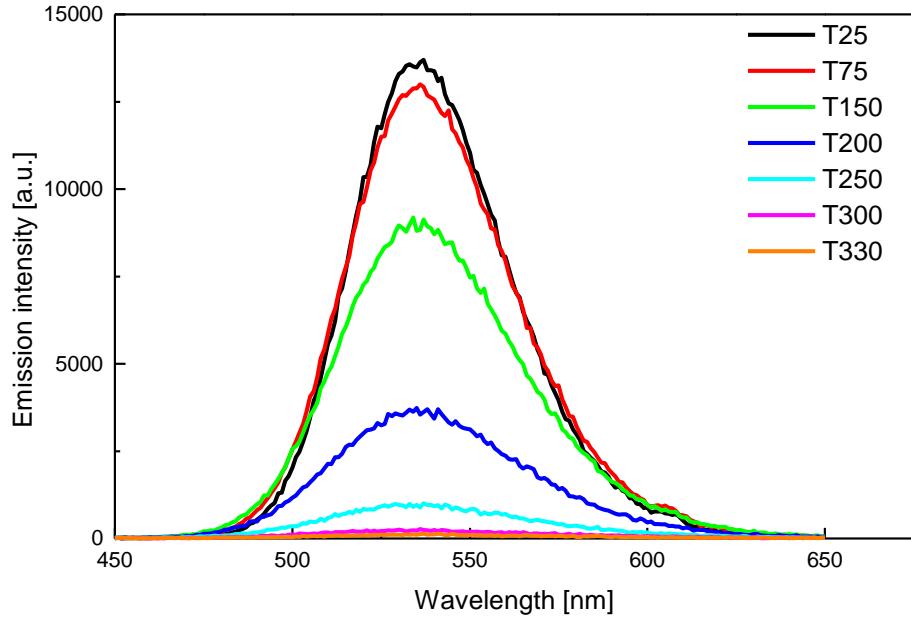
$$N_{NR} = N_{NR,max} * e^{-\frac{E}{kT}}$$

$$IQE = \frac{N_R}{N_R + N_{NR}} = \frac{1}{1 + \frac{N_{NR}}{N_R}}$$

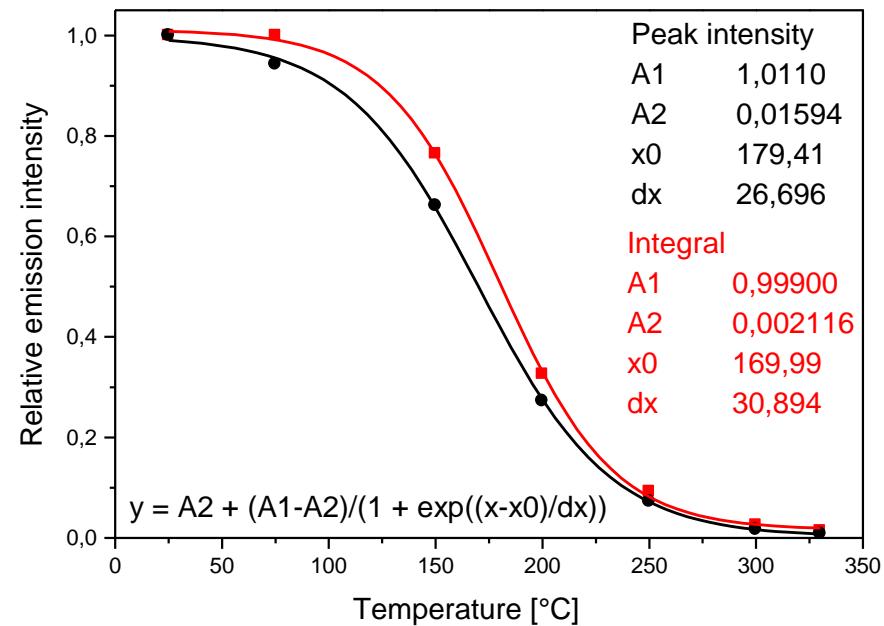
4.2.7 Temperature Resolved Spectroscopy

Example of SrGa₂S₄:Eu

Emission spectrum (460 nm excitation)



Emission intensity

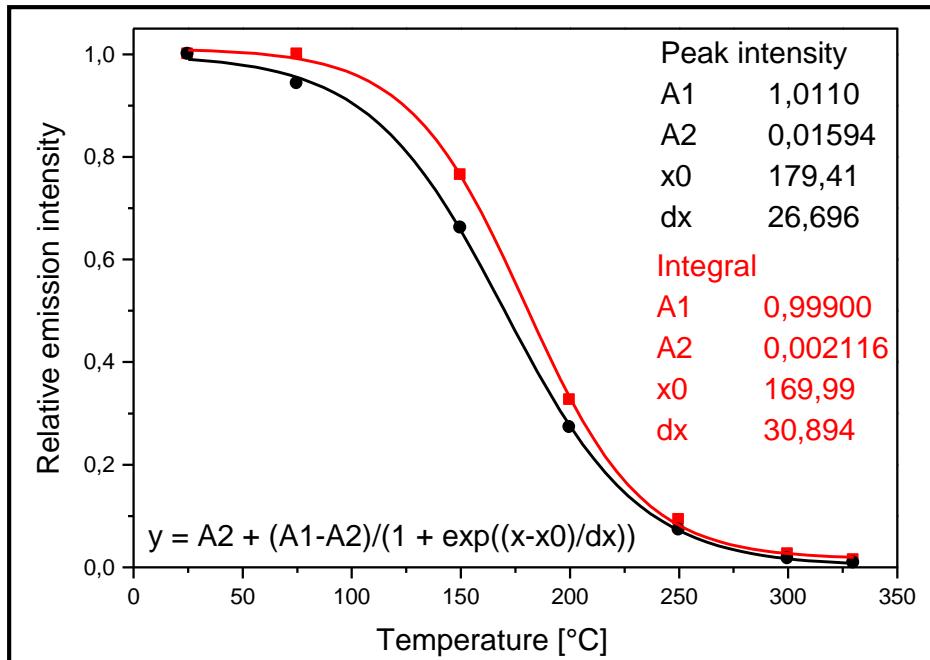


- $T_{1/2}$ = Temperature at which the luminescence process declines by 50% of its low temperature emission intensity ($\text{SrGa}_2\text{S}_4:\text{Eu} \sim 170^\circ\text{C}$)
- In many applied inorganic phosphors the quantum yield starts increasing distinctly from $100 - 150^\circ\text{C}$ upwards

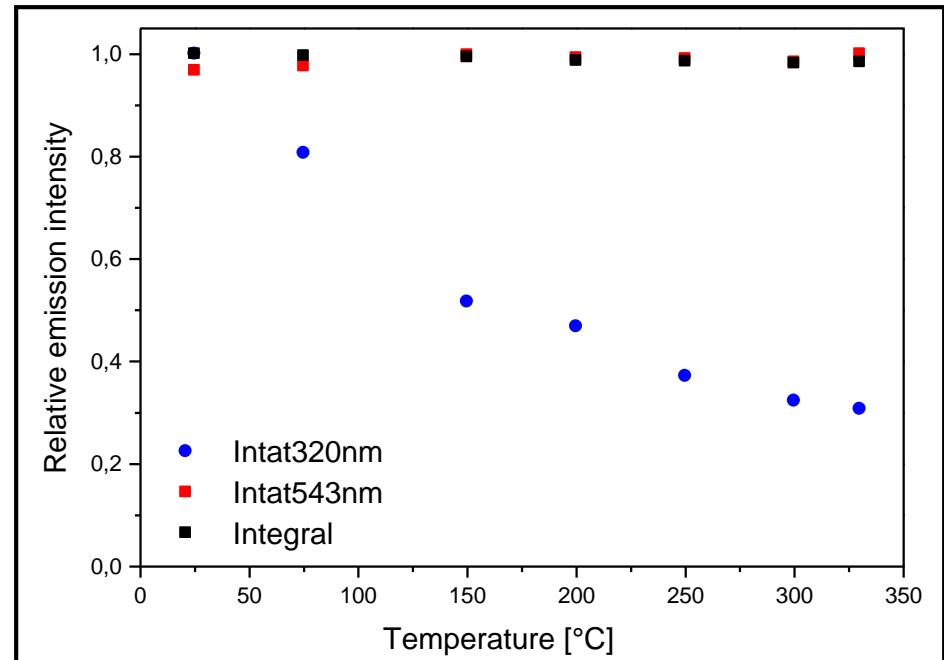
4.2.7 Temperature Resolved Spectroscopy

Comparison between a line and a band emitter

SrGa₂S₄:Eu (460 nm excitation)



LaPO₄:Ce,Tb (254 nm excitation)



Temperature resolved spectra for the investigation of

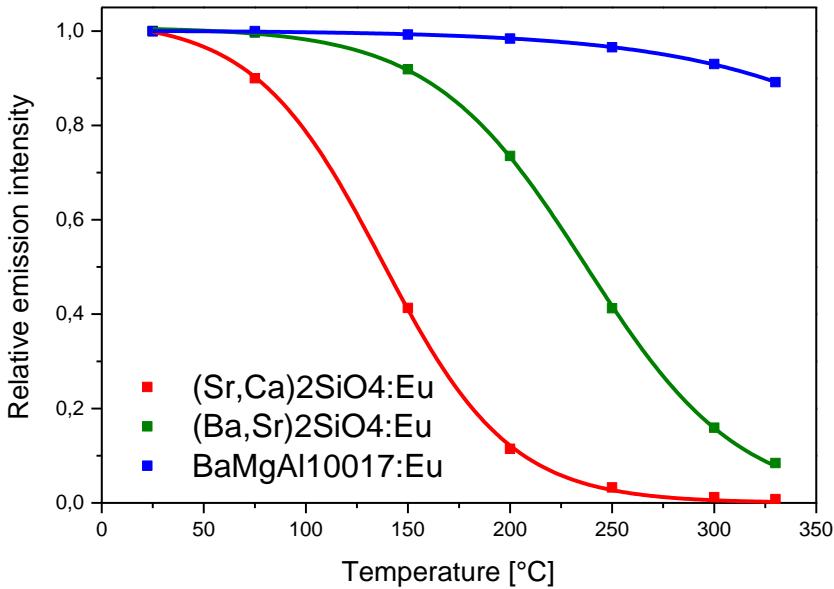
- Thermal quenching of the luminescence
- Energy transfer, e.g. in LaPO₄:Ce,Tb



4.2.7 Temperature Resolved Spectroscopy

Thermal quenching of Eu²⁺ phosphors

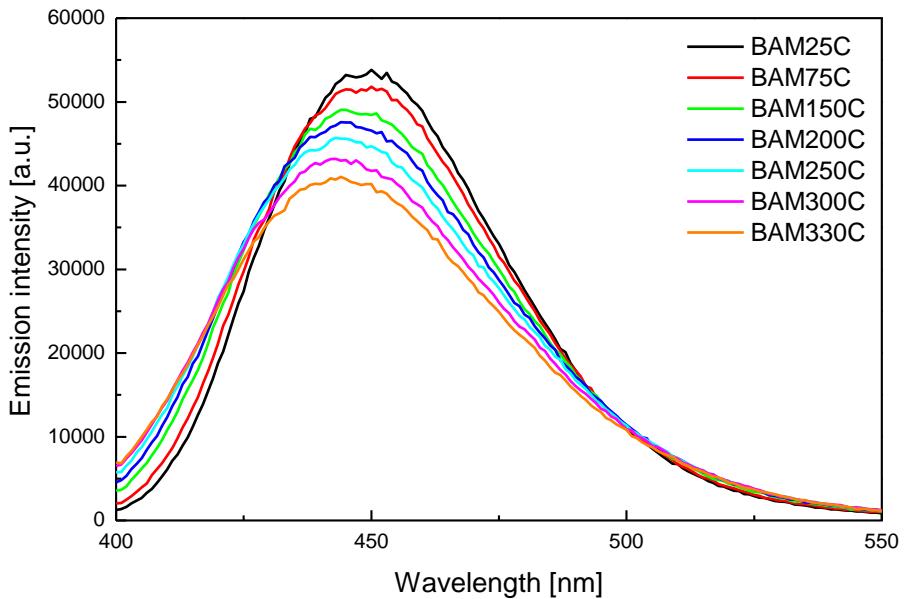
Quantum yield as a function of temperature



Stokes shift

Thermal quenching

Spectral width of the emission band of BaMgAl₁₀O₁₇:Eu as a function of T

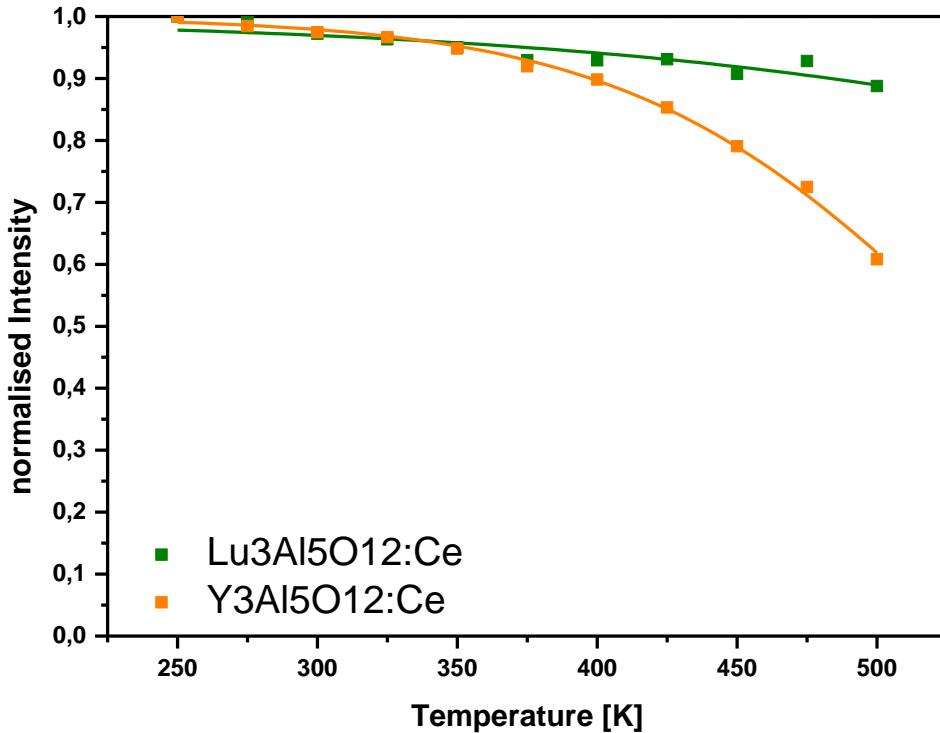
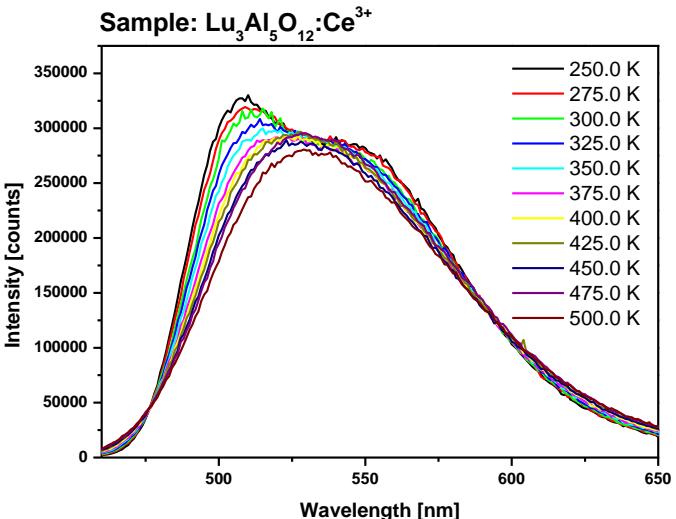
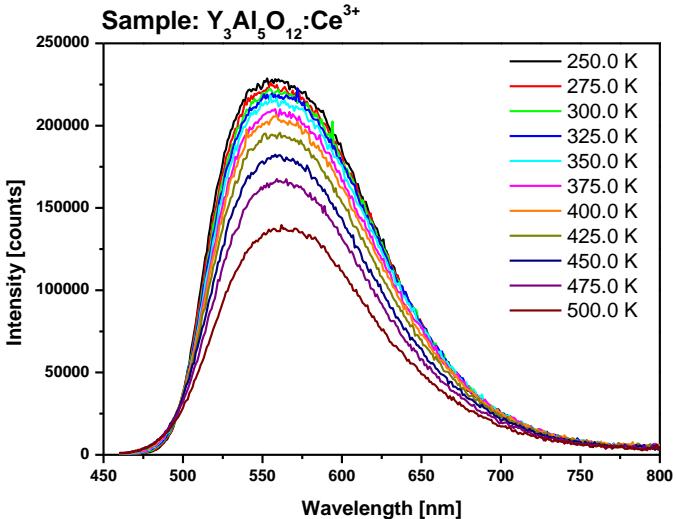


$\text{BaMgAl}_{10}\text{O}_{17}:\text{Eu} < (\text{Ba},\text{Sr})_2\text{SiO}_4:\text{Eu} < (\text{Sr},\text{Ca})_2\text{SiO}_4:\text{Eu}$
 $\text{BaMgAl}_{10}\text{O}_{17}:\text{Eu} < (\text{Ba},\text{Sr})_2\text{SiO}_4:\text{Eu} < (\text{Sr},\text{Ca})_2\text{SiO}_4:\text{Eu}$

Blue shift due to thermal expansion of the host and thus reduction in crystal field splitting!

4.2.7 Temperature Resolved Spectroscopy

Thermal quenching of Ce³⁺ phosphors



Stokes shift:

Thermal quenching:

Lu₃Al₅O₁₂:Ce < Y₃Al₅O₁₂:Ce

Lu₃Al₅O₁₂:Ce < Y₃Al₅O₁₂:Ce

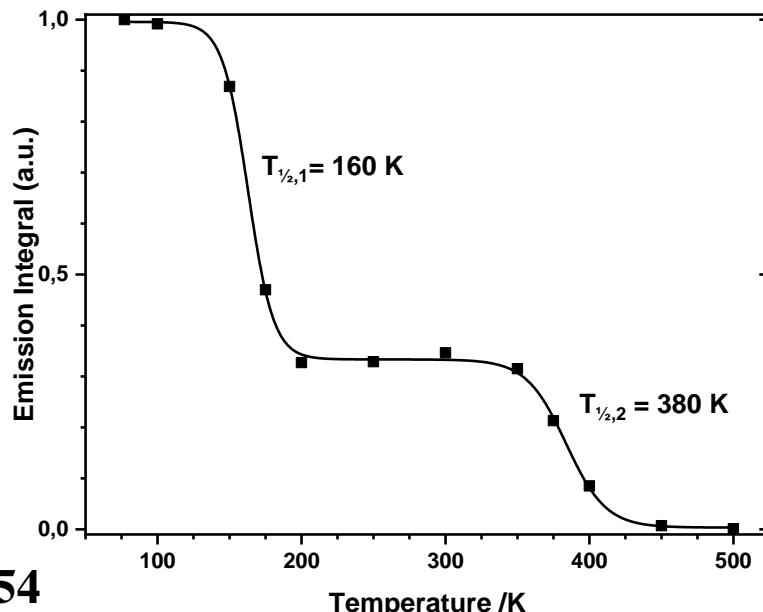
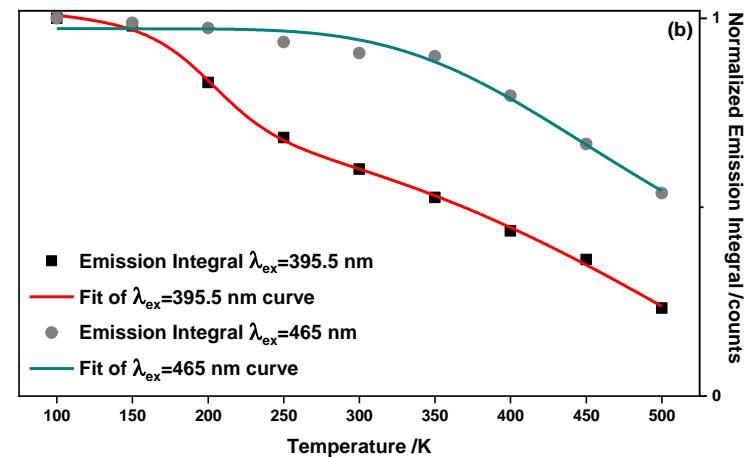
Re-absorption leads to the observed color point shift!

4.2.7 Temperature Resolved Spectroscopy

Special cases

- Simultaneous excitation via band gap and direct excitation of the activator can result in a bi-sigmoidal curve (example: $\text{LiEuMo}_2\text{O}_8$)
- Simultaneous emission of two luminescent centers in the same spectral region can also result in a bi-sigmoidal curve (example: $\text{Rb}_2\text{Ge}_4\text{O}_9:\text{Mn}^{4+}$)
- The curve can be fitted via weighted addition of two fitting functions. This yield two $T_{1/2}$ values:

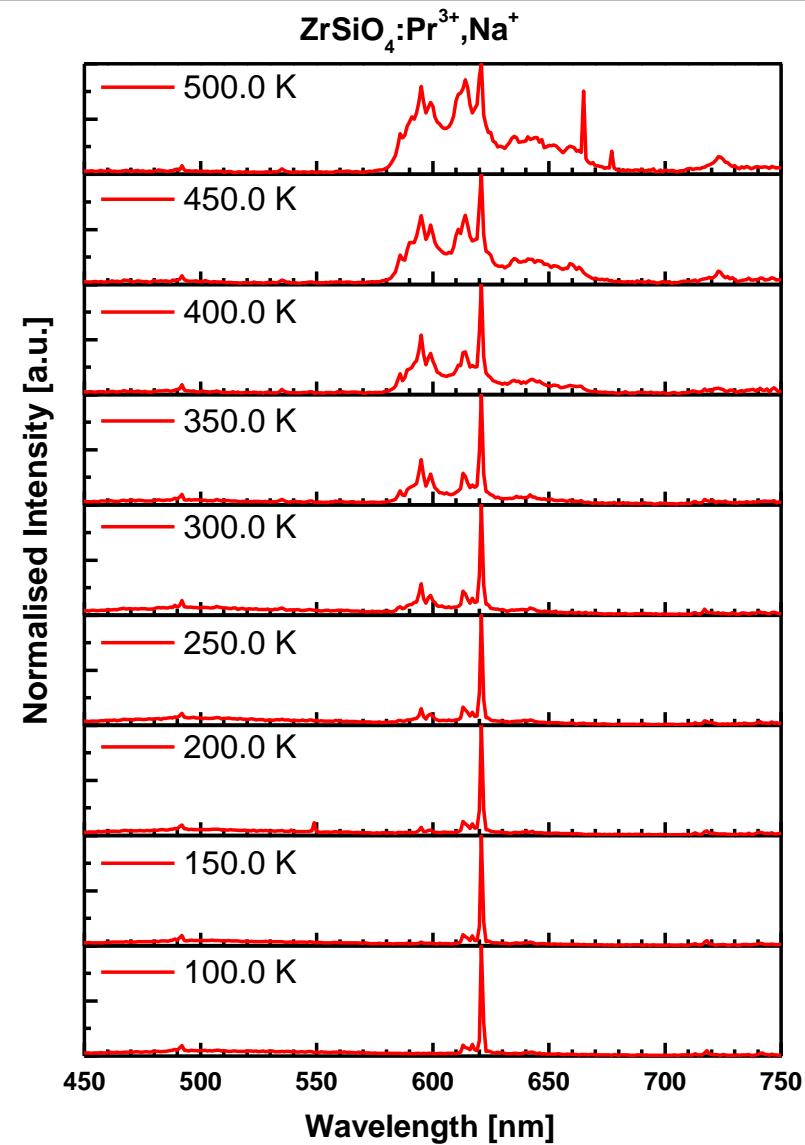
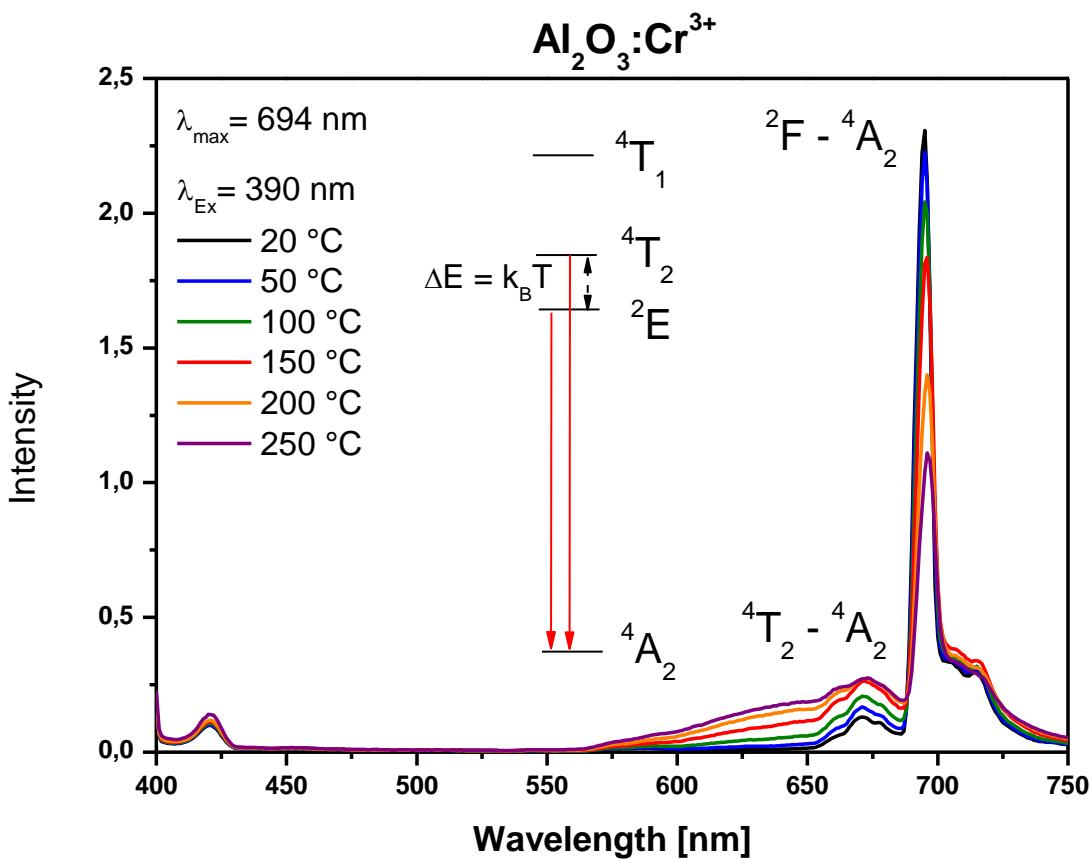
$$I_{\text{total}}(T) = A \cdot I_1(T) + (1-A) \cdot I_2(T)$$



Lit.: F. Baur, T. Jüstel, J. Luminescence 177 (2016) 354

4.2.7 Temperature Resolved Spectroscopy

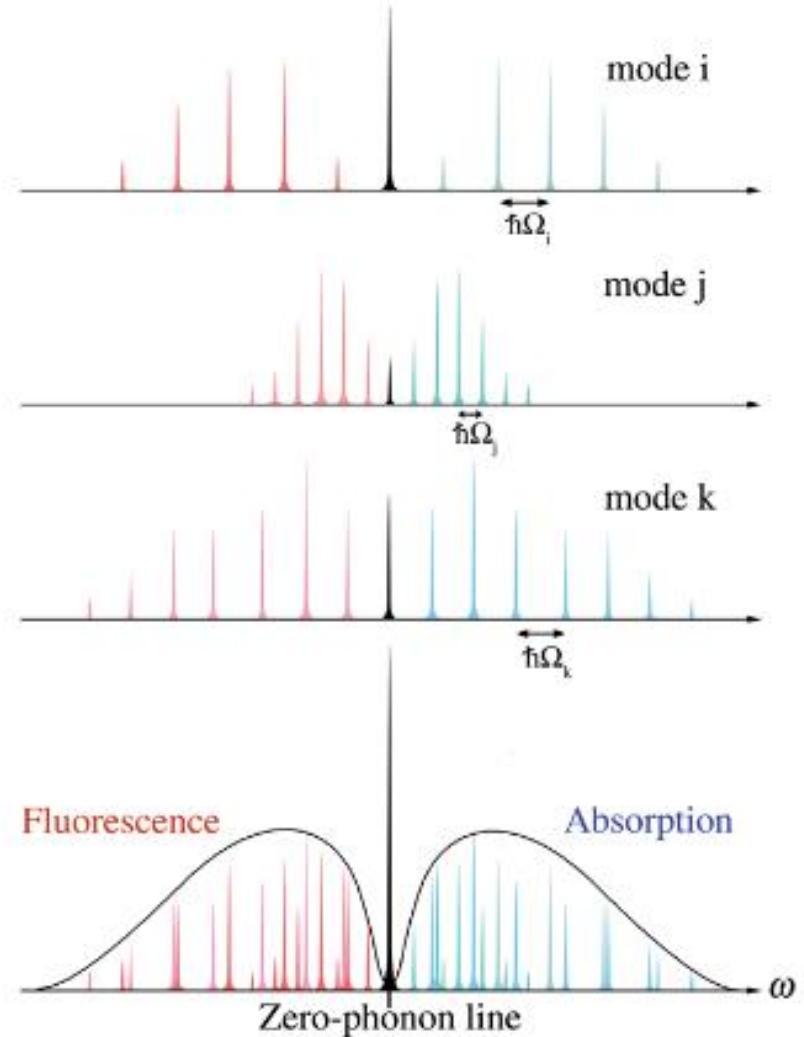
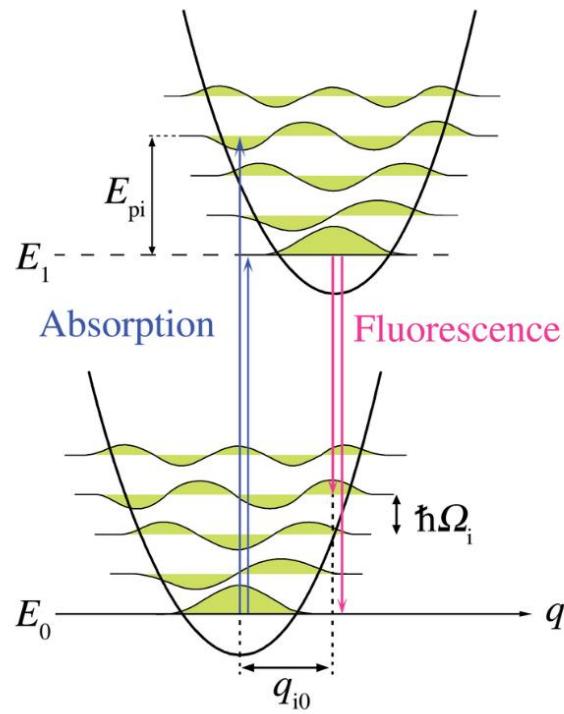
Thermal excitation of higher radiative states



4.2.7 Temperature Resolved Spectroscopy

Zero Phonon Line

The higher the temperature, the more is the luminescence influenced by lattice vibrations (phonons) – instead of a line, a band will be observed

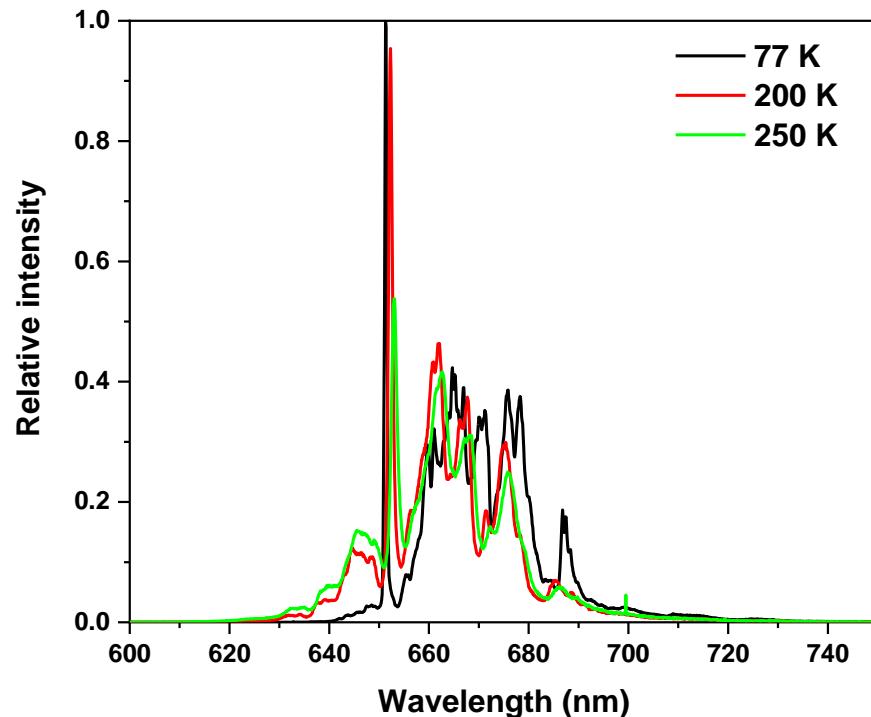


4.2.7 Temperature Resolved Spectroscopy

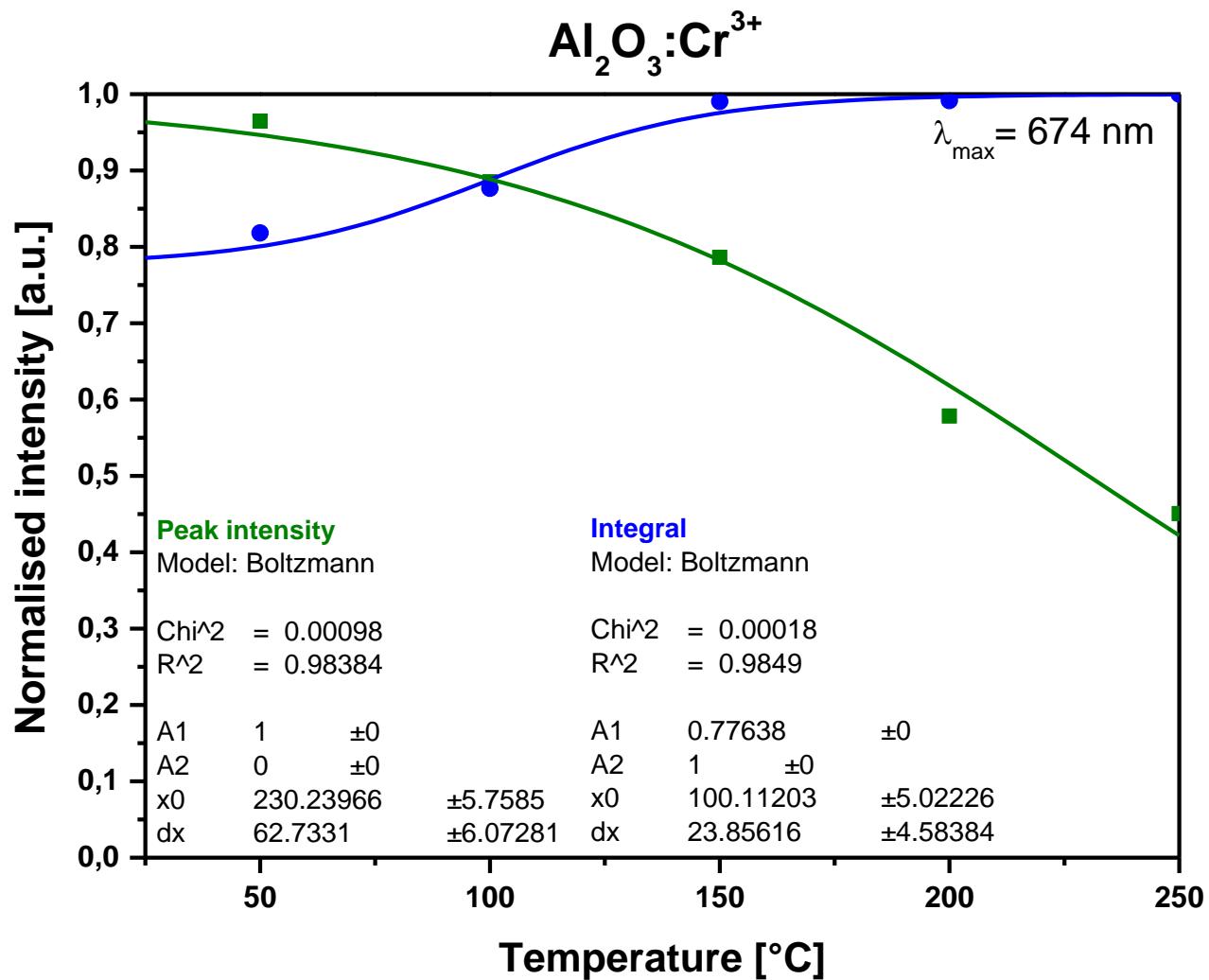
Zero Phonon Line

In emission spectra, the Stokes transitions are at a lower energy (generation of phonons / heat), the anti-Stokes transitions are at higher energy (extinction of phonons / cooling)

- The intensity and number of ZPLs depend on the site symmetry and number of dopant sites
- The spectral position of vibrational lines depends on the phonon energy

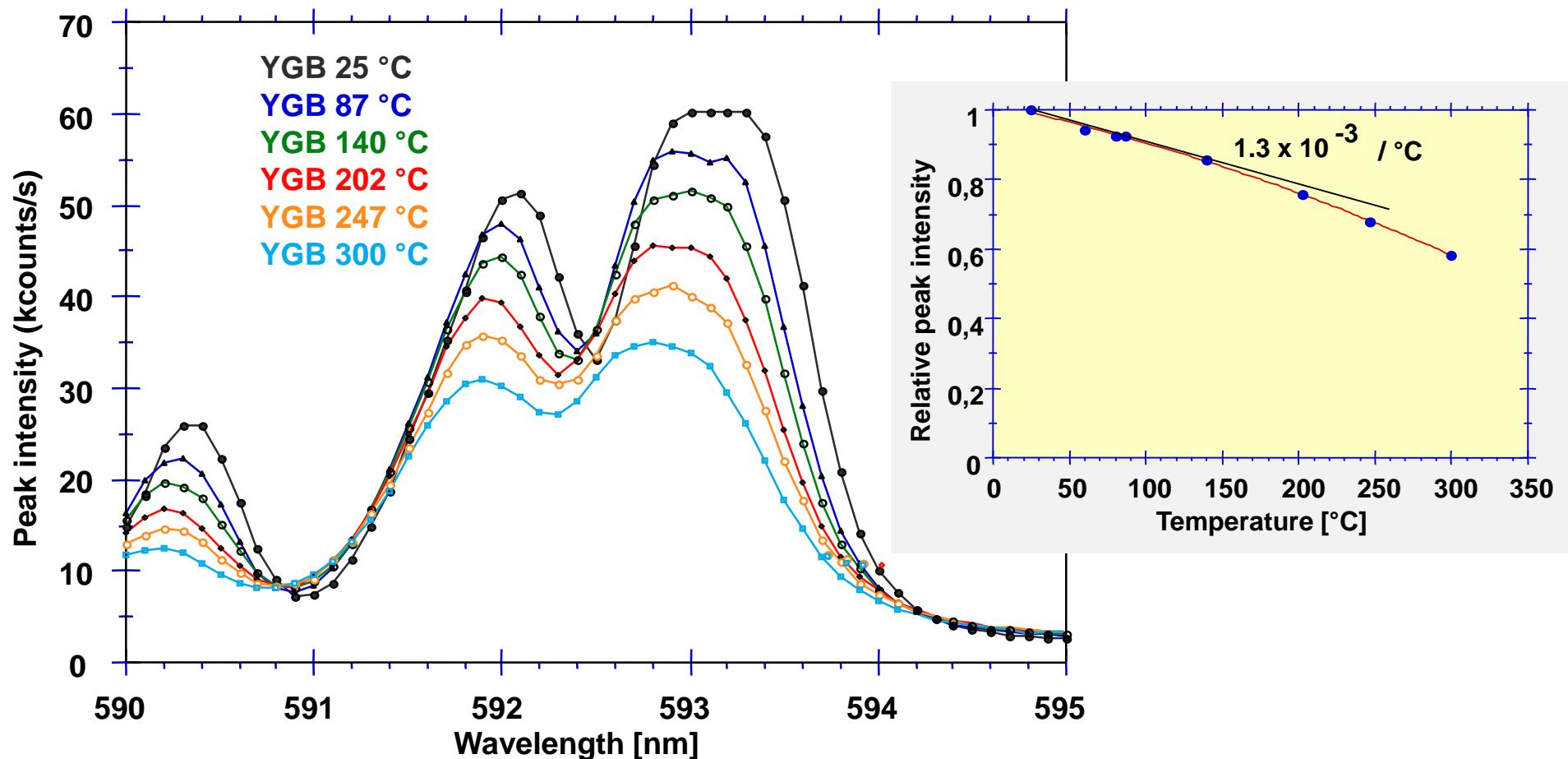


4.2.7 Temperature Resolved Spectroscopy



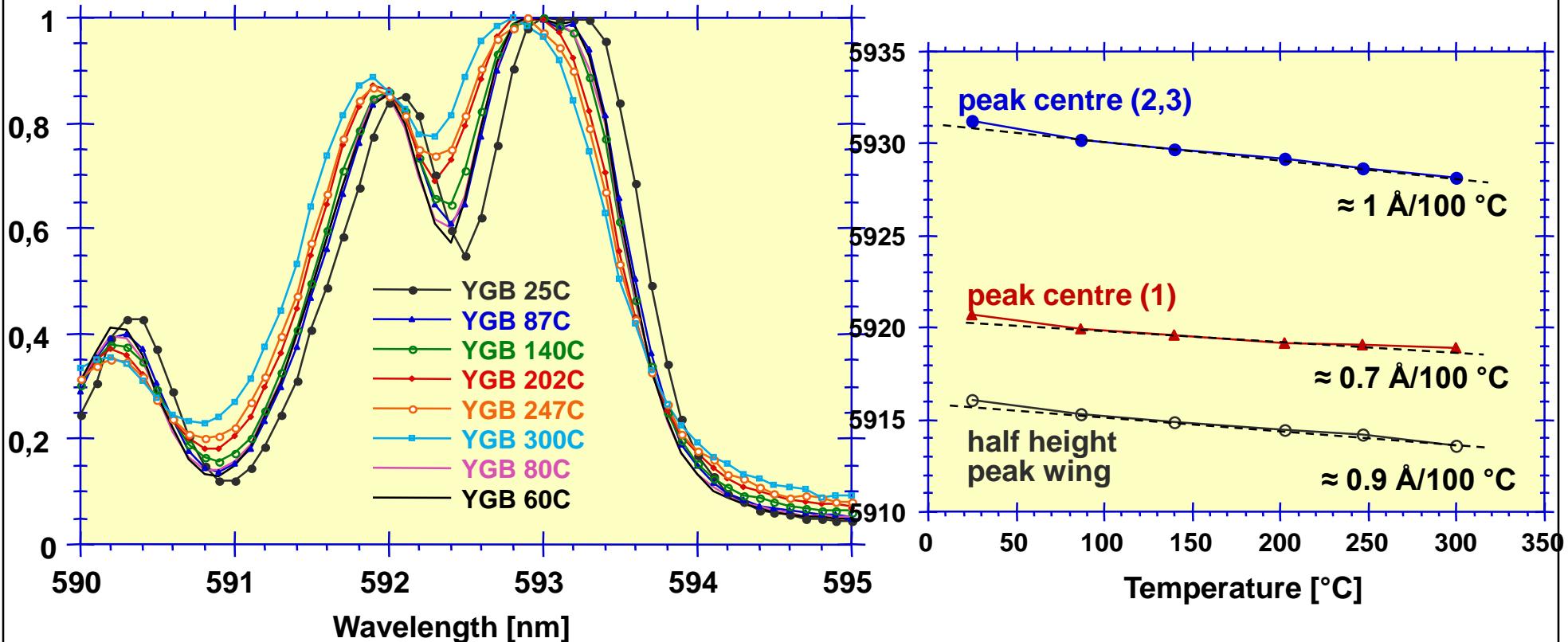
4.2.7 Temperature Resolved Spectroscopy

Red PDP phosphor $(Y,Gd)BO_3:Eu^{3+}$ (YGB:Eu) $\rightarrow ^5D_0 \rightarrow ^7F_1$ transition as a temperature sensor since it shows reversible change of the emission spectrum



4.2.7 Temperature Resolved Spectroscopy

Red PDP phosphor $(Y,Gd)BO_3:Eu^{3+}$ (YGB:Eu) $\rightarrow ^5D_0 \rightarrow ^7F_1$ transition as a temperature sensor since it shows reversible change of the emission spectrum



Optical temperature sensing:

- Line/band shift
- Line/band broadening
- Line/band intensity (ratio)

4.2.7 Temperature Resolved Spectroscopy

Thermoluminescence

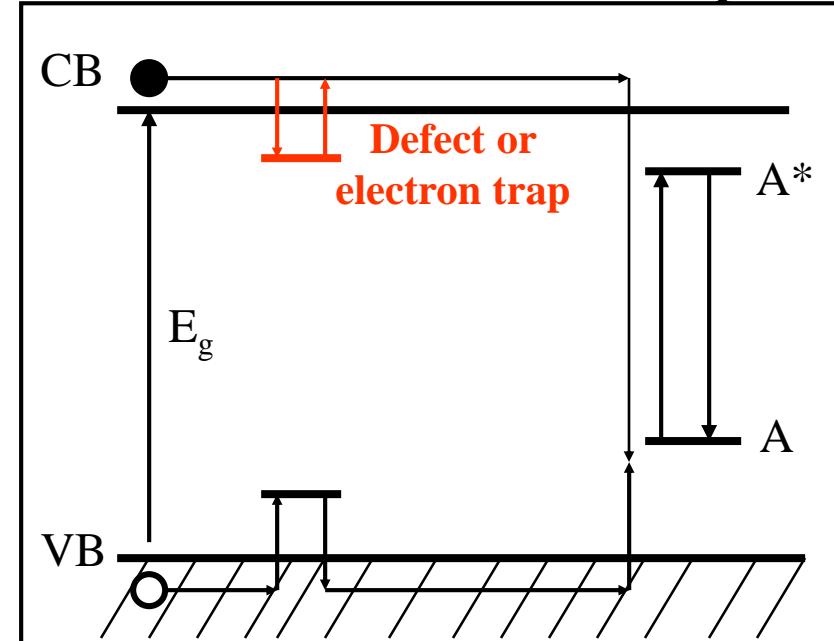
Technology

- Sample is cooled e.g. down to -196°C (77 K) or -269°C (4.2 K)
- Irradiation with UV-, VUV-, X-ray radiation or electrons
- The intensity of the emitted radiation is, during up-heating, measured with a constant time-dependent heating rate

Thermoluminescence (TL) (glow curves)

- Integral of the bands: Defect density
- Position of the bands: Energetic position of the defects
- Number of bands: Number of defect types

Band model of a solid state compound

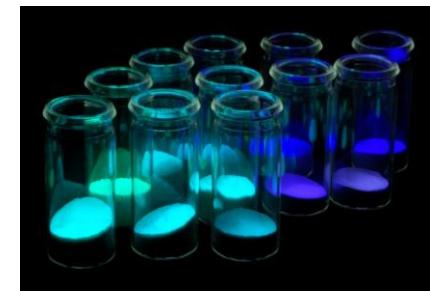


⇒ Insights in charge carrier trapping and luminescence quenching mechanisms

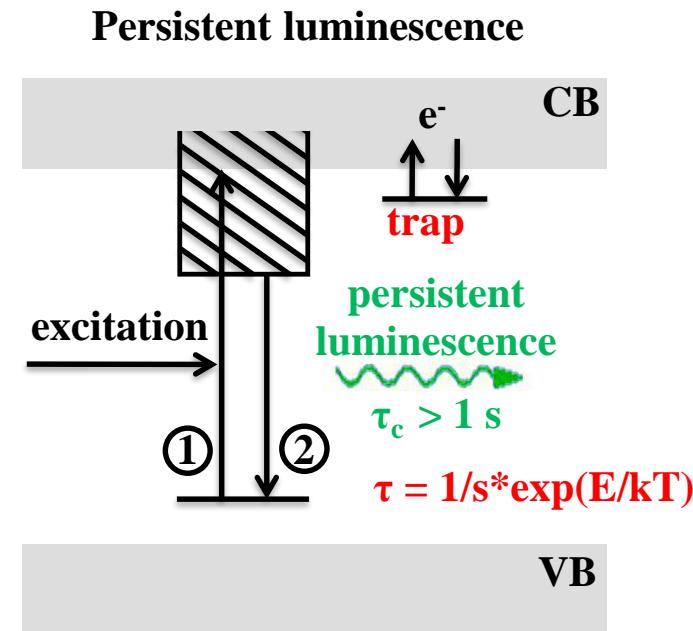
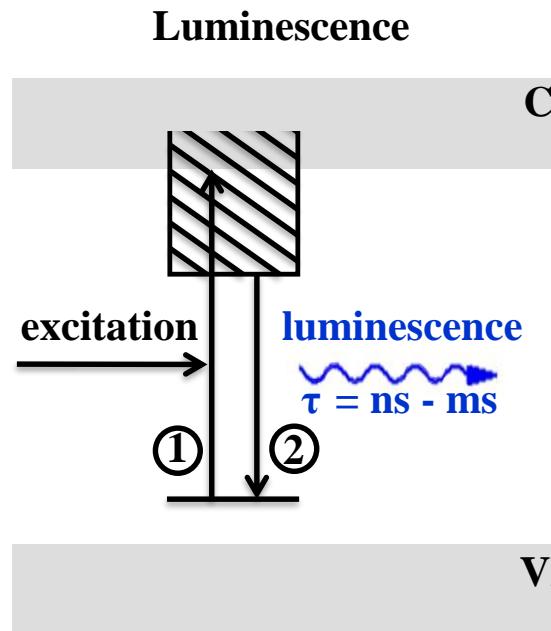
4.2.7 Temperature Resolved Spectroscopy

Thermoluminescence: Application Areas

- Age determination
- Dosimetry and x-ray diffraction
- Characterization of (persistent) luminescent materials
 - Afterglow and display phosphors, scintillators



$\text{Sr}_4\text{Al}_{14}\text{O}_{25}:\text{Eu}^{2+},\text{Dy}^{3+}$
samples



4.2.7 Temperature Resolved Spectroscopy

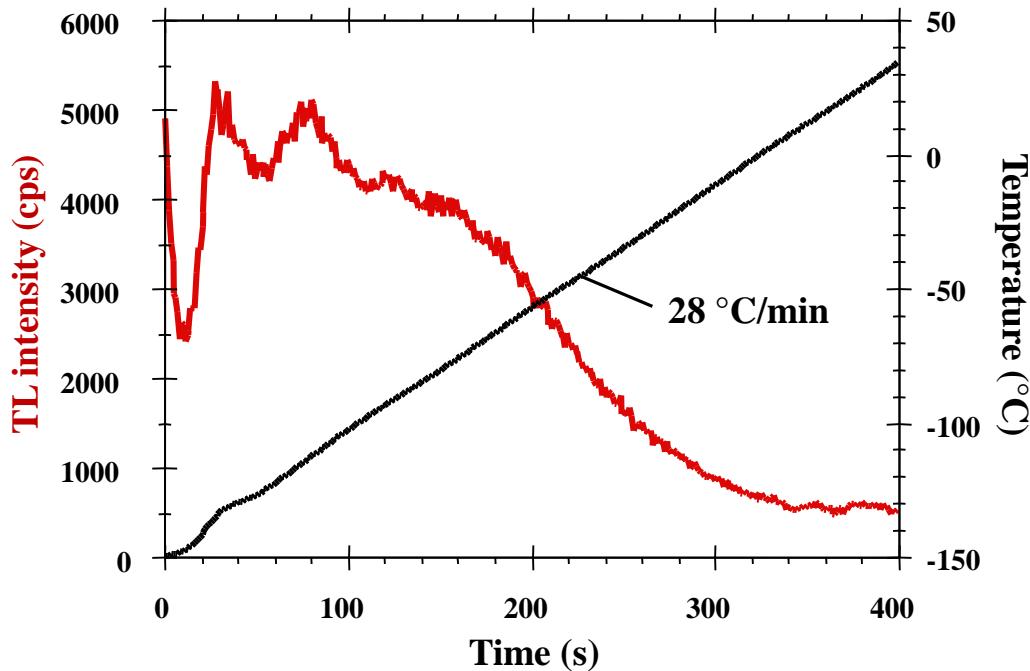
Thermoluminescence: Persistent luminescent materials

Host matrix	Luminescence center	Co-dopant	Emission maximum [nm]	Afterglow duration
MgAl_2O_4	Eu^{2+}	Dy^{3+}	480	> 1 min
CaAl_2O_4	Eu^{2+}	Nd^{3+}	440	> 5 h
$\text{Ca}_{12}\text{Al}_{14}\text{O}_{33}$	Eu^{2+}	Nd^{3+}	440	> 5 min
$\text{SrAl}_{12}\text{O}_{19}$	Eu^{2+}	Dy^{3+}	400	> 2 h
SrAl_4O_7	Eu^{2+}	Dy^{3+}	480	> 1 h
$\text{Sr}_4\text{Al}_{14}\text{O}_{25}$	Eu^{2+}	Dy^{3+}	405, 490	> 10 h
SrAl_2O_4	Eu^{2+}	Dy^{3+}	445, 520	> 10 h
$\text{Sr}_3\text{Al}_2\text{O}_6$	Eu^{2+}	Dy^{3+}	535, 620	> 10 min
BaAl_2O_4	Eu^{2+}	Dy^{3+}	505	> 2 h
$\text{SrMgAl}_{10}\text{O}_{17}$	Eu^{2+}	Dy^{3+}	460, 515	> 3 min
$\text{BaMgAl}_{10}\text{O}_{17}$	Eu^{2+}	Co^{3+}	450	> 5 min
$\text{BaCa}_2\text{Al}_8\text{O}_{15}$	Eu^{2+}	Dy^{3+}	435	> 5 min
$\text{Ca}_2\text{Si}_5\text{N}_8$	Eu^{2+}	Tm^{3+}	610	> 1 h
$\text{CaAl}_2\text{B}_2\text{O}_7$	Eu^{2+}	Nd^{3+}	464	> 1 h
$\text{SrAl}_{1.7}\text{B}_{0.3}\text{O}_4$	Eu^{2+}	-	520	> 2 h
$\text{Ca}_2\text{P}_2\text{O}_7$	Eu^{2+}	Y^{3+}	415	> 6 h
$\text{Sr}_2\text{P}_2\text{O}_7$	Eu^{2+}	Y^{3+}	420	> 8 h
$\text{SrMg}_2\text{P}_2\text{O}_8$	Eu^{2+}	$\text{Ce}^{3+}/\text{Gd}^{3+}$	400	> 2 h
CaS	Eu^{2+}	$\text{Tm}^{3+}/\text{Pr}^{3+}$	650	> 1 h
CaGa_2S_4	Eu^{2+}	Ho^{3+}	555	> 30 min
Ca_2SiS_4	Eu^{2+}	Nd^{3+}	660	> 30 min
$\text{Sr}_3\text{Al}_2\text{O}_5\text{Cl}_2$	Eu^{2+}	$\text{Tm}^{3+}/\text{Dy}^{3+}$	620	> 5 min

4.2.7 Temperature Resolved Spectroscopy

Thermoluminescence

Example: BaMgAl₁₀O₁₇:Eu



Arrhenius equation

$$A = s * \exp(-E/kT)$$

s = frequency factor [s⁻¹]

E = trap depth [eV]

Activation i.e., emptying
the electrons in the traps
leads to emission of radiation

First order kinetics i.e.,
no retrapping:

$$I(t) = -dn_t/dt = -n_t s * \exp(-E/kT)$$

“Glow curve” = Luminescence intensity as a function of time or temperature

4.2.7 Temperature Resolved Spectroscopy

Randall-Wilkins first-order glow peak

Intensity at a given temperature

$$I(t) = -dn_t/dt = -n_t s * \exp(-E/kT)$$

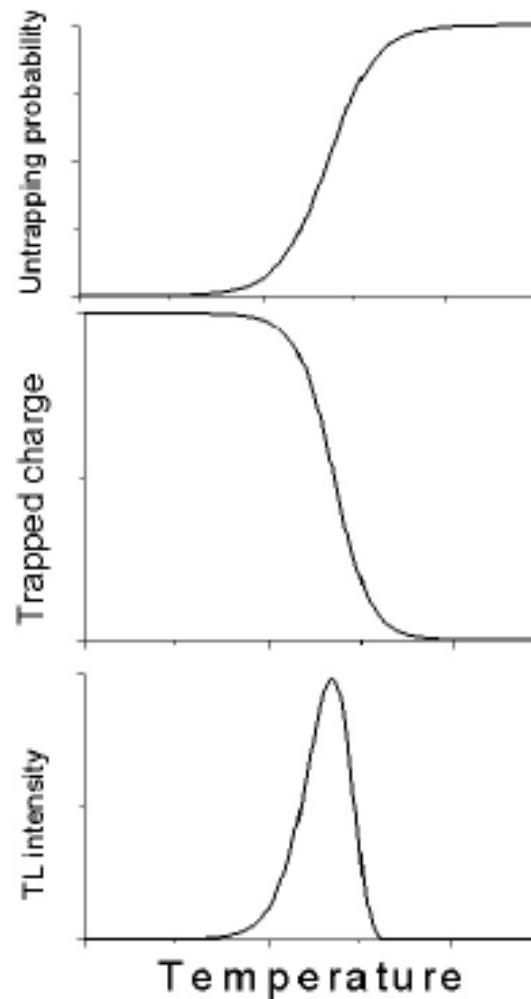
Temperature profile in TL experiments

$$T(t) = T_0 + \beta \cdot t$$

T_0 = start temperature

β = heating rate

k or $k_B = 1.38 \cdot 10^{-23} \text{ J K}^{-1}$ (Boltzmann constant)



4.2.7 Temperature Resolved Spectroscopy

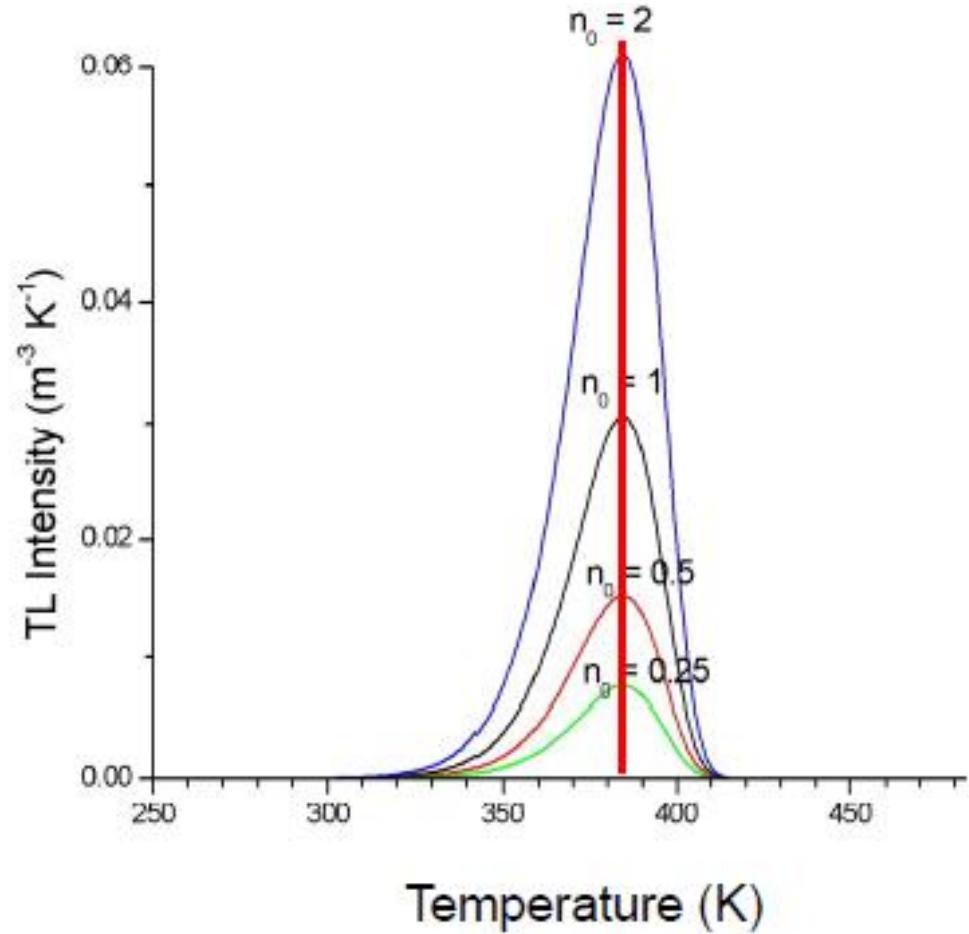
First-order glow peak - Variation in n_0 (Number of traps)

$$\int_0^{\infty} I(t) dt = - \int_0^{\infty} \frac{dn}{dt} dt = - \int_{n_0}^{n_{\infty}} dn = n_0$$

Findings for n_0

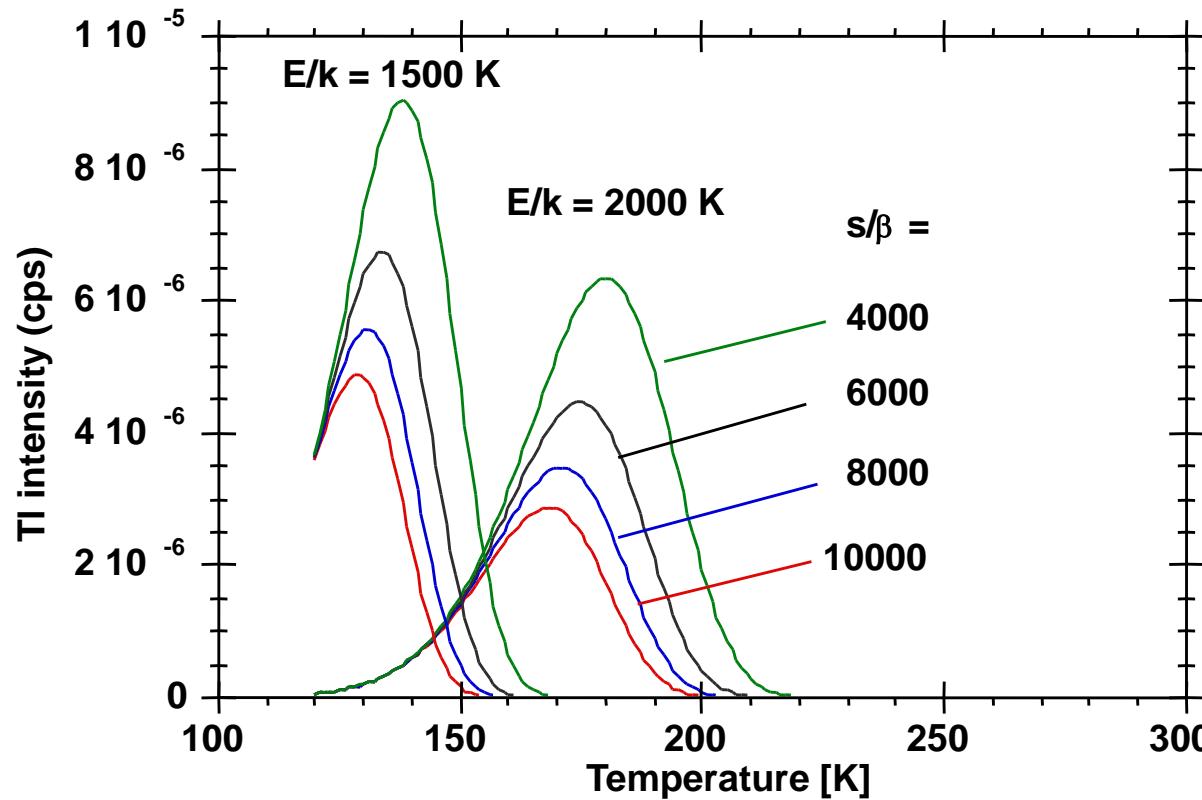
- ~ Peak area
- ~ Peak height
- Independent of peak maximum

→ relates to defect density of material



4.2.7 Temperature Resolved Spectroscopy

First-order glow peak - Variation in E and s or β



Assumption: 1st order kinetics

- N_{T0} : Number of traps
- s : Frequency factor
- β : Heating rate
- E : trap depth

$$I_{TL} = N_{T0} s \exp(-E/kT) \cdot \exp\left[-s/\beta \int_{T_0}^T \exp(-E/kT)dT\right]$$

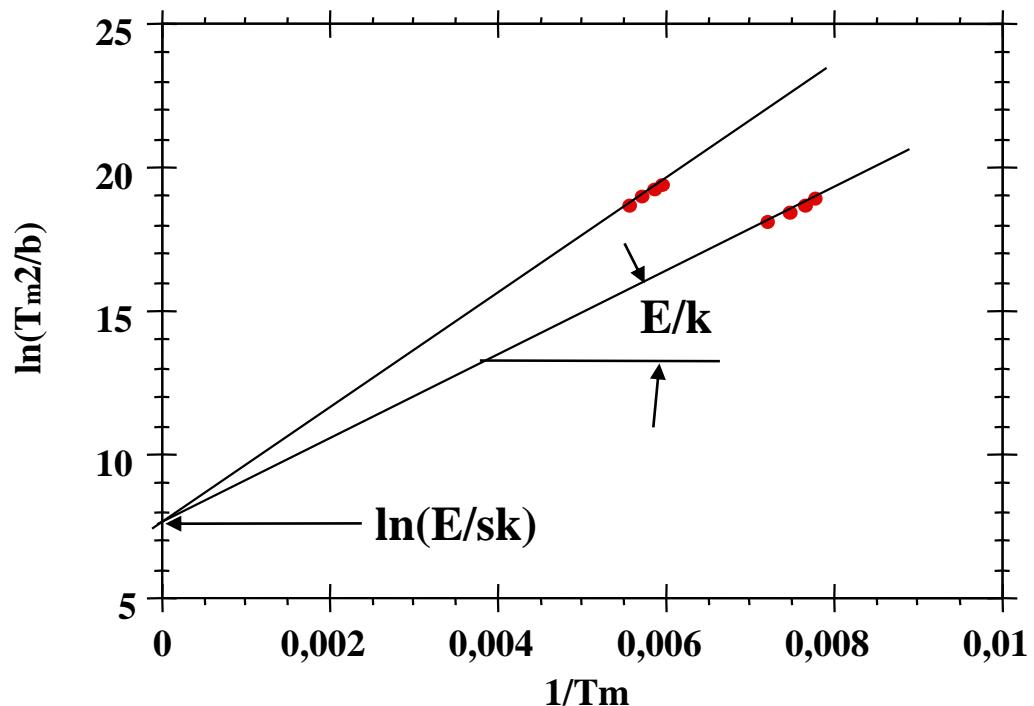
Randall, Wilkins 1945

4.2.7 Temperature Resolved Spectroscopy

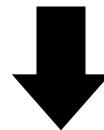
Analysis of thermoluminescence measurements

$$I_{TL} = N_{T_0} s \exp(-E/kT) \cdot \exp\left[-s/\beta \int_{T_0}^T \exp(-E/kT)dT\right]$$

First order kinetics



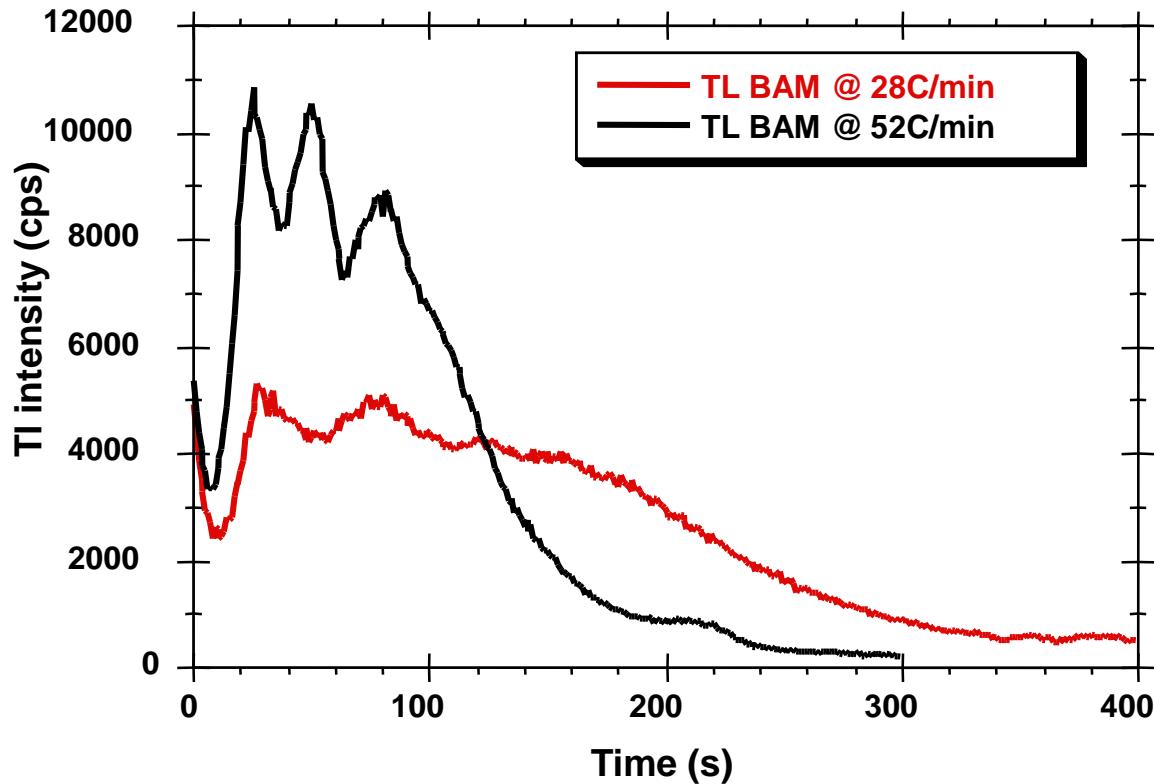
$$\frac{dI_{TL}}{dT} = 0 \text{ at } T = T_m$$



$$\frac{\beta \cdot E}{k T_m^2} = s \cdot \exp(-E/kT_m)$$

4.2.7 Temperature Resolved Spectroscopy

a) Thermoluminescence analysis of $\text{BaMgAl}_{10}\text{O}_{17}:\text{Eu}$ (BAM:Eu)



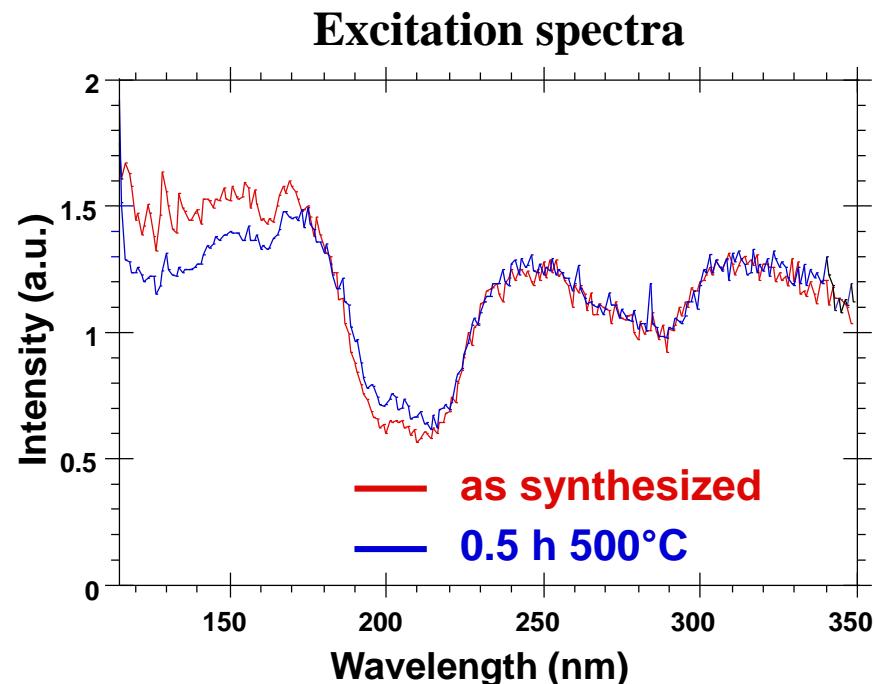
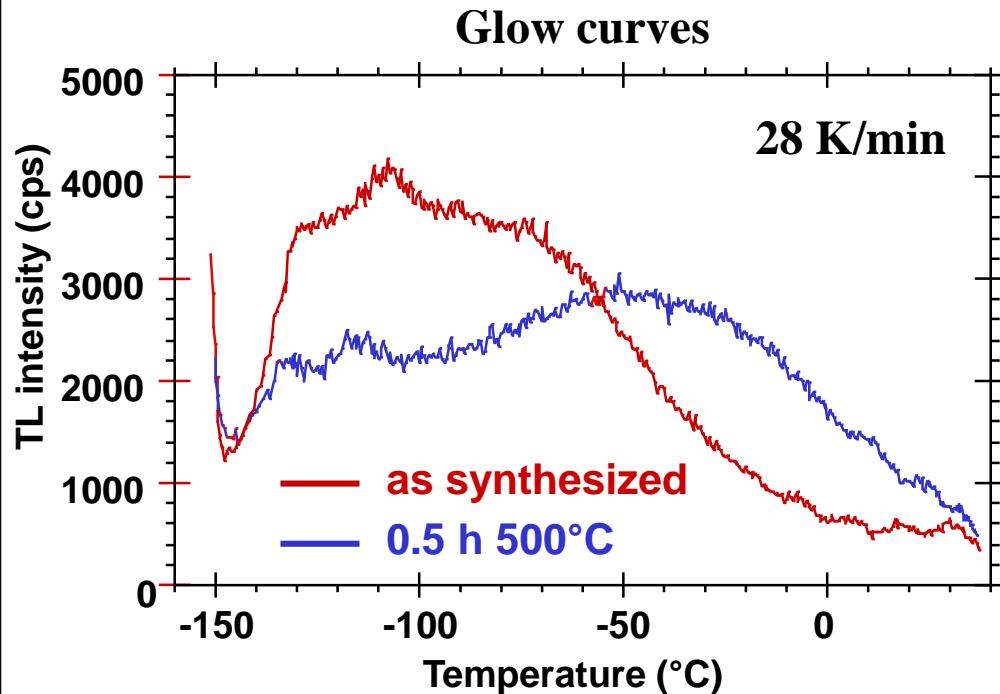
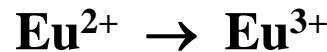
As a function of heating rate

β	$\int I_{\text{TL}} dt$
[°C/min]	[Mcps*s]
11	0.898
15	1.117
22	1.104
28	1.050
45	1.027
52	1.078

4.2.7 Temperature Resolved Spectroscopy

a) Thermoluminescence analysis of $\text{BaMgAl}_{10}\text{O}_{17}:\text{Eu}$ (BAM:Eu)

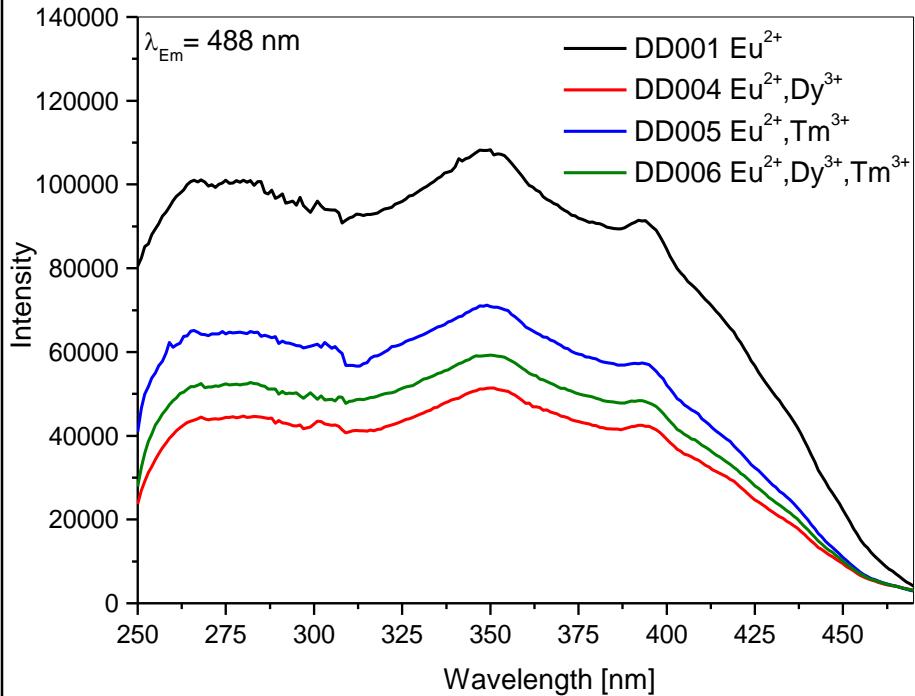
BAM:Eu is at $T > 300$ °C sensitive to oxidation:



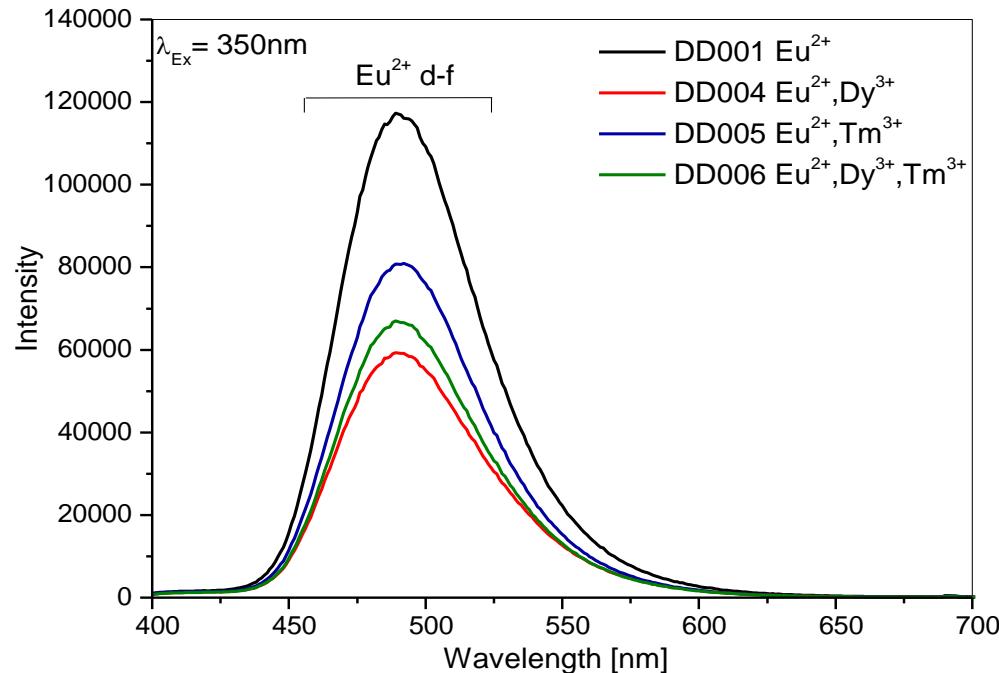
"High temperature" TL traps are generated by the oxidation (direct evidence for the formation of Eu^{3+})

4.2.7 Temperature Resolved Spectroscopy

b) Luminescence spectra of $\text{Sr}_4\text{Al}_{14}\text{O}_{25}:\text{Eu}(\text{X})$ with $\text{X} = \text{Dy}, \text{Nd}, \text{Tm}$



Excitation spectrum
of $\text{Sr}_4\text{Al}_{14}\text{O}_{25}:\text{Eu}^{2+},\text{X}$



Emission spectrum
of $\text{Sr}_4\text{Al}_{14}\text{O}_{25}:\text{Eu}^{2+},\text{X}$

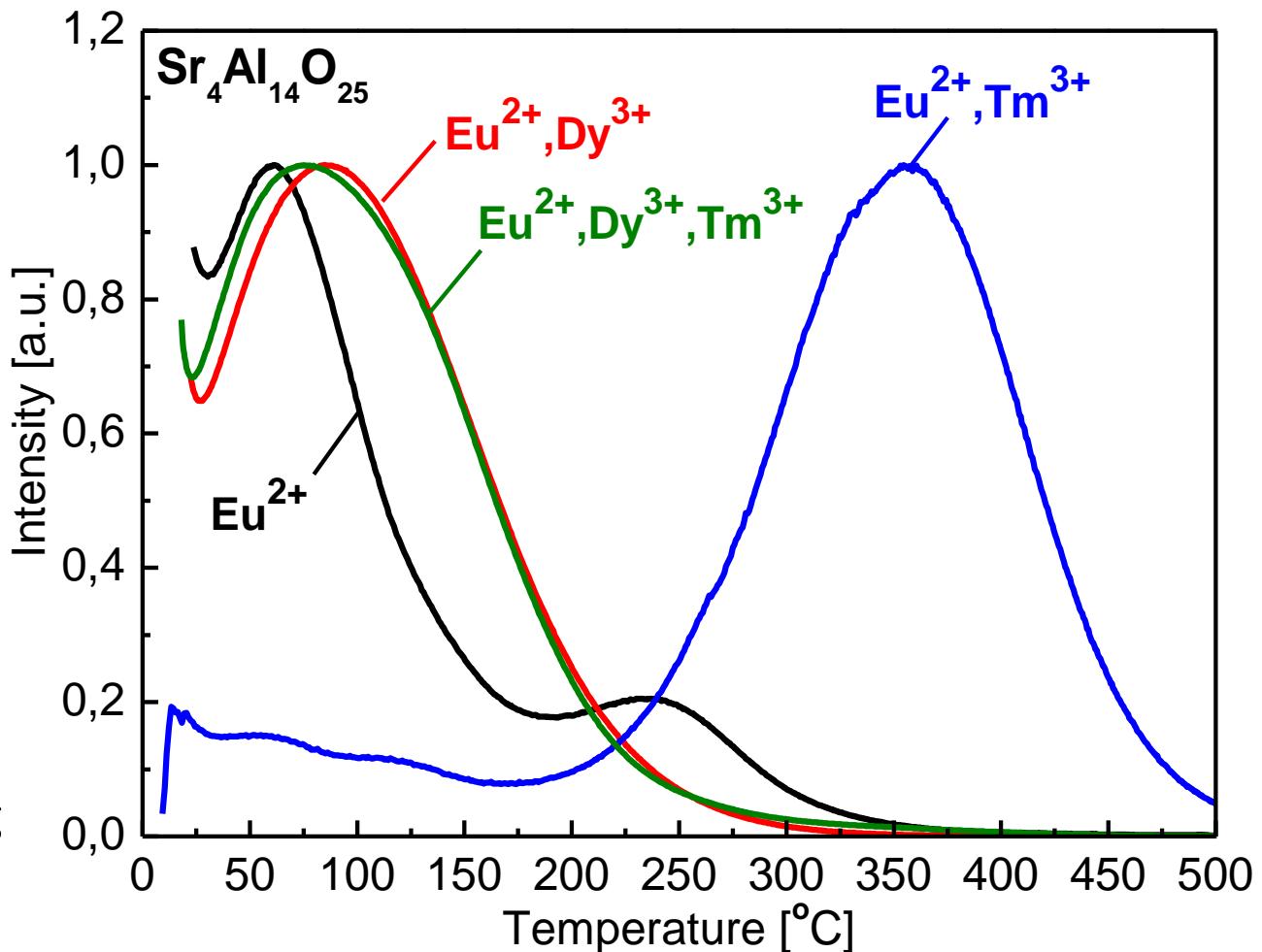
4.2.7 Temperature Resolved Spectroscopy

b) Thermoluminescence analysis of $\text{Sr}_4\text{Al}_{14}\text{O}_{25}:\text{Eu}$

Co-doping with Dy^{3+} results in „afterglow“ at room temperature

Further doping with Tm^{3+} results in „afterglow“ also at high temperature
(→ Storage phosphors)

As soon as Dy^{3+} is incorporated, Tm^{3+} hardly plays a role in the trapping of charge carriers

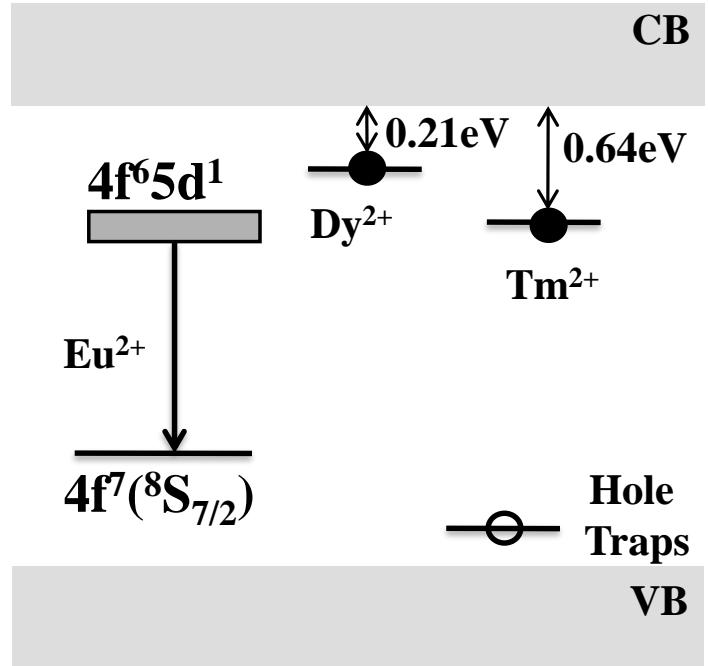


4.2.7 Temperature Resolved Spectroscopy

b) Thermoluminescence analysis of $\text{Sr}_4\text{Al}_{14}\text{O}_{25}:\text{Eu}$

Data from the "peak shape" method

Compound	„Trap“ depth [eV]
$\text{Sr}_4\text{Al}_{14}\text{O}_{25}:\text{Eu}^{2+}$	0.26, 0.55
$\text{Sr}_4\text{Al}_{14}\text{O}_{25}:\text{Eu}^{2+}, \text{Dy}^{3+}$	0.21
$\text{Sr}_4\text{Al}_{14}\text{O}_{25}:\text{Eu}^{2+}, \text{Tm}^{3+}$	0.64
$\text{Sr}_4\text{Al}_{14}\text{O}_{25}:\text{Eu}^{2+}, \text{Dy}^{3+}, \text{Tm}^{3+}$	0.23



- Dy^{2+/3+} is energetically closer to the conduction band than Tm^{2+/3+}
- Energy gap between Dy and Tm is about 0.4 eV
(see also publications by Pieter Dorenbos, TU Delft)

4.2.8 Time Resolved Spectroscopy

Requirement: Pulsed, highly intense light source

- μ s-flash lamps
- ns-flash lamps
- Laser: N₂, excimer, Al₂O₃:Cr, YAG:Nd,
- LEDs: (Al,In,Ga)N, (Al,In,Ga)P, Ga(As,P),

Measurement Procedure

- Pulsed or continuous wave excitation of the sample
- Measurement of the luminescence intensity as the function of time after the excitation source has been switched off (or single photon counting)
- Fitting of the obtained decay curve with one or several e-functions:

$$I(t) = A_0 + B_1 * \exp(-t/\tau_1) + B_2 * \exp(-t/\tau_2) + \dots \quad \text{with } B_1 + B_2 + \dots + B_n = 1.0$$

4.2.8 Time Resolved Spectroscopy

Lifetime of the excited state

First order kinetics (two level system, no quenching)

$$dN_e/dt = - N_e * k_{eg}$$

$$\Rightarrow dN_e/N_e = -k_{eg} * dt$$

$$\Rightarrow \ln(N_e(t)/N_e(0)) = -k_{eg} * t$$

$$\Rightarrow N_e(t) = N_e(0) * \exp(-t/\tau)$$

Integration

$$\tau = 1/k_{eg} = 1/k_r$$

Allowed transitions

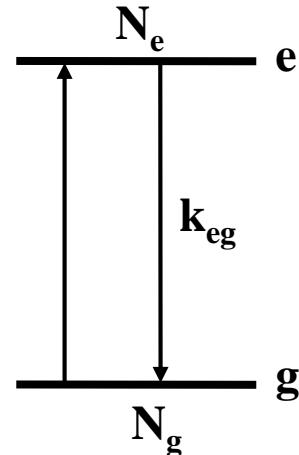
10^{-6} to 10^{-9} s (4f - 5d, 6s - 6p)

Eu^{2+}

Forbidden transitions

$\sim 10^{-3}$ s (4f - 4f, 5d - 5d)

Eu^{3+}



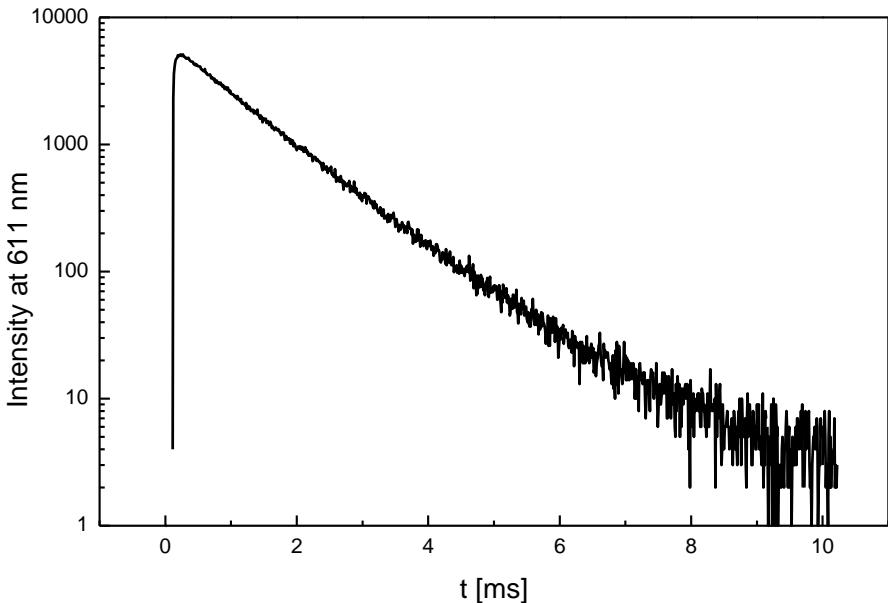
4.2.8 Time Resolved Spectroscopy

Mono- and bi-exponential decay

mono-exponential behavior, e.g. Eu³⁺, Gd³⁺, Tb³⁺

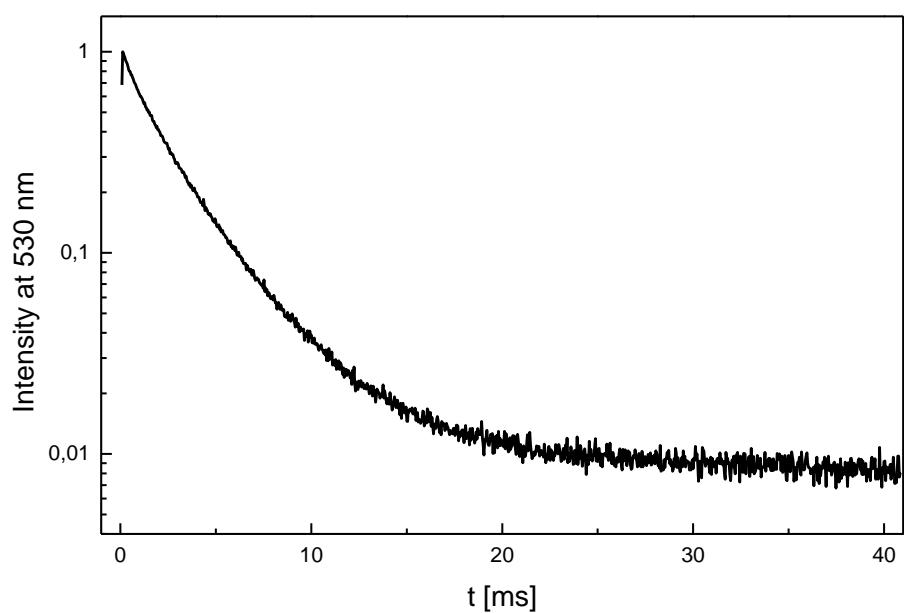
bi-exponential behavior, e.g. Mn²⁺ due to Mn²⁺-O-Mn²⁺ antiferromagnetic coupling

(Y,Gd)₂O₃:Eu³⁺



$$B_1 = 1.0 \quad \tau_1 = 1.1 \text{ ms}$$

Zn₂SiO₄:Mn²⁺



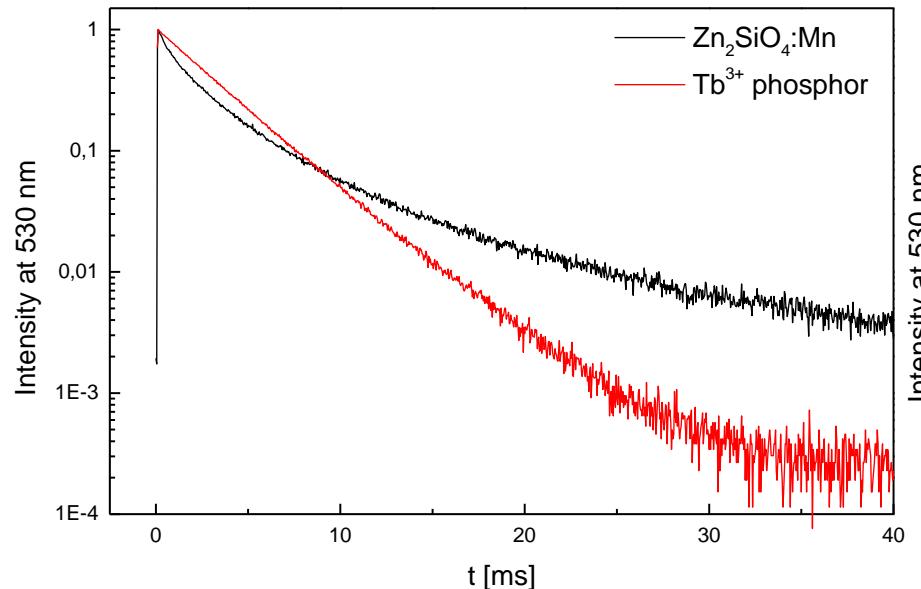
$$B_1 = 0.44 \quad \tau_1 = 5.6 \text{ ms}$$

$$B_2 = 0.56 \quad \tau_2 = 2.3 \text{ ms}$$

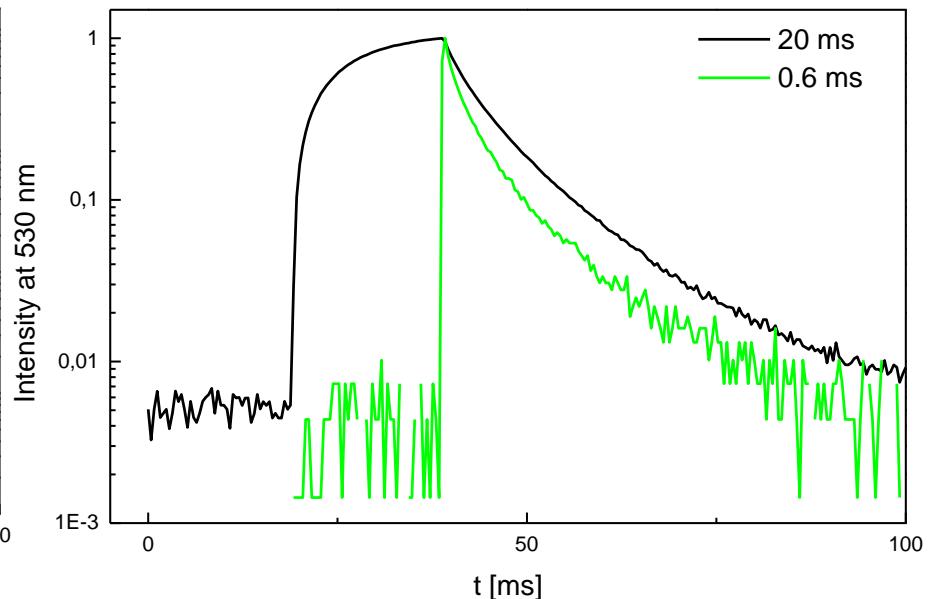
4.2.8 Time Resolved Spectroscopy

Mono- and bi-exponential decay

Single-pulse excitation of
 $\text{Zn}_2\text{SiO}_4:\text{Mn}^{2+}$ and $(\text{Y},\text{Gd})\text{BO}_3:\text{Tb}^{3+}$



Decay curve of $\text{Zn}_2\text{SiO}_4:\text{Mn}^{2+}$
(Single pulse and continuous wave
excitation)



Phosphor

- $(\text{Y},\text{Gd})\text{BO}_3:\text{Tb}^{3+}$
- $\text{Zn}_2\text{SiO}_4:\text{Mn}^{2+}$

Decay curve

- mono-exponential
- bi-exponential

Single pulse

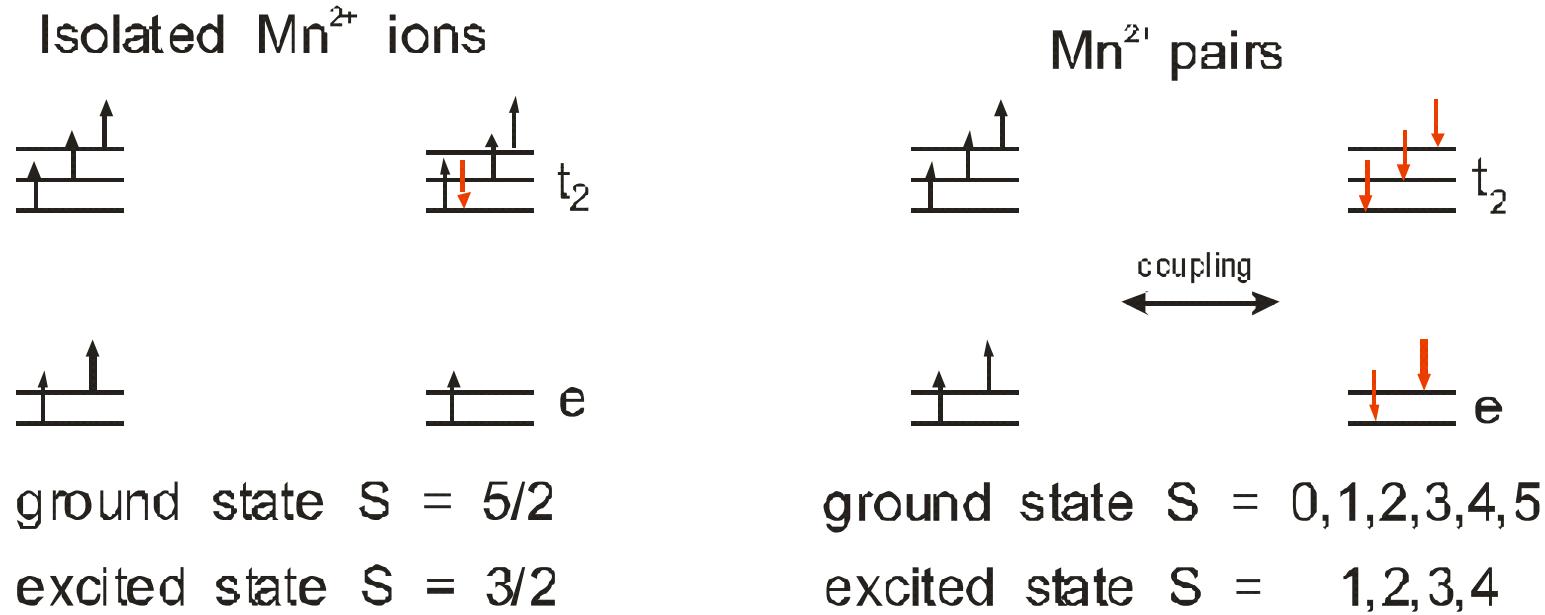
- 9 ms
- 11 ms

Continuous wave

- 10 ms
- 17 ms

4.2.8 Time Resolved Spectroscopy

Exchange interaction in Mn^{2+} phosphors



High Mn^{2+} concentration \Rightarrow Exchange interaction between Mn^{2+} -ions

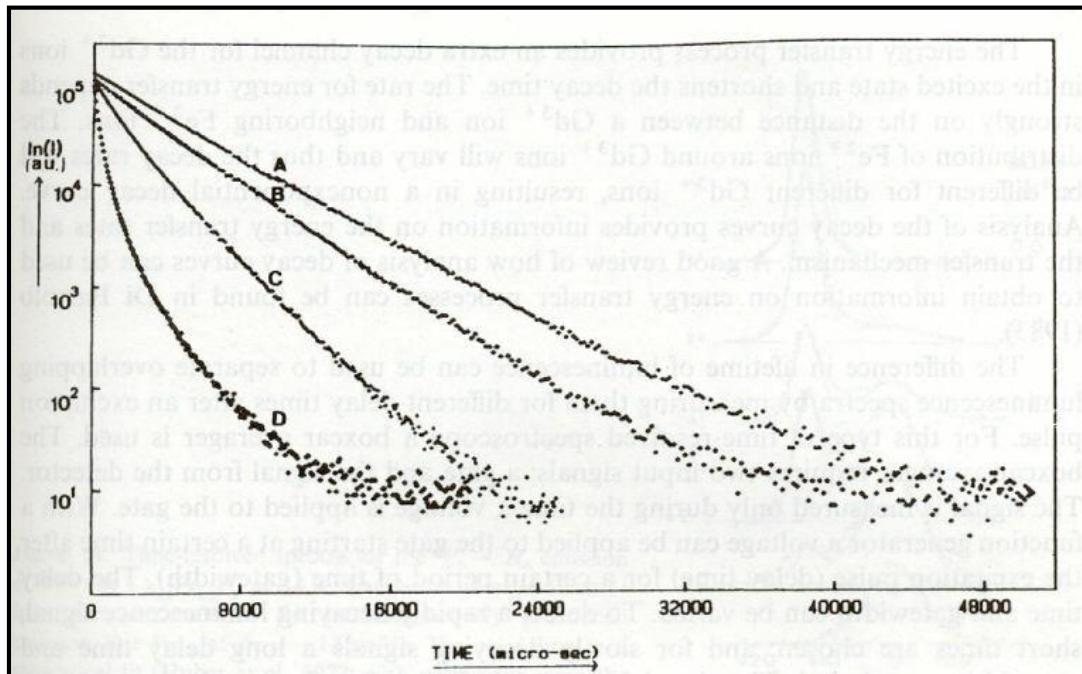
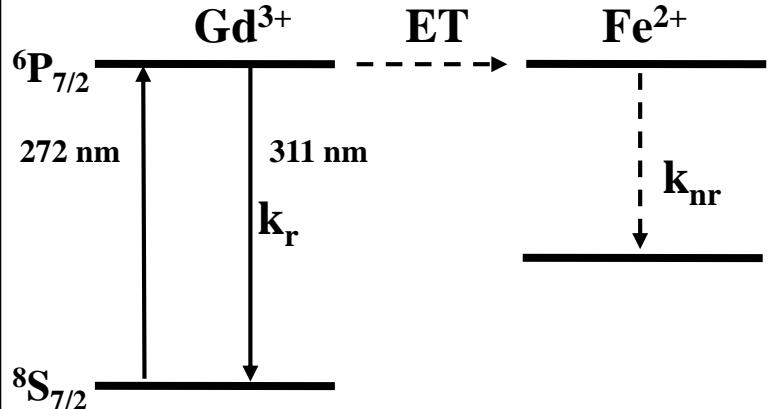
The decay time is reduced by eliminating the spin selection rule

Low excitation density \Rightarrow Excitation of Mn^{2+} pairs \Rightarrow Short decay time

High excitation density \Rightarrow Excitation of isolated Mn^{2+} ions \Rightarrow Long decay time

4.2.8 Time Resolved Spectroscopy

Example: $\text{La}_{0.4}\text{Gd}_{0.6}\text{MgB}_5\text{O}_{10}:\text{x}\%\text{Fe}$



Decay of the Gd^{3+} emission (intraconfigurational transition $^6\text{P}_{7/2} - ^8\text{S}_{7/2}$) at 311 nm

- A: $\text{La}_{0.4}\text{Gd}_{0.6}\text{MgB}_5\text{O}_{10}$
- B: $\text{La}_{0.4}\text{Gd}_{0.6}\text{MgB}_5\text{O}_{10}:\text{0.01\% Fe}$
- C: $\text{La}_{0.4}\text{Gd}_{0.6}\text{MgB}_5\text{O}_{10}:\text{0.1\% Fe}$
- D: $\text{La}_{0.4}\text{Gd}_{0.6}\text{MgB}_5\text{O}_{10}:\text{1\% Fe}$

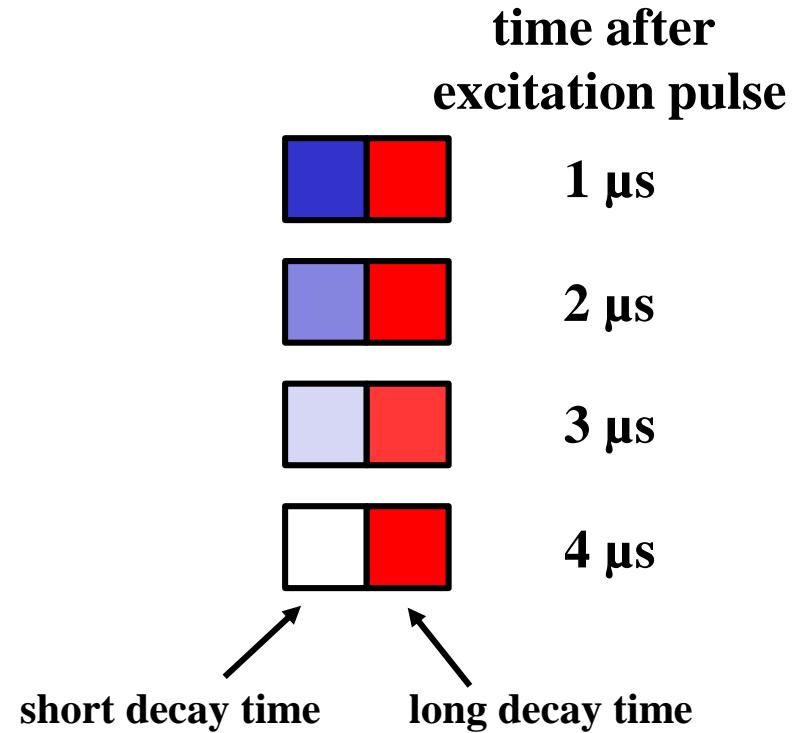
\Rightarrow Energy transfer to Fe^{2+} reduces decay time

4.2.8 Time Resolved Spectroscopy

Time Resolved Emission Spectra

Emission spectrum and/or intensity can change over time e.g. due to

- Two or more emitting components emission with different decay times
- saturation effects
- energy transfer



In order to observe these effects, the system must not be in equilibrium
→ Pulsed excitation!

Example for a system comprising a fast and a slow decaying component

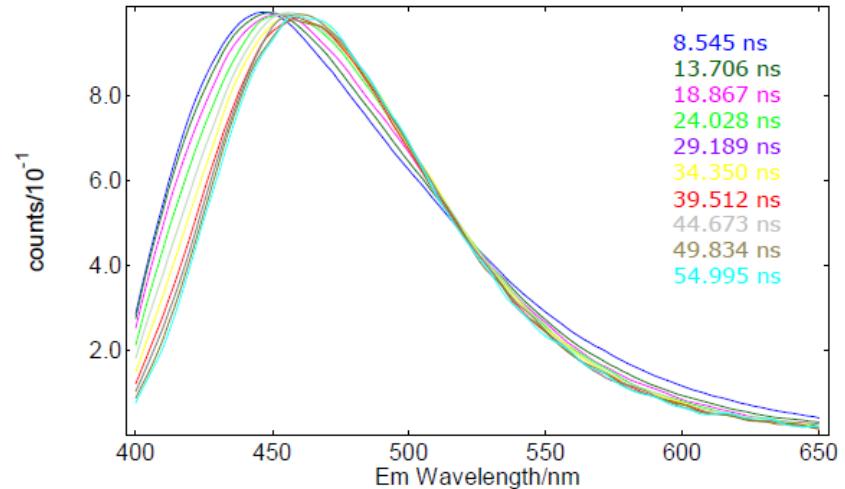
e.g. Eu^{2+} and Eu^{3+}

4.2.8 Time Resolved Spectroscopy

Time Resolved Emission Spectra

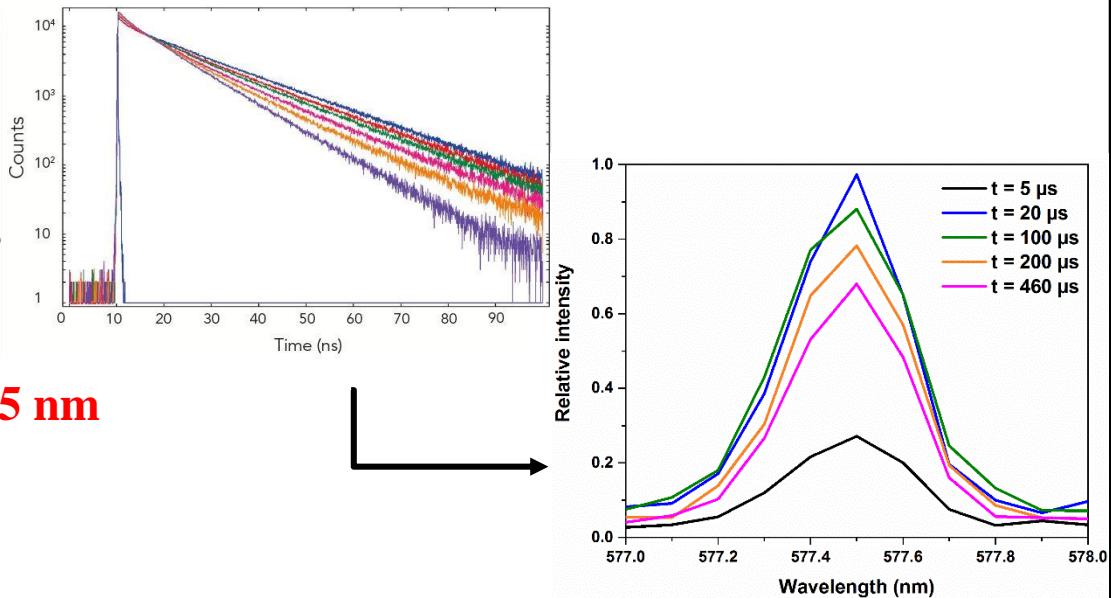
Recording by ICCD (intensified CCD)

- CCD with amplifier
- all wavelengths are recorded simultaneously
- down to **200 ps** temporal resolution
- wavelength resolution c. **0.5-1 nm**
- **quick** measurement



Recording via decay curves

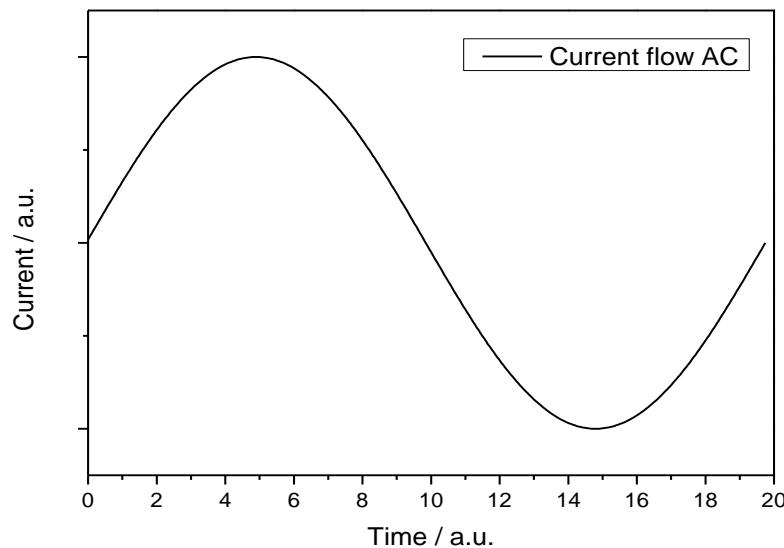
- A decay curve is recorded for each emission wavelength
- A time resolved emission spectrum is constructed from the curves
- down to **few ps** temporal resolution
- wavelength resolution of down to **0.05 nm**
- **(very) slow** measurement



4.2.9 Flicker Measurements

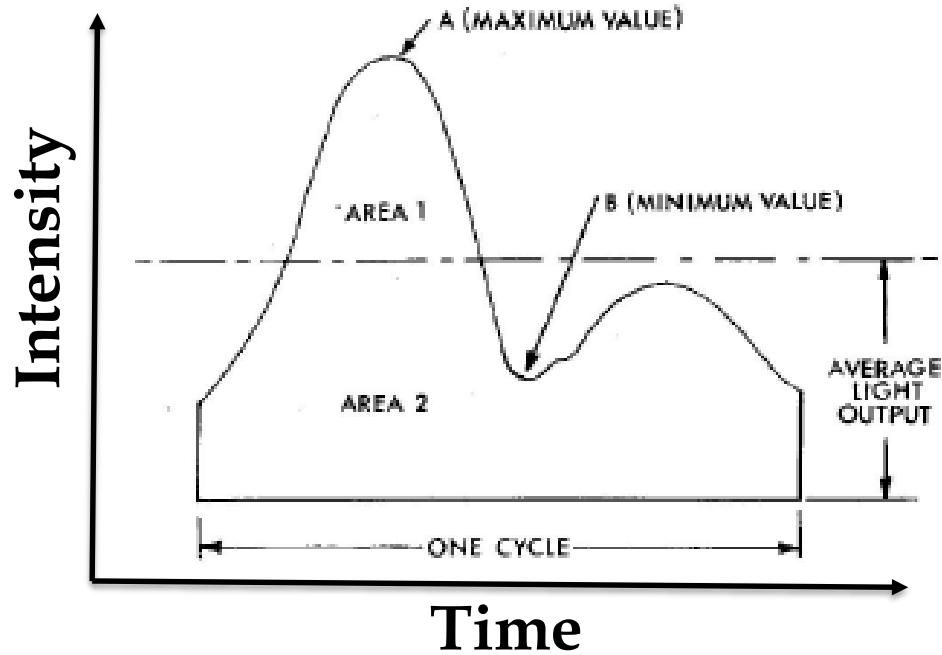
Definition of flicker

- Flicker is the fast, periodic change of the intensity of a light source
- It is an inherent problem of all kind of electrical light sources
- Alternating current (50 Hz) from the electric grid causes flicker of light sources
- Incandescent lamps show solely little flicker due to thermal inertia of glowing tungsten



4.2.9 Flicker Measurements

Quantification of flicker



$$\text{Percent Flicker} = \frac{\text{Max Value} - \text{Min Value}}{\text{Max Value} + \text{Min Value}} \cdot 100\% = \frac{A - B}{A + B} \cdot 100\%$$

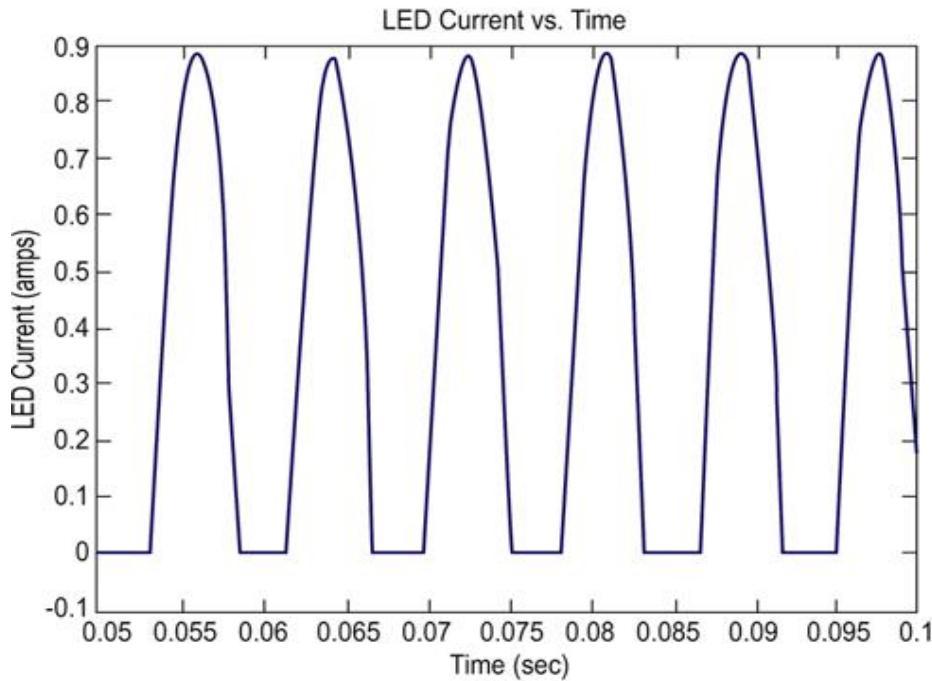
$$\text{Flicker index} = \frac{\text{Area above Mean}}{\text{Total area}} = \frac{\text{Area 1}}{\text{Area 1} + \text{Area 2}}$$

M. Rea, IESNA Lighting Handbook, 9th Edition, (2000)

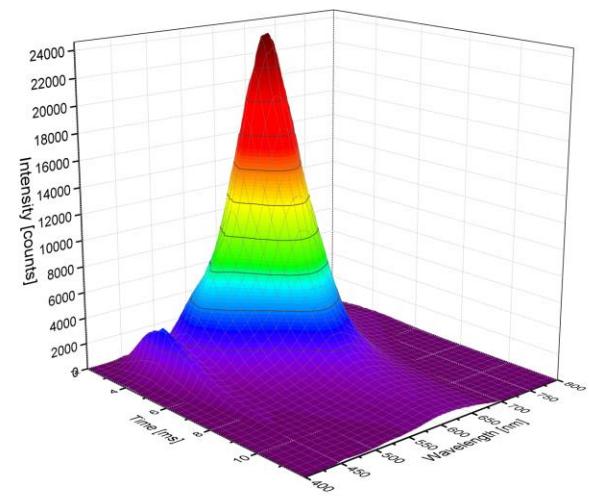
4.2.9 Flicker Measurements

Discharge lamps and particularly LEDs show strong flicker

- Flicker with 100 Hz for 50 Hz alternating current
- Electronic driver are able to flatten the current and intensity curve
- Alternative solution: Phosphors with long decay time > 10 ms



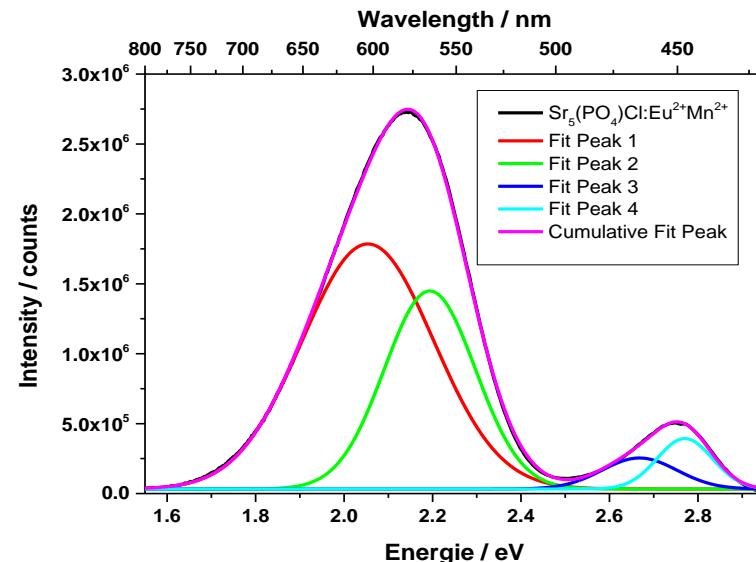
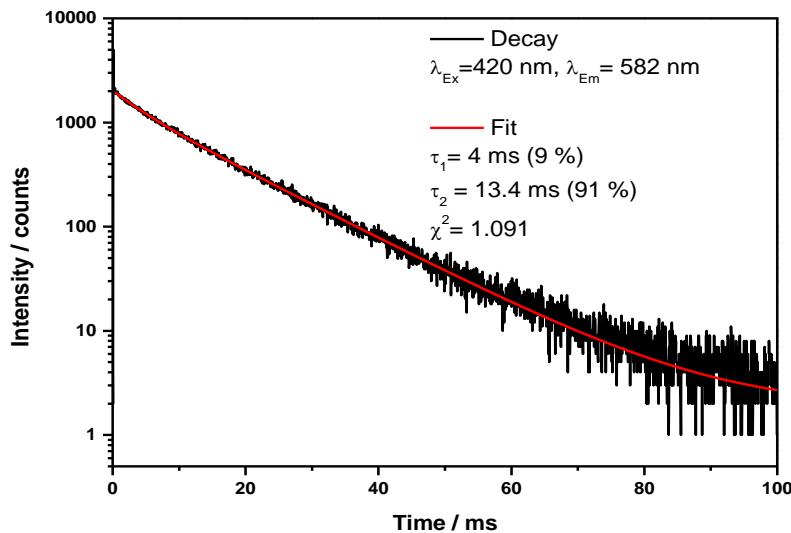
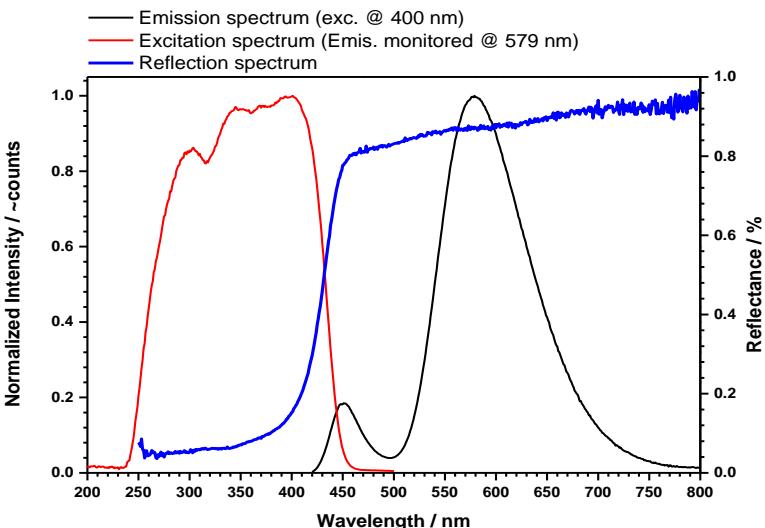
Example: Osram Parathom E14
Percent Flicker @ 100 Hz = 96%
Flicker Index = 0.45



4.2.9 Flicker Measurements

Example: $\text{Sr}_5(\text{PO}_4)_3\text{Cl}:\text{Eu}^{2+}\text{Mn}^{2+}$

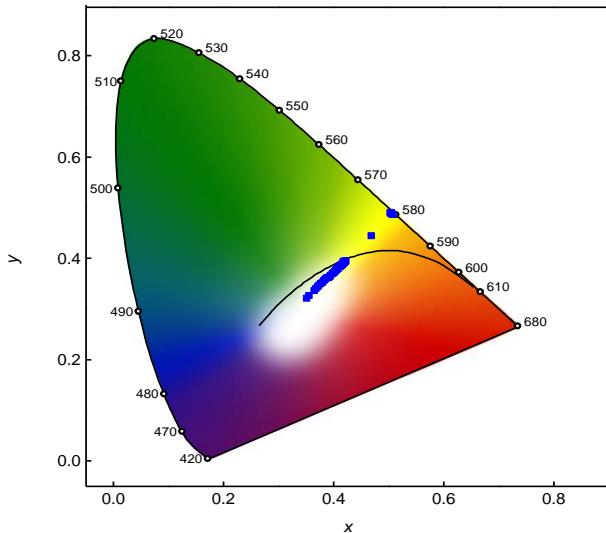
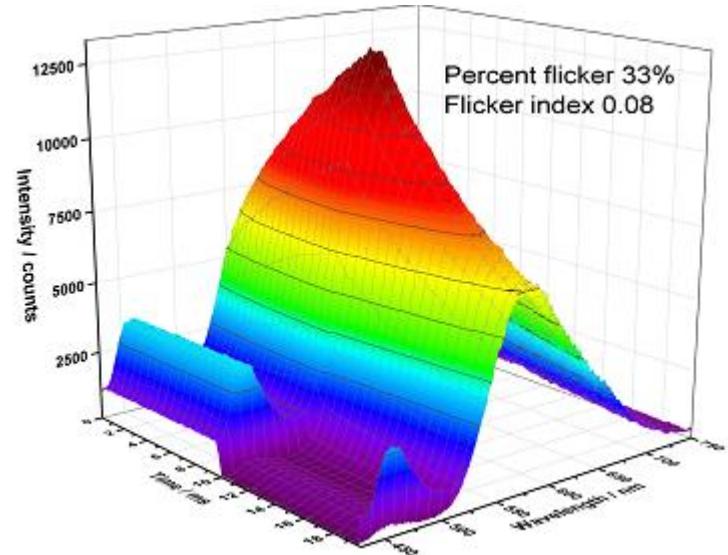
- Excitation 280 – 425 nm
- Emission 579 nm
 - Sr1 451 nm ($\text{CN} = 9$)
 - Sr2 579 nm ($\text{CN} = 7$)
- Decay time $\sim 12.5 \text{ ms}$



4.2.9 Flicker Measurements

Example: $\text{Sr}_5(\text{PO}_4)_3\text{Cl}:\text{Eu}^{2+}\text{Mn}^{2+}$

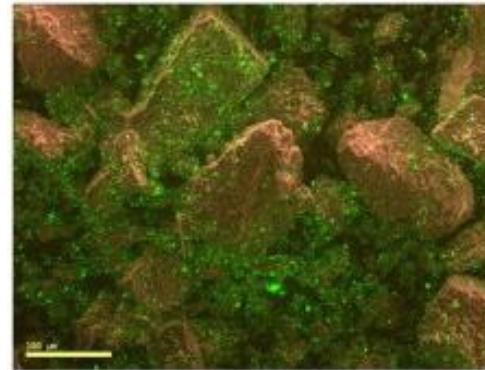
- Percent Flicker 33%
- Flicker Index 0.08
- Color point shifts from white to yellow
- Color point (average) :
 - $x = 0.453; y = 0.430; \text{LE} = 378 \text{ lm/W}_{\text{opt}}$
- CCT = 2916 K



4.2.10 X-ray and Cathode Luminescence

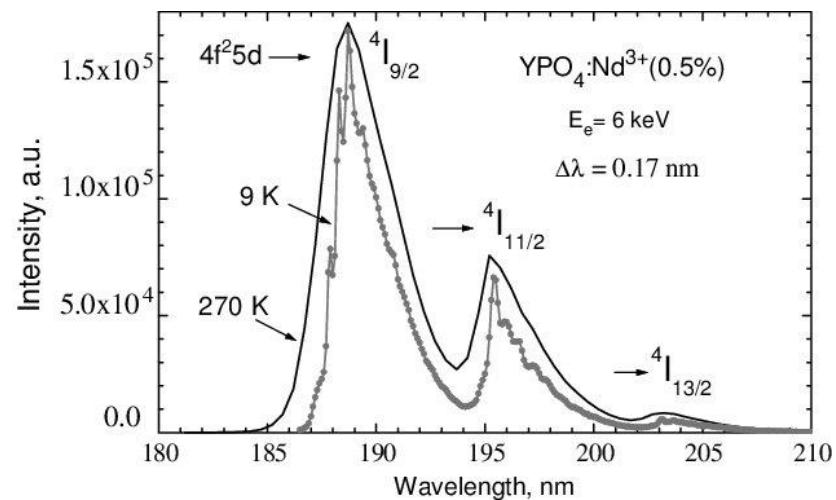
Excitation with high-energy radiation

- Unspecific excitation of all luminescent species (defects, band transitions, activators)
- Impurities in the ppm range can be determined qualitatively (e. g. Gd^{3+} in Y_2O_3)
- The sample can be excited with a spatial resolution in the μm range
- VUV or UV-C emission can be observed, such as Nd^{3+} bands at 190 nm



RGB image revealing the distribution of the API (green) and the lactose (brown) using CL spectroscopic imaging

Lit.: DOI: 10.1016/j.ejps.2011.10.017



Lit.: DOI: 10.1016/S0168-9002(02)00749-0

4.3 Reflection Spectroscopy

4.3 Reflection Spectroscopy

4.3.1 Physical Basis

4.3.2 Reflection Spectrometer

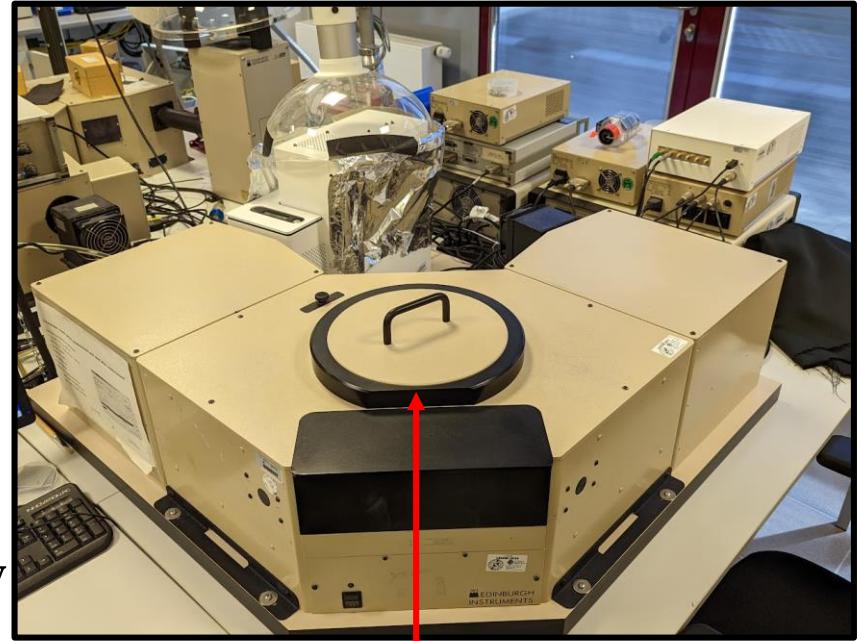
4.3.3 The Integrating Sphere

4.3.4 VUV Reflection

4.3.5 Band Gap Determination

4.3.6 The Kubelka-Munk Function

4.3.7 Application of Reflection Spectroscopy

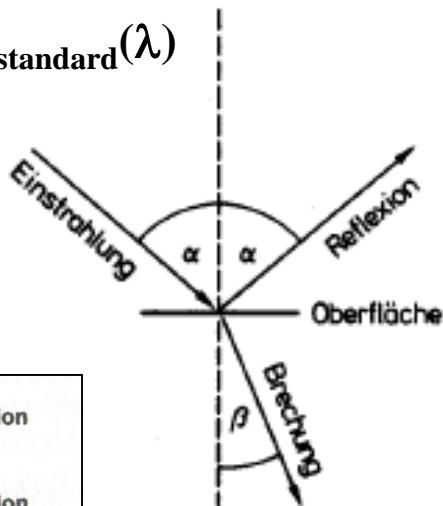


4.3.1 Physical Basis

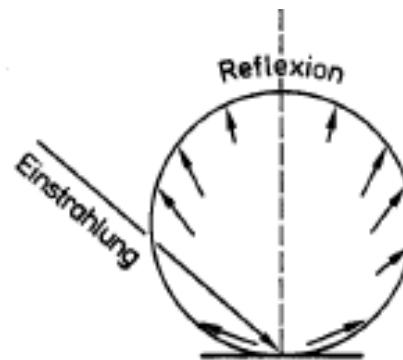
Types of Reflection

Measurement is made relative to a white standard, e.g. BaSO₄, CaCO₃, or Teflon powder

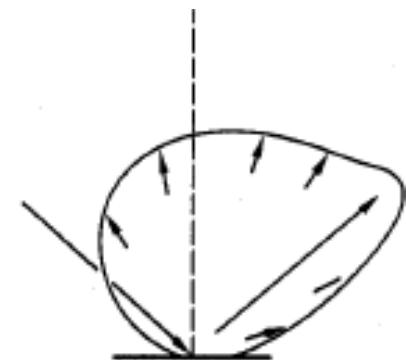
$$R_{\text{diffuse}}(\lambda) = I_{\text{sample}}(\lambda)/I_{\text{white standard}}(\lambda)$$



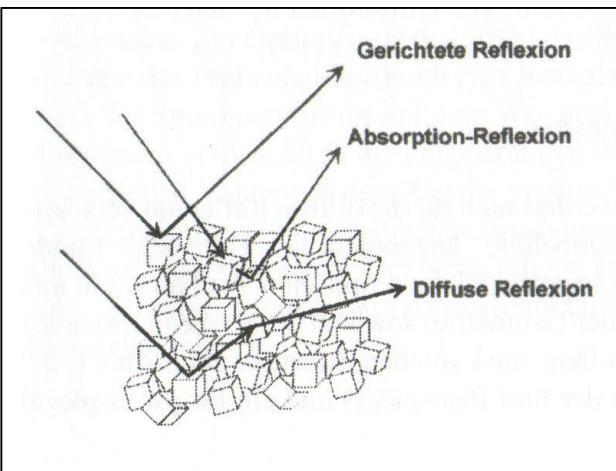
Regular
reflection



Diffuse
reflection



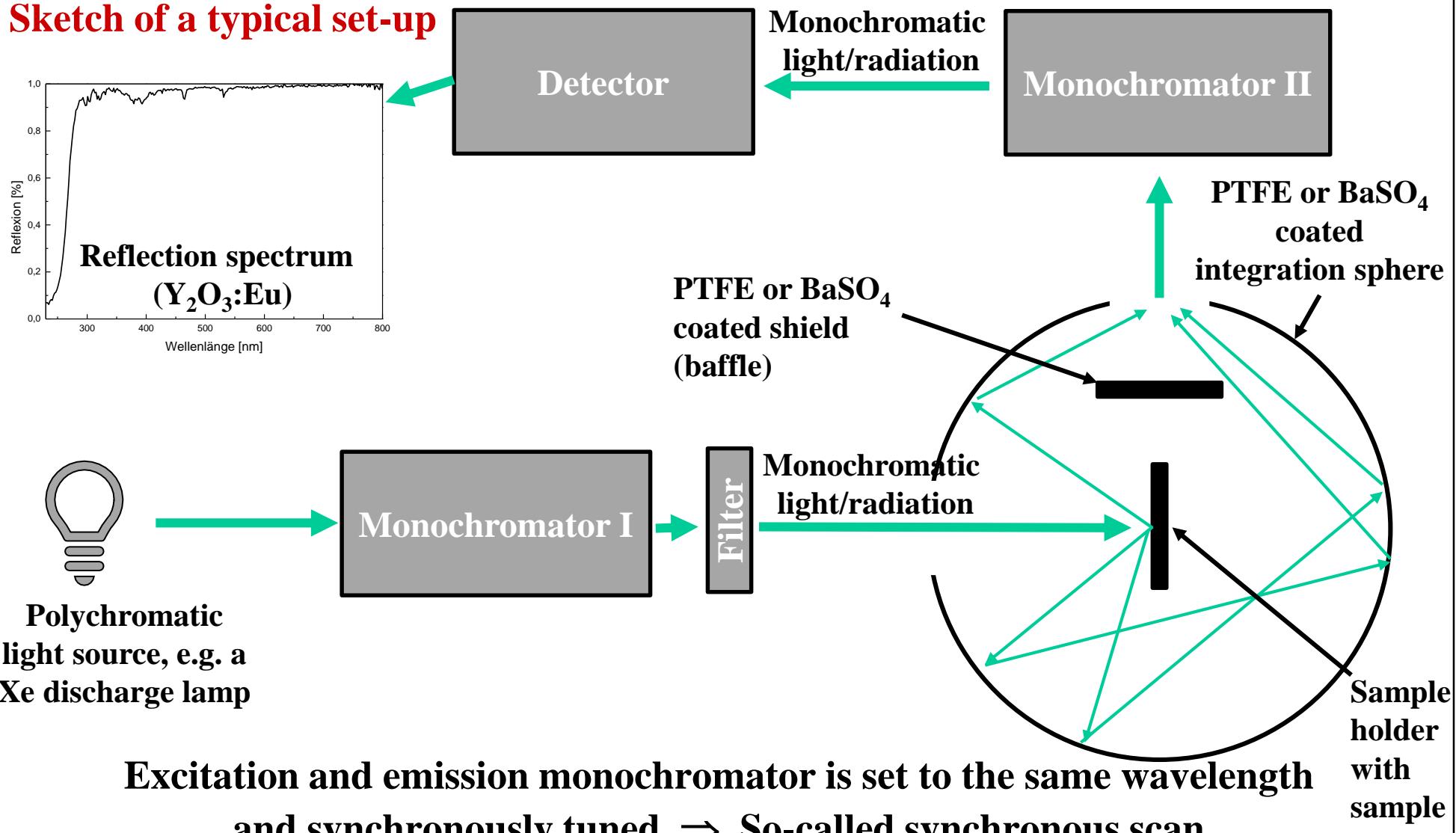
Overlapping
of two
processes



Lit.: D. Oelkrug in H. Naumer, W. Heller, Untersuchungsmethoden in der Chemie, Thieme, 3rd edition, 1997

4.3.2 Reflection Spectrometer

Sketch of a typical set-up



4.3.3 The Integrating Sphere

Image of an Ulbricht Sphere

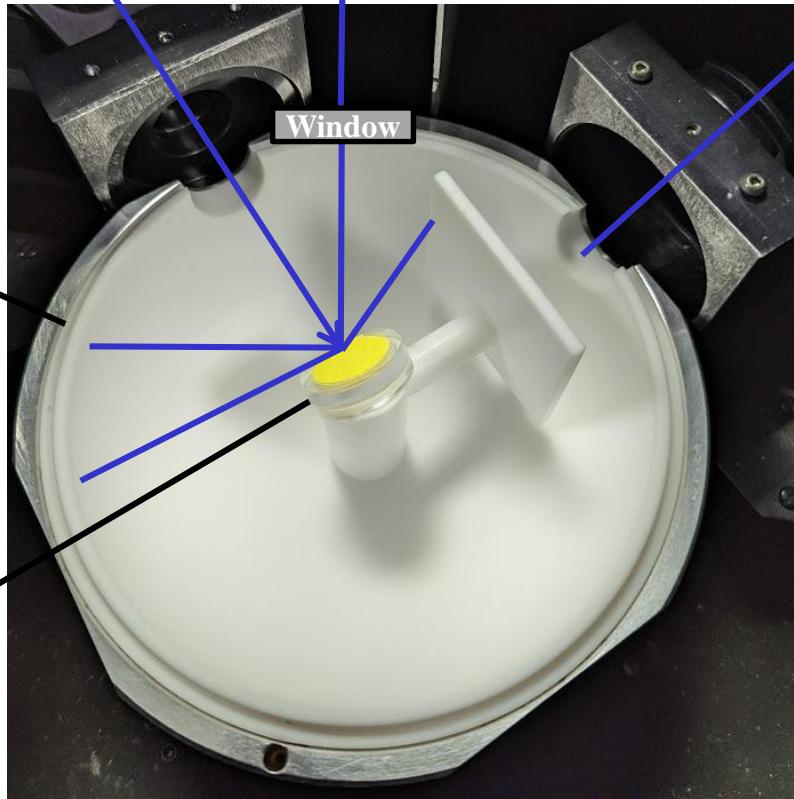
Light from monochromator 1

Coated
integrating
sphere
(BaSO_4)

Sample holder
with YAG:Ce

Optional output
to analysis the
specular reflexes

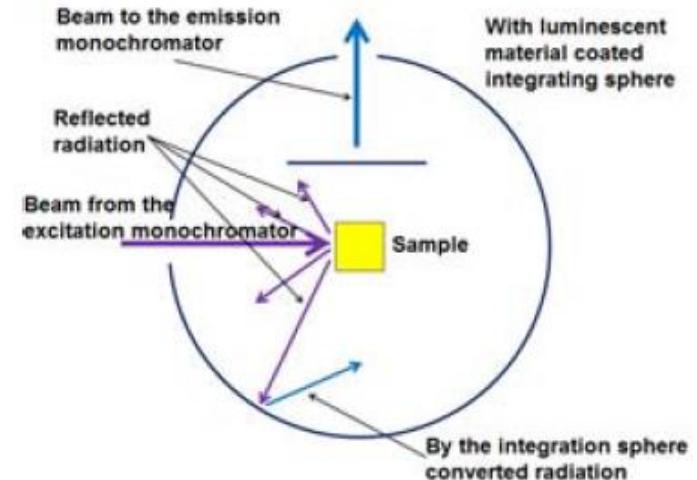
to mono-
chromator 2



4.3.4 VUV Reflection

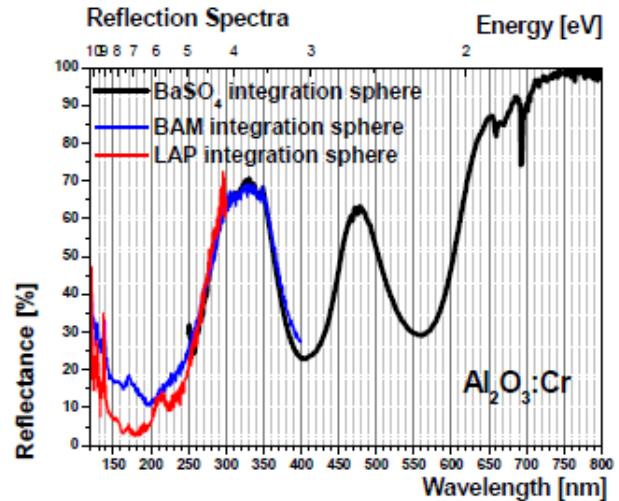
Specific Challenges in the VUV Range

- PMTs cannot detect VUV radiation
- Teflon and BaSO_4 absorb VUV ($< 200 \text{ nm}$)
→ Scintillation detector + LiF or MgF_2 coating (hygroscopic!)



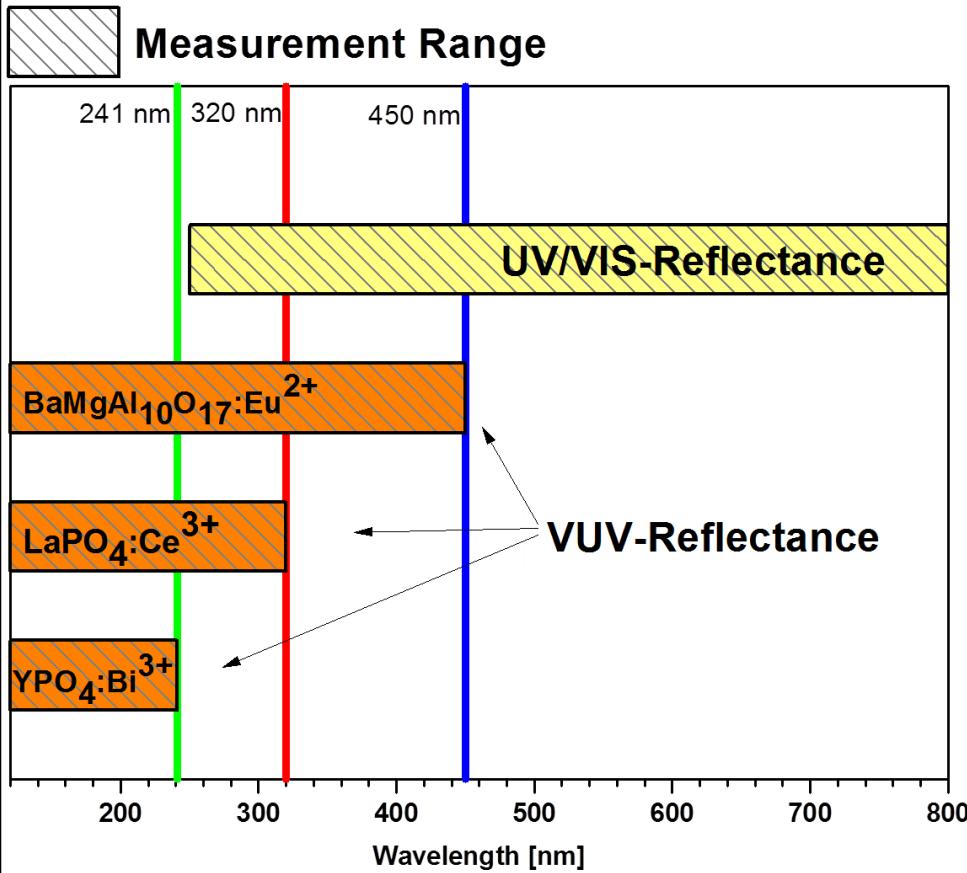
Alternative: Coating of the integration sphere with a phosphor which converts the VUV to visible radiation

Drawbacks: A phosphor coating is more prone to contamination than a teflon sphere. The converter must not emit in the same range as the sample!



4.3.4 VUV Reflection

Specific Challenges in the VUV Range



BaMgAl₁₀O₁₇:Eu coated integration sphere under daylight or upon VUV exc.



LaPO₄:Ce coated integration sphere

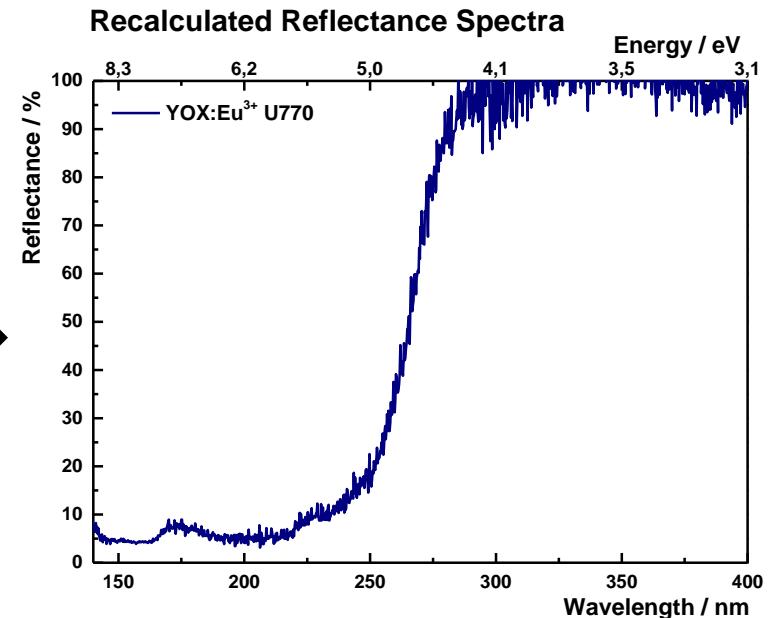
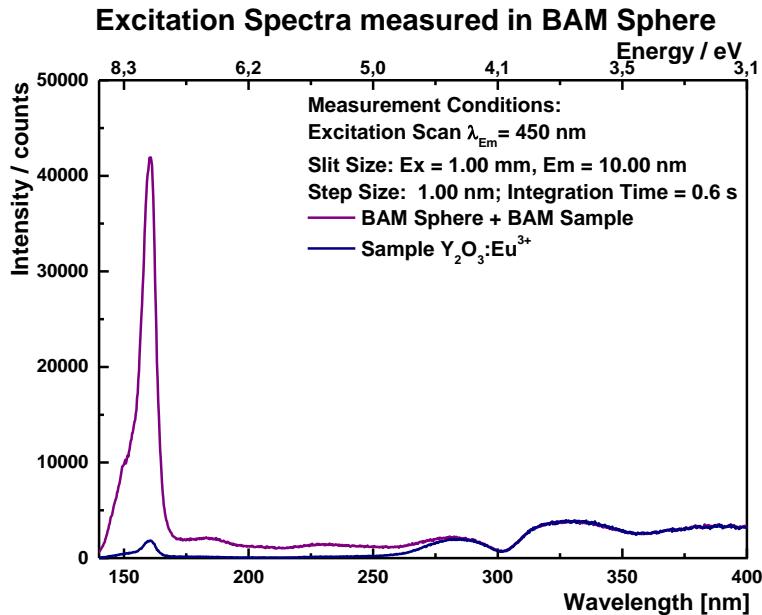
Reference

B. Herden, A. Katelnikovas, S. Möller, H. Winkler, R. Petry, H.-J. Meyer, T. Jüstel, Powder Reflection Spectroscopy in the Vacuum UV Range, D. Enseling, J. Appl. Spectrosc. 81 (2014) 341-346

4.3.4 VUV Reflection

Procedure for recording a diffuse reflection spectrum $R_{\text{diffuse}}(\lambda)$ in the VUV range

1. Irradiation of the sample by VUV photons outcoupled from the evacuated excitation monochromator into a N₂ flushed BaMgAl₁₀O₁₇:Eu²⁺ (BAM:Eu) coated Ulbricht sphere
2. Conversion of reflected VUV photons into blue light by BAM:Eu coating and integration over the sphere
3. Light output is collected by a PMT via the emission monochromator
4. Determination of reflection spectrum by dividing I(Sample)/I(BAM:Eu)



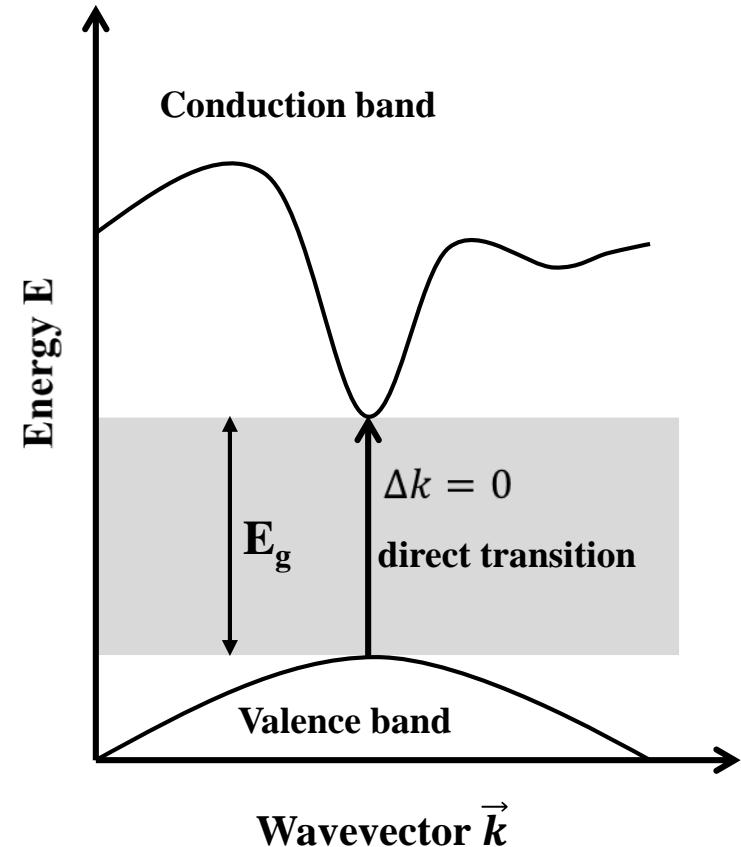
4.3.5 Band Gap Determination

Motivation

- The optical band gap is of large interest to judge potential application areas of optical materials
- It indicates the difference in energy between the top of the valence band filled with electrons and the bottom of the conduction band devoid of electrons

Material class	E_g [eV]
Metals	0.0
Semiconductors	0.0 - ~3.0
Insulators	> ~3.0

$$\mathbf{p} = \hbar \cdot \mathbf{k}$$

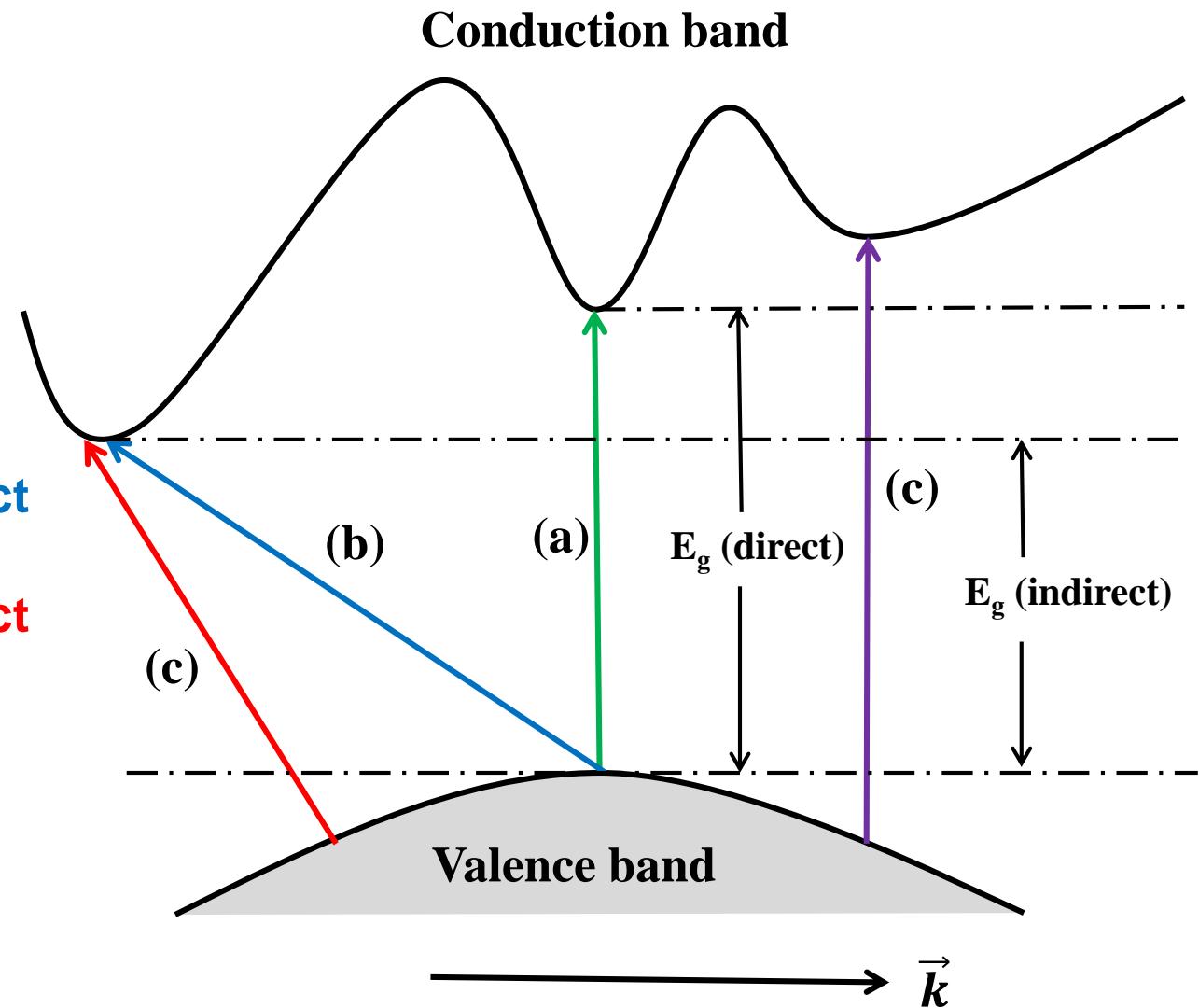


4.3.5 Band Gap Determination

Motivation

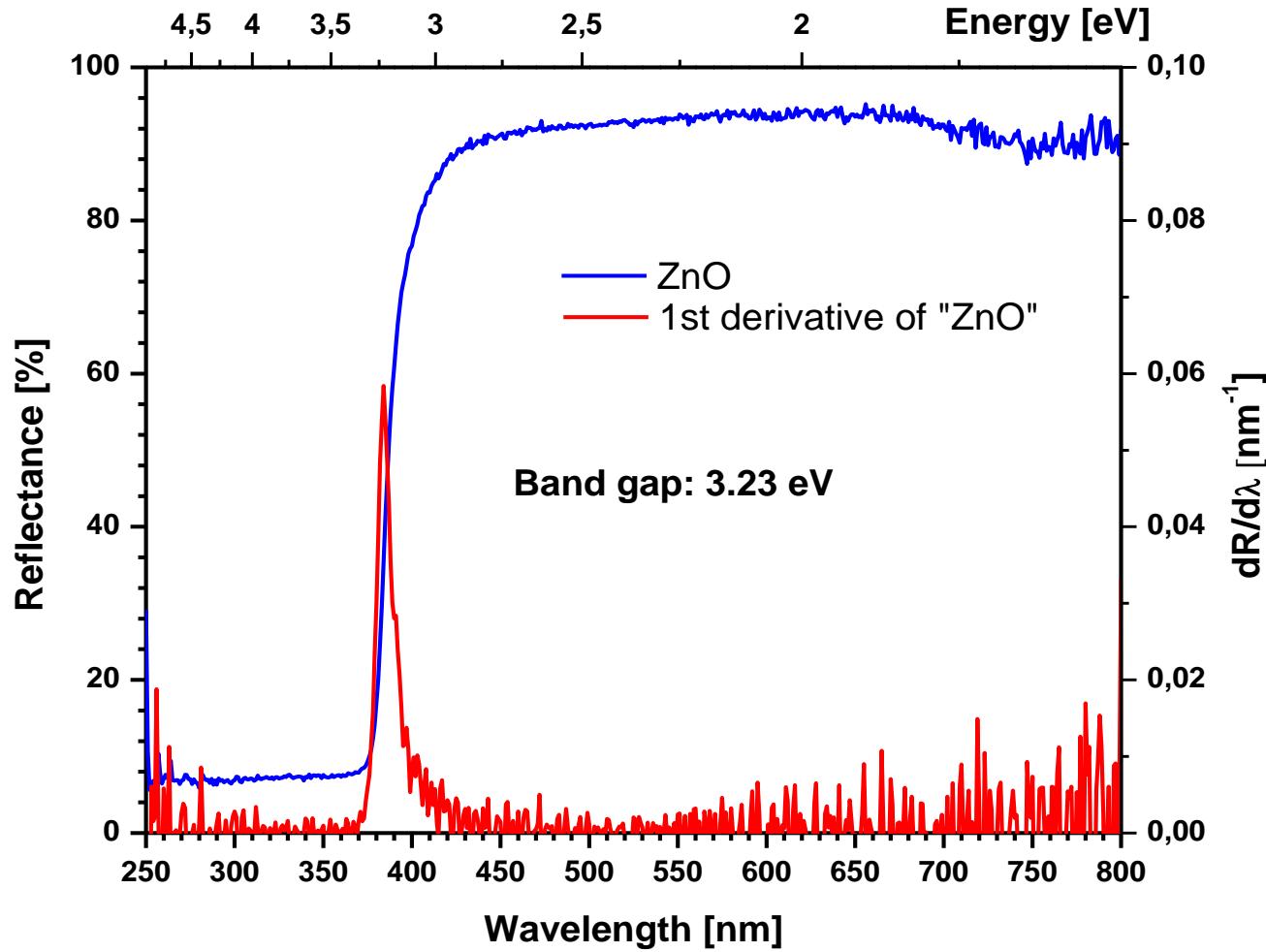
Types of transitions

- (a) Allowed direct
- (b) Allowed indirect
- (c) Forbidden direct
- (d) Forbidden indirect



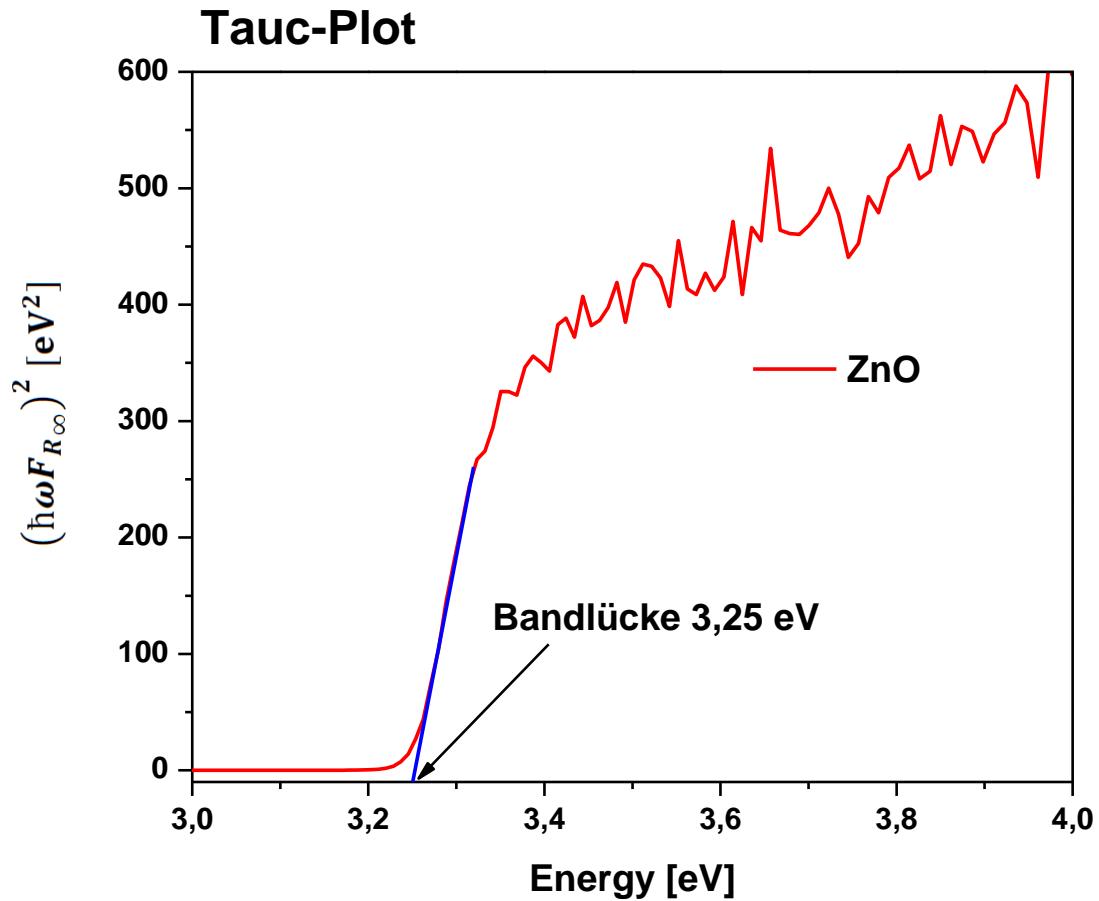
4.3.5 Band Gap Determination

Method by the turning point



4.3.5 Band Gap Determination

Method by the Tauc plot



Tauc's equation

$$(\alpha h\nu)^{1/n} = A(h\nu - E_g)$$

For allowed direct transitions..... $n = 1/2$

For forbidden direct transitions..... $n = 3/2$

For allowed indirect transitions..... $n = 2$

For forbidden indirect transitions..... $n = 3$

Lit.: J. Tauc (F. Abeles ed.), Optical Properties of Solids, North-Holland (1972)

4.3.6 The Kubelka-Munk Function

Diffuse Reflection of Powder Samples Depend on Ratio of Absorption to Scattering

Under certain assumptions, i.e.

- Diffuse radiation
- Suppression of the regular reflection
- Approximately spherical particles with $d \gg$ than the wavelength of the incident light
- Weak bands in a non-absorbing matrix
- No overlapping of bands
- Infinite layer thickness

there is a mathematical relationship between absorption, scattering and reflection

⇒ The Kubelka-Munk-function:

$$F(R_\infty) = \frac{A}{S} = \frac{(1-R_\infty)^2}{2 \cdot R_\infty}$$

S = Scattering coefficient
R_∞ = Diffuse scattering
A = Absorption = ε·c

An ideally black surface $R_\infty = 0$ is therefore difficult to realize, because for $R_\infty = 0$ applies A/S = ∞, i.e. extremely high absorption or virtually no scattering

The scattering coefficient S can be estimated from the particle size of the powder:

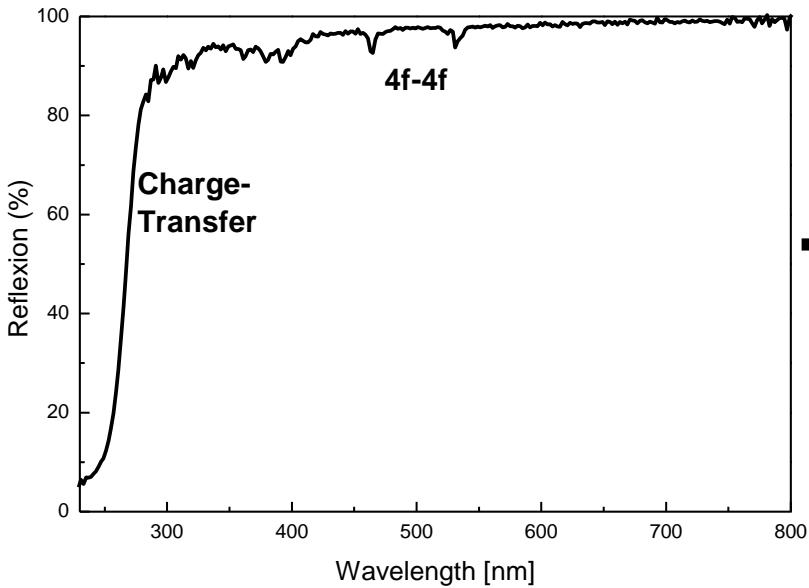
$$S \sim 1/\text{grain size} \sim 1/d_{50} [\text{cm}^{-1}]$$

4.3.6 The Kubelka-Munk Function

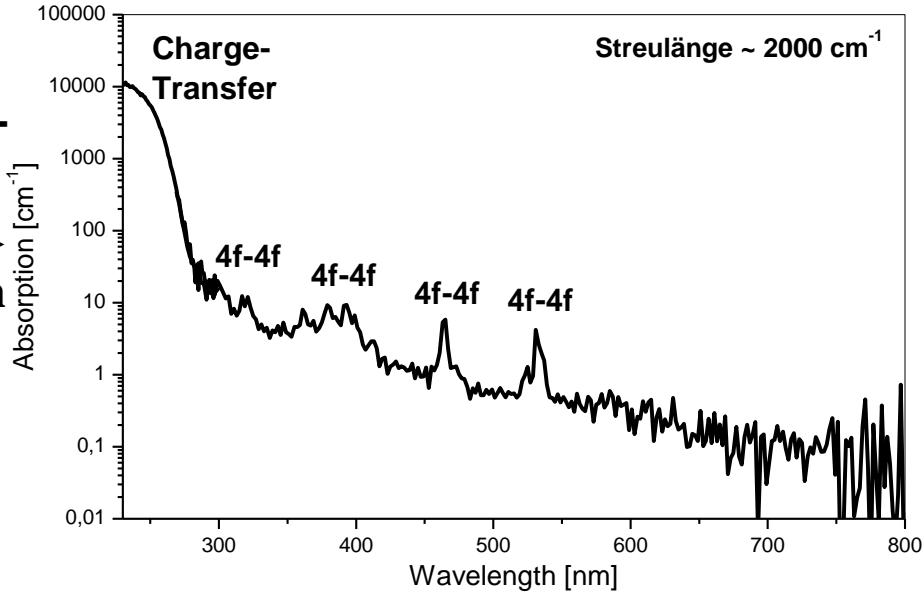
Determination of Qualitative Absorption Spectra of Powder Samples

Boundary condition: $T = 0$, „semi-infinitely thick layer“ $\Rightarrow R + A = 1$

Example: Phosphor $Y_2O_3:Eu^{3+}(5\%)$ with $d_{50} = 5 \mu m \Rightarrow S \sim 0.2 \mu m^{-1} = 2000 \text{ cm}^{-1}$



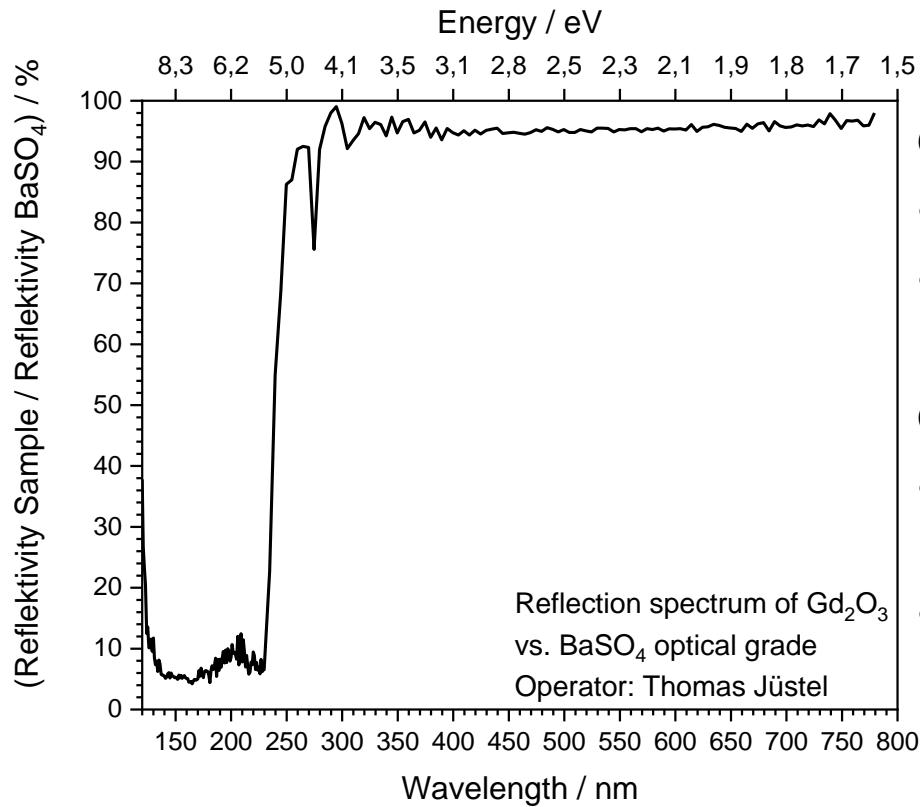
Kubelka-
Munk
Function



4.3.7 Application of Reflection Spectroscopy

Determination of Optical Transitions of Powder Samples

Example: Gd_2O_3



Optical transitions

- Band-to-band transition $\sim 225 \text{ nm}$ (5.5 eV)
- $[\text{Xe}]4\text{f}^7 ({}^8\text{S}_{7/2}) - [\text{Xe}]4\text{f}^7 ({}^6\text{P}_{5/2}) \sim 275 \text{ nm}$

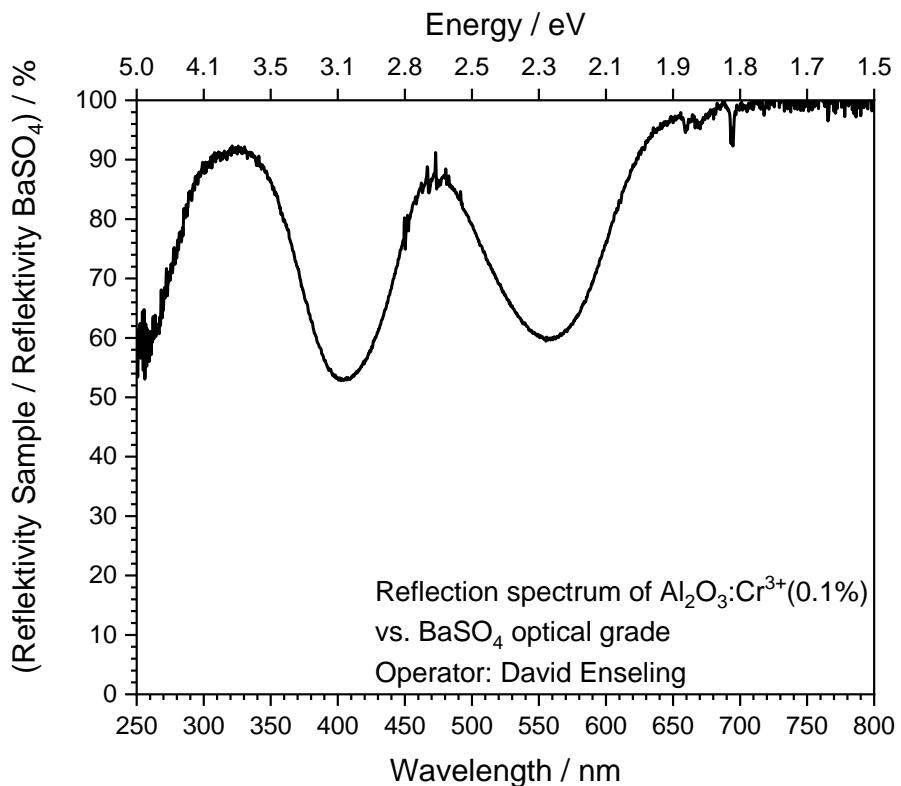
Optical band gap according to literature

- 5.5 eV (IOP Conf. Series: Journal of Physics: Conf. Series 917 (2017) 062001)
- 5.4 eV (J. All. Comp. 242 (1996) 41)

4.3.7 Application of Reflection Spectroscopy

Determination of Optical Transitions of Powder Samples

Example: $\text{Al}_2\text{O}_3:\text{Cr}^{3+}$



Optical transitions

- Band-to-band transition < 200 nm
- Cr^{3+} : $[\text{Ar}]3d^3$ (4F) to $[\text{Ar}]3d^3$ (4F , 4P)
 - $^4A_2 - ^4T_2$ 405 nm band
 - $^4A_2 - ^4T_1$ 560 nm band
 - $^4A_2 - ^2E$ 694 nm lines

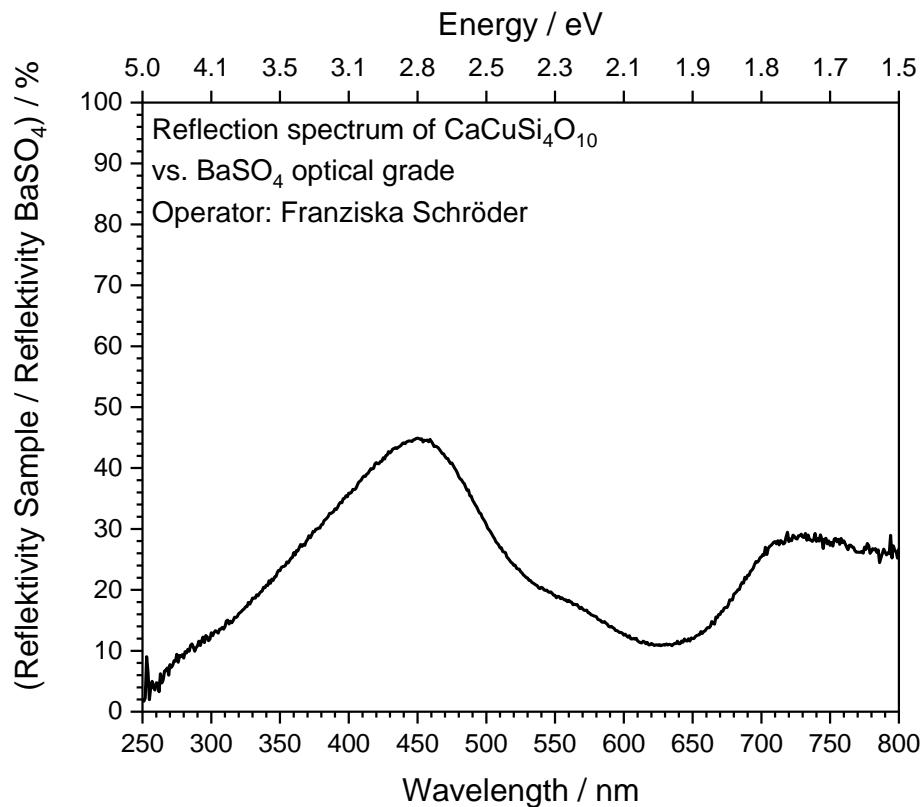
Optical band gap

- 8.7 eV (Appl. Phys. Lett. 55 (1989) 1955)
- 8.8 eV (Chem. Phys. Lett. 637 (2015) 172)
- 8.9 eV (JETP Letters 85 (2007) 165)

4.3.7 Application of Reflection Spectroscopy

Determination of Optical Transitions of Inorganic Pigments

Example: $\text{CaCuSi}_4\text{O}_{10}$



Optical transitions

- Band-to-band transition $\sim 340 \text{ nm (3.6 eV)}$
- Cu^{2+} : $[\text{Ar}]3\text{d}^9 \rightarrow [\text{Ar}]3\text{d}^9$ transitions

$${}^2\text{E}_g \rightarrow {}^2\text{B}_{1g} \quad 500 \text{ nm}$$

$${}^2\text{A}_{1g} \rightarrow {}^2\text{B}_{1g} \quad 625 \text{ nm}$$

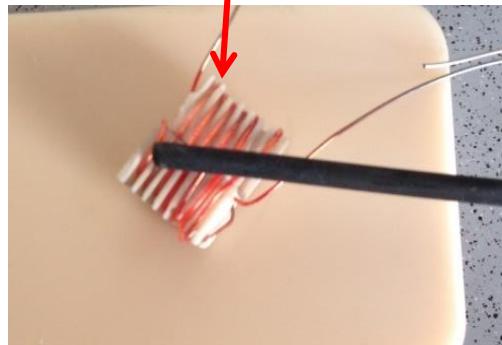
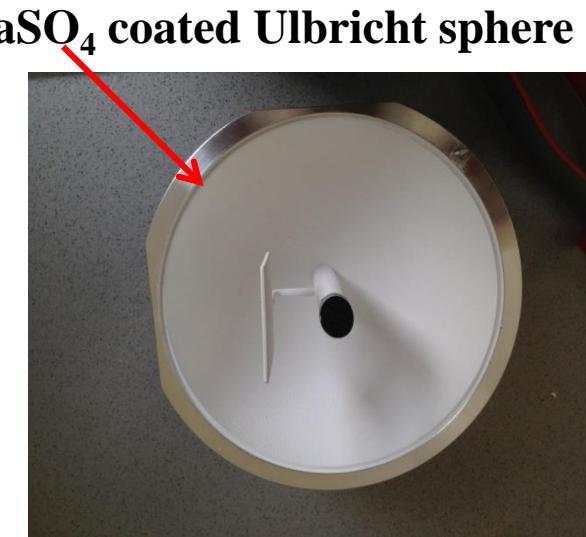
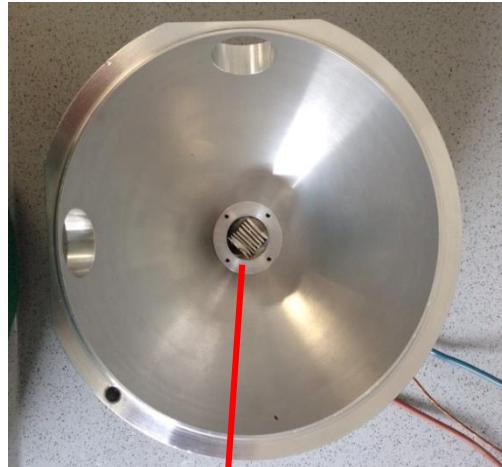
$${}^4\text{A}_{2g} \rightarrow {}^2\text{T}_{1g} \quad 750 \text{ nm}$$

Optical band gap according to literature

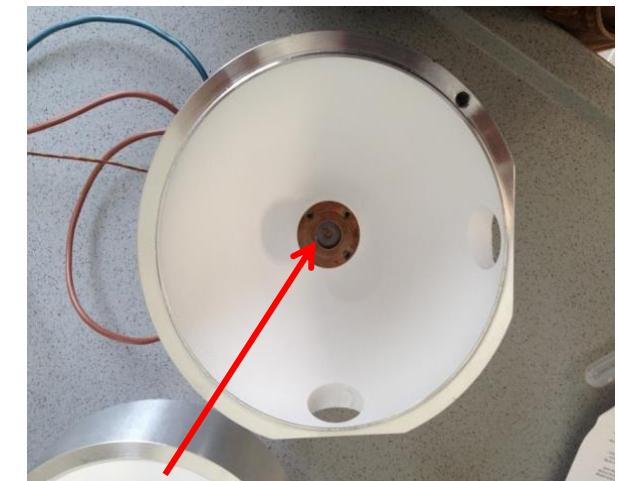
- 3.6 eV (J. Phys. Chem. Lett. 7 (2016) 399)

4.3.7 Application of Reflection Spectroscopy

Thermochromic Properties of Inorganic Pigments



Heating element
up to 300 °C



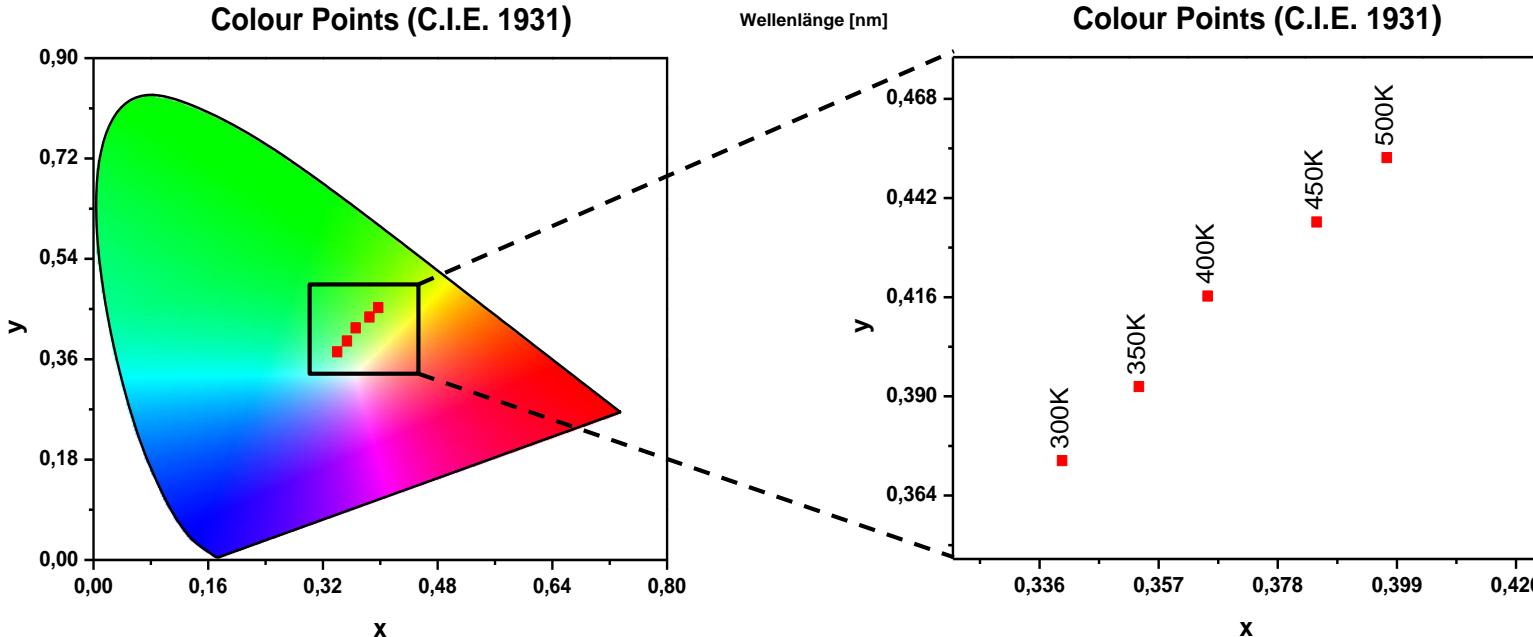
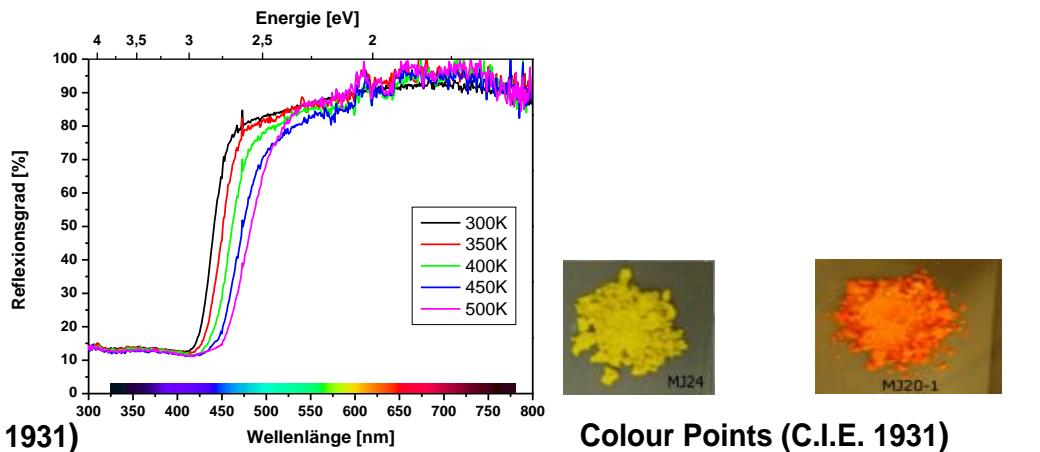
Cu sample holder

4.3.7 Application of Reflection Spectroscopy

Thermochromic Properties of Inorganic Pigments

Example:

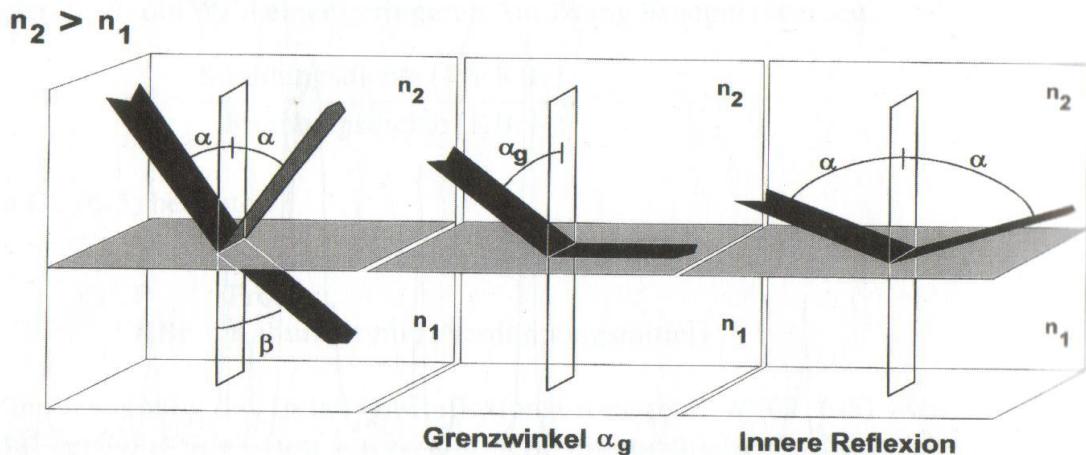
Yellow pigment BiVO_4



4.3.7 Application of Reflection Spectroscopy

Attenuated Total Reflection (ATR)

Phenomenon of total reflection



But:

The total reflected radiation also penetrates into the optically less dense medium (Goos, Hänchen 1947) and can be weakened by absorption
→ ATR-spectroscopy

Note:

The penetration depth of radiation is wavelength dependent:
The larger the wavelength, the more intensive bands appear,
furthermore their positions are changed.
→ ATR correction is required (e.g. OPUS software from Bruker)

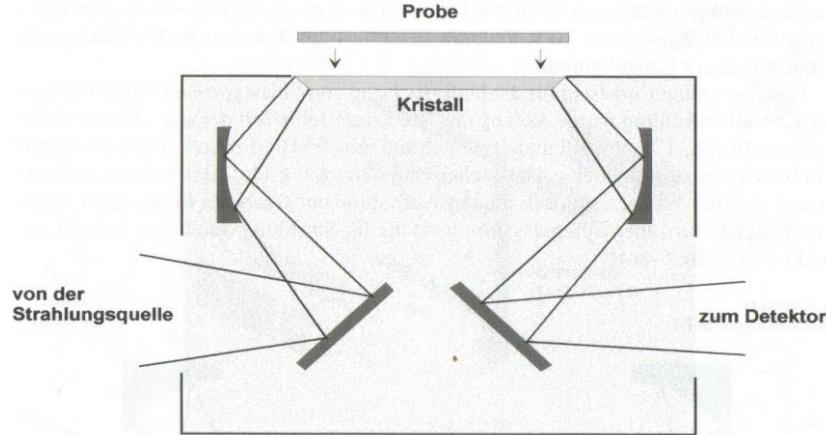
$$d_p = \frac{\lambda}{2\pi \sqrt{(n_K^2 \sin^2 \Theta - n_P^2)}}$$

4.3.7 Application of Reflection Spectroscopy

Attenuated Total Reflection (ATR)



ATR Unit for FT-IR Spectroscopy



Advantages

- No sample preparation is necessary for powders or thin crystals and liquids

Disadvantages

- Intimate contact between the material and crystal is necessary
- Crystal may be neither chemically or mechanically damaged
- Only room temperature
- Range of application is limited by the crystal material, e.g. 17000 - 650 cm⁻¹ for ZnSe

